

## Unit-1 Quantum Mechanics.

### Quantum Mechanics:- (Microscopic)

" Science dealing with the behaviour of Matter & light on atomic & subatomic scale."

It attempts to describe the Properties of Molecules atoms and their constituents (Electron, Protons, Neutrons and other subatomic Particle).

### Classical Mechanics:-

Branch of Physics deals with the motion of object smaller as well as large objects.

OR Classical Mechanics deals with the motion of bodies under the influence of force.

#### Classical Mechanics

- I) Deals with Macroscopic objects
- II) Based on Newton's Law of Motion.
- III) In classical mechanics Behaviour of Particle can be completely known.
- IV) classical mechanics deals with certainties

#### Quantum Mechanics

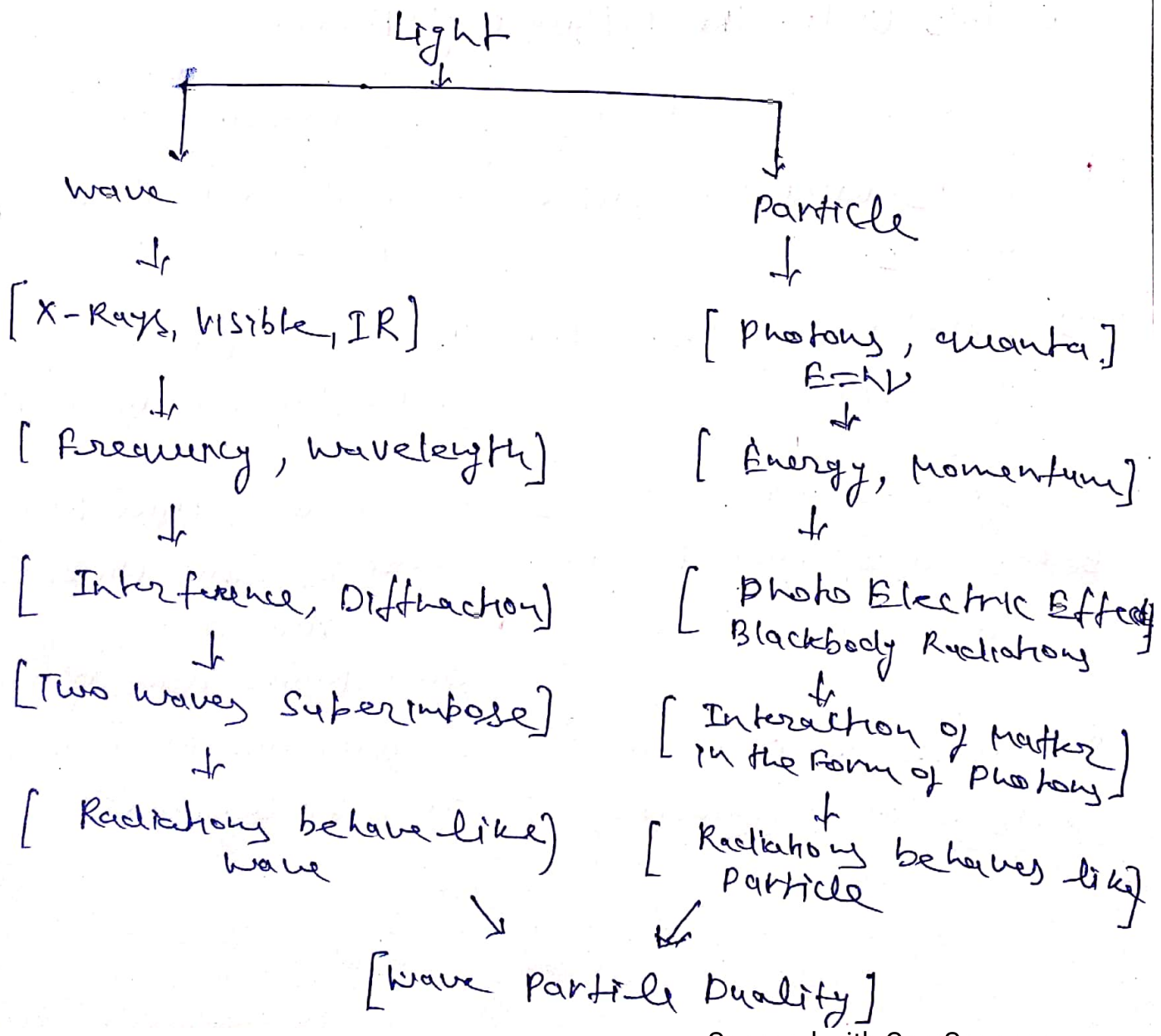
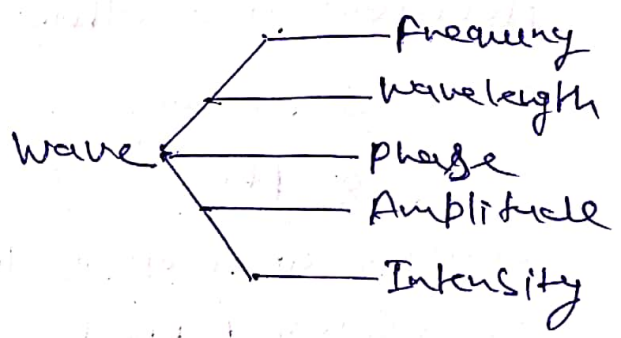
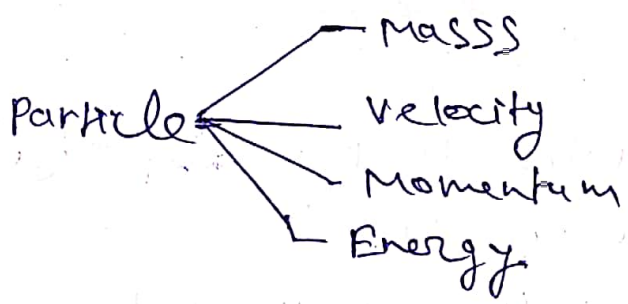
- I) Deals with Microscopic objects
- II) Based on Schrodinger wave Equation.
- III) In Quantum Mechanics there is some Uncertainty in determination of position and momentum of Particle.
- IV) Quantum Mechanics deals with probabilities.

# Wave Particle Duality :-

"Wave Particle Duality Describe the properties of both photons and subatomic particles to exhibit the properties of both waves & particles"

OR

"It is the concept of Quantum mechanics that every particle or entity may be described as a particle or wave"



## De Broglie Hypothesis - (4)

Light wave can act as a wave sometimes and as a Particle at other times this is known as de Broglie Hypothesis.

### Matter Waves :-

According to de Broglie Hypothesis any moving Particle is associated with a wave. The wave associated with a Particle known as de Broglie waves or matter waves.

The wavelength,  $\lambda$  of matter waves associated with a Particle moving with velocity  $v$  is inversely proportional to the magnitude of momentum of Particle. Thus

$$\boxed{\lambda = \frac{h}{mv}} = \frac{h}{p} \quad (p = mv)$$

The relation  $\lambda = \frac{h}{mv}$  is known as de-Broglie Equation and wavelength  $\lambda$  is called de-Broglie wavelength.

### De-Broglie wavelength associated with an accelerated charge Particle :-

If a charge Particle, say an electron is accelerated by a potential difference of  $V$  volts, then its Kinetic Energy is given by  $K.E = eV$

$$\frac{1}{2} mv^2 = eV$$

$$v = \sqrt{\frac{2eV}{m}}$$

Then electron wavelength is given by - (5)

$$\lambda = \frac{h}{mv} = \frac{h}{m} \sqrt{\frac{m}{2eV}}$$

$$\lambda = \frac{h}{\sqrt{2emV}}$$

De Broglie wavelength Expressed in term of Kinetic Energy:-

If a particle has K.E, Kinetic Energy

$$K.E = \frac{1}{2} mv^2$$

$$= \frac{1}{2} mv^2 \times \frac{m}{m}$$

$$= \frac{m^2 v^2}{2m} = \frac{p^2}{2m}$$

$$p = \sqrt{2m(K.E)}$$

(Substituting value of  $h/mv$ )  $\lambda = \frac{h}{p} = \frac{h}{\sqrt{2m(K.E)}} = \frac{12.28 \text{ \AA}}{\sqrt{V}}$

De Broglie wavelength associated with a particle in thermal Equilibrium:-

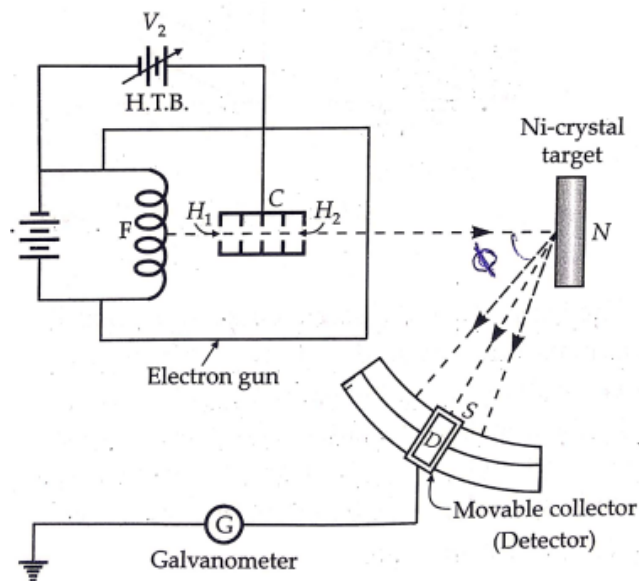
For a particle in thermal Equilibrium at temperature  $T$ , then their Kinetic Energy is given by -

$$K.E = \frac{3}{2} kT$$

$$\lambda = \frac{h}{\sqrt{2m(K.E)}} = \frac{h}{\sqrt{3m kT}}$$

## Davisson & Germer Experiment:-

( Experiment to show the Existence of Matter waves)  
 Davisson & Germer shows that electron beam can undergo diffraction when passed through an atomic crystal. It confirms the wave nature of electron and De-Broglie wave length.



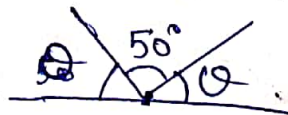
- \* In Experimental Arrangement of Davisson & Germer
- The electron beam is generated from a hot Tungsten filament F. The filament is connected to variable voltage source (V).
- The electron emerges through an opening and falls normally on the surface of Nickel crystal.
- The intensity of scattered electron can be measured by detector D in different direction. The detector can be moved in different directions.

Note:-

$$\theta + 50 + \theta = 180^\circ$$

$$2\theta = 130^\circ$$

$$\theta = 65^\circ$$

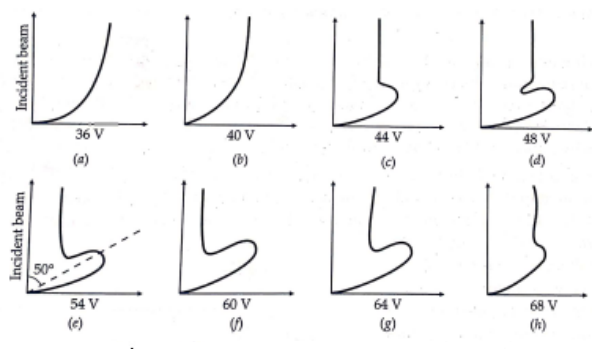


~~Answers~~

\* First of all accelerating voltage  $V$  is given a low value and detector on the circular scale moved to various position and value of current was measured.

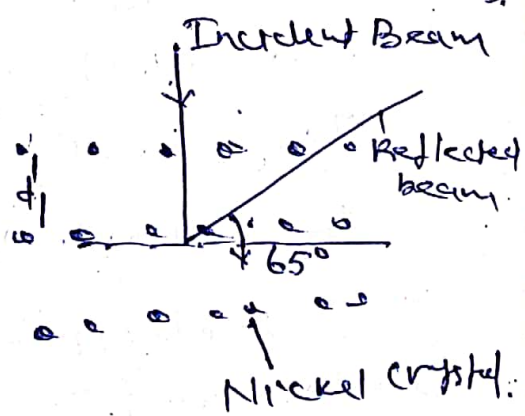
\* The current (which measures the intensity of diffracted beam) is plotted against angle  $\phi$ .

( $\phi$  is the angle between Incident beam on the crystal and the beam entering in the detector).



\* It is seen that bump begins to appear in the curve for 44-volt electron. With increasing potential ( $V$ ) the bump becomes most prominent (Maximum Intense) for 54 volt at  $\phi = 50^\circ$ .

\* As Nickel is made up of crystal using Bragg's law of Diffraction wave length ( $\lambda$ ) can be calculated



$$\lambda = 2d \sin \theta$$

$$\lambda = 2 \times 0.91 \text{ \AA} \times \sin 65^\circ$$

$$\lambda = 1.65 \text{ \AA} \quad \text{--- (1)}$$

( $d = 0.91 \text{ \AA}$ ,  $d$  is the interlayer spacing).

The wavelength of electron can also be calculated from accelerating potential  $V$  using de-Broglie equation.

$$\lambda = \frac{h}{\sqrt{2meV}}$$

Note:-

$$\lambda = \frac{h}{p}$$

$$p = mv$$

$$p^2 = m^2 v^2 = m \times m v^2$$

$$= m \times 2E \quad (K.E = \frac{1}{2}mv^2)$$

$$\Rightarrow p^2 = 2mE$$

$$E = \frac{p^2}{2m} \Rightarrow p = \sqrt{2mE}$$

$$\lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE}}$$

Hence 
$$\lambda = \frac{h}{\sqrt{2meV}} = \frac{6.6 \times 10^{-34}}{\sqrt{2 \times 9.1 \times 10^{-31} \times 54 \times 1.6 \times 10^{-19}}}$$

(Substituting)  
 $h = 6.6 \times 10^{-34}$   
 $m = 9.1 \times 10^{-31}$   
 $e = 1.6 \times 10^{-19}$

(Note  $E = 54 \text{ eV}$   
 $= 54 \times 1.6 \times 10^{-19} \text{ J}$ )

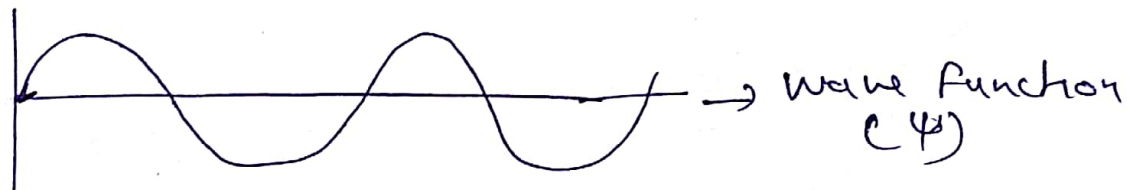
$$\lambda = 1.66 \text{ \AA} \quad \text{--- (2)}$$

It is seen that the values obtained experimentally using Bragg's equation and de-Broglie equation agreed well. Therefore Davisson - Germer conclude that electrons exhibit diffraction properties.

Wave Function:- Function describing the probability of each possible observation such as position & momentum.

Represented by  $(\psi) - \Psi_i$

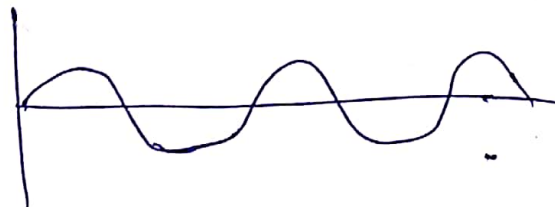
$\psi(x, t)$   
Position      time



Note- A wave is associated with a particle.

Characteristics:-

(I) If a wave has higher energy it rotate with high frequency.



(II) maximum Amplitude of a wave has maximum probability of finding the particle. ( $\psi^2$ )

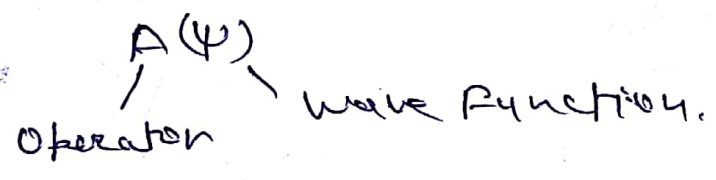
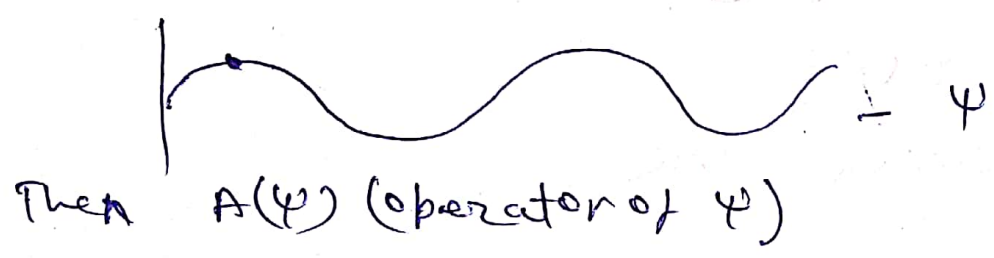


(III) Momentum of a wave is inversely proportional to wavelength.

$$p \propto \frac{1}{\lambda} \quad (p - \text{Momentum})$$

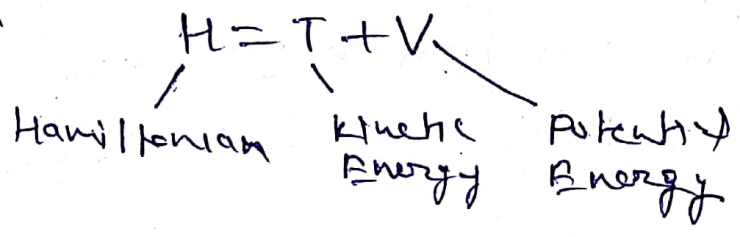
Operator: An operator is a mathematical operation on the wave function that produce another wave function.

Suppose  $\psi$  be any wave function



E.g - Hamiltonian operator (H)

Hamiltonian operator (H) represent total Energy of a system



Normalization of a wave function:

Normalization. when the probability of finding a particle is 1 then the function is said to be Normalized

$$\int \psi^* \psi dx = 1$$

Suppose ~~if~~ a function is not normalized

(1)

$$\int \psi^* \psi dx = N \quad (N \neq 1)$$

Then

$$\frac{1}{N} \int \psi^* \psi dx = 1$$

$$\int \frac{\psi^*}{\sqrt{N}} \frac{\psi}{\sqrt{N}} dx = 1$$

↓  
Normalized

~~It~~ means Probability of finding Particle is 1

( $\psi^* \psi$  - Represent Probability of finding Particle)  
(Probability density)

2.5 Normalize the wave function

$\psi(x) = 0$  outside the box of size  $l$

$\psi(x) = A \sin kx$   $0 < x < l$

where  $k = \frac{\pi}{l}$

Hint:  $\int_{-\infty}^{+\infty} |\psi|^2 dx = 1,$

Here  $\psi(x) = A \sin kx = A \sin \frac{\pi x}{l} = \int_0^l A^2 \sin^2 \frac{\pi x}{l} dx = 1$

From above expression,  $A = \sqrt{\frac{2}{l}}$

Then  $\psi_n = \sqrt{\frac{2}{l}} \sin \frac{\pi x}{l}$

[GGSIPU, May 2012 (4 marks)]

## Wave Function & Probability Interpretation:- (12)

Micro particles exhibit wave properties, quantity  $\psi$  represent a de broglie wave, This quantity ( $\psi$ ) is called Wave Function.

$\psi$  describe wave as a function of position & time. It is an observable quantity.

$\psi$  is a complex valued function.

We can only know the probable value in a measurement. The probability cannot be negative.

## Probability Interpretation of wave Function:-

$|\psi|^2$  (Square of magnitude of wave function) represent the probability of finding particle in that region.

Probability of finding the particle ~~is~~ is proportional to  $|\psi(x,y,z)|^2 dx dy dz$  at time  $t$ .

$$P \propto |\psi(x,y,z)|^2 dv$$

$|\psi|^2$  is called Probability density

$\psi$  is Probability Amplitude.

## Physical Interpretation of Wave Function $\psi$ :

A satisfactory interpretation of wave function  $\psi$  associated with a moving particle was given by Born in 1926.

He postulated that square of magnitude of wave function  $|\psi|^2$  at a particular point is proportional to probability of finding the particle at that point.

Also 
$$\int_{-\infty}^{\infty} |\psi|^2 d\tau = 1.$$

A wave function that obey this equation said to be normalized. Beside being normalized  $\psi$  must fulfill the following requirement.

i)  $\psi$  must be finite everywhere.

If  $\psi$  is ~~not~~ infinite it would have infinite probability of finding the particle at that point. This would violate the uncertainty principle. Hence  $\psi$  must be finite.

ii)  $\psi$  must be single valued.

If  $\psi$  has more than one value at any point, it would have more than one value of probability of finding the particle at that point. Hence it must be single valued.

iii)  $\psi$  must be continuous and have continuous first derivative everywhere.

$d^2\psi/dx^2$  must be finite everywhere. This possible is so only if  $d\psi/dx$  has no discontinuity. Further  $d\psi/dx$  as a continuous function implies that  $\psi$  too continuous across boundary.

## Heisenberg's Uncertainty Principle:-

(18)

This Principle was discovered by Heisenberg in 1927. It states that we can measure either position or momentum of a particle with any desired degree of accuracy (within limit of experimental equipment). OR

It is impossible to measure both, position and momentum simultaneously with accuracy.

$$\Delta x \cdot \Delta p_x \geq \frac{h}{4\pi}$$

$\Delta x \rightarrow$  Uncertainty in position.

$\Delta p \rightarrow$  Uncertainty in momentum.

### Physical Significance!

- I) Uncertainty Principle explain why it is possible for radiation and matter to have dual (wave-particle) character.
- II) It also make it clear that we can predict only the probable behaviour of quantum mechanical system not the exact behaviour.
- III) It helps in understanding many phenomenon like absence of electron within nuclei, natural broadening of spectral lines etc.

## Energy Momentum Uncertainty!

The uncertainty relation for simultaneous measurement of Energy  $E$  & time  $t$  expressed as.

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{2} \quad \left( \hbar = \frac{h}{2\pi} \right)$$

Derivation let us consider a microparticle moving with velocity  $v$ , Its kinetic Energy will be -

$$E = \frac{1}{2} m v^2$$

$$\Delta E = \Delta \left( \frac{1}{2} m v^2 \right)$$

$$= m v \Delta v = v \Delta p$$

$$\left( v^2 \frac{\Delta k}{\Delta t} \right)$$

$$\Rightarrow \Delta E = \frac{\Delta k}{\Delta t} \Delta p$$

$$\Delta E \cdot \Delta t = \Delta k \cdot \Delta p$$

$$\Delta E \cdot \Delta t \geq \frac{\hbar}{2}$$

$$\Delta k \cdot \Delta p \geq \frac{\hbar}{2}$$

**Example 1.5** Find the smallest possible uncertainty in position of the electron moving with velocity  $3 \times 10^7$  m/s (Given  $h = 6.63 \times 10^{-34}$  Js,  $m_0 = 9.1 \times 10^{-31}$  kg) [GGSIPU, May 2007 (2.5 marks)]

**Solution.** Given  $v = 3 \times 10^7$  m/s

Let  $\Delta x_{\min}$  be the minimum uncertainty in position of the electron and  $\Delta p$  the maximum uncertainty in the momentum of the electron.

Thus we have, 
$$\Delta x_{\min} \cdot \Delta p_{\max} = \frac{h}{2\pi} \quad \dots(i)$$

or

$$\Delta p_{\max} = p = mv$$

$$\Delta p_{\max} = \frac{m_0 v}{\sqrt{1 - \frac{v^2}{c^2}}} \quad \dots(ii)$$

$$\Delta x_{\min} = \frac{h \sqrt{1 - \frac{v^2}{c^2}}}{2\pi m_0 v} = \frac{6.63 \times 10^{-34} \sqrt{1 - \left(\frac{3 \times 10^7}{3 \times 10^8}\right)^2}}{2 \times 3.14 \times 9.1 \times 10^{-31} \times 3 \times 10^7} \text{ m}$$

$$= 0.03867 \times 0.9949 \times 10^{-10} \text{ m} = 3.8 \times 10^{-12} \text{ m}$$

**Example 1.7** Show that the uncertainty in the location of the particle is equal to de-Broglie wavelength the uncertainty in its velocity is equal to its velocity. [GGSIPU, April 2014 (2 marks)]

**Solution.** Given  $\Delta x = \lambda$ ,

Since we know that the uncertainty principle

$$\Delta x \Delta p_x = h \quad \text{or} \quad \lambda \Delta p_x = h$$

or 
$$\Delta p_x = \frac{h}{\lambda} \quad \text{or} \quad \Delta p_x = p$$

or 
$$m \Delta v_x = m v_x \quad \text{or} \quad \Delta v_x = v_x$$

$p$

- 1.23 A typical atomic nucleus is about 5 Fermi in radius. Use the uncertainty principle to place a lower limit on the energy an electron must have if it is part of nucleus. [GGSIPU, Feb. 2012 (2 marks)]

Hint : We know that 1 Fermi =  $10^{-15}$  m

The uncertainty principle in electron's position is

$$\Delta x = 5 \times 10^{-15} \text{ m}$$

$$\Delta p \geq \frac{\hbar}{\Delta x} = \frac{6.626 \times 10^{-34} \text{ Js}}{2\pi \times 5 \times 10^{-15} \text{ m}} = 2.11 \times 10^{-20} \text{ kg m/s}$$

The momentum would also be of the same order if this is the uncertainty in it. This suggests that the kinetic energy of electron is far greater than its rest energy and it can be written as

KE =  $pc$ , so that

$$pc \geq (2.11 \times 10^{-20} \text{ kg m/s}) \times (3 \times 10^8 \text{ m/s}) \geq 6.33 \times 10^{-12} \text{ J} \geq 39 \text{ MeV}$$

Thus, the kinetic energy of an electron must exceed 39 MeV, for it to be a nucleus constituent. Experiments indicate that the electrons in an atom have only a fraction of this energy. Thus, we can conclude that the electrons are present in the nucleus.

- 1.24 An electron has a speed of  $2 \times 10^4 \text{ ms}^{-1}$  within the accuracy 0.01%. Calculate the uncertainty in the position of the electron. [GGSIPU, May 2017 (2.5 marks)]

Hint :  $\Delta x \Delta p_x = \Delta x (m \Delta v_x) = \frac{h}{n\pi}$

$$\Rightarrow \Delta x = \frac{h}{n\pi m \Delta v_x} = \frac{6.63 \times 10^{-34}}{4\pi \times 9.1 \times 10^{-31} \times (2 \times 10^4 \times 0.0001)} = 29 \mu\text{m}$$

1.2 Calculate the de-Broglie wavelength of (i) a 46 g golf ball moving with velocity 30 m/s and (ii) an electron moving with velocity  $10^7$  m/s. Which one is measurable? [GGSIPIU, May 2010 (2.5 marks)]

Hint:  $\lambda = \frac{h}{mv}$

(i) For golf ball:  $\lambda = \frac{6.63 \times 10^{-34}}{46 \times 10^{-3} \times 30} = 4.8 \times 10^{-34}$  m [Not measurable]

(ii) For electron:  $\lambda = \frac{6.63 \times 10^{-34}}{9.1 \times 10^{-31} \times 10^7} = 72$  nm [measurable].

1.5 A ball of mass  $10^{-3}$  kg moves with a velocity of  $10^{-2}$  ms<sup>-1</sup>. What is the de-Broglie wavelength of the ball?

Hint:  $\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{6.626 \times 10^{-34} \text{ Js}}{10^{-3} \text{ kg} \times 10^{-2} \text{ ms}^{-1}} = 6.626 \times 10^{-29}$  m

1.6 An electron and a proton have the same de Broglie wavelength. Prove that the energy of electron is greater. [GGSIPIU, April 2015 (2 marks)]

Hint:  $\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{h}{\sqrt{2mE}}$  then  $\lambda_e = \lambda_p \Rightarrow \frac{h}{\sqrt{2m_e E_e}} = \frac{h}{\sqrt{2m_p E_p}}$

$m_e E_e = m_p E_p \Rightarrow E_e = \frac{m_p E_p}{m_e}$  since  $m_e < m_p$ , so  $E_e = \frac{m_p E_p}{m_e}$

1.7 An electron and a proton are moving with same velocity. Find the ratio of their (i) de Broglie wavelength, (ii) phase velocity and (iii) group velocity. [GGSIPIU, May 2015 (6 marks)]

Hint: (i)  $\frac{\lambda_e}{\lambda_p} = \frac{h}{m_e v_e} \times \frac{m_p v_p}{h} = \frac{m_p}{m_e}$

(ii)  $\frac{v_{se}}{v_{sp}} = \frac{v_e}{v_p} = 1$  and (iii)  $v_p v_g = c^2$  so  $\frac{v_{pe}}{v_{pp}} = \frac{v_e}{v_p} = 1$

1.8 Calculate the de Broglie wavelength of basket ball of mass 1 kg, moving at a speed of 10 ms<sup>-1</sup>. Discuss the reason, why we cannot observe its wave nature. [GGSIPIU, May 2015 (4.5 marks)]

Hint:  $\lambda = \frac{h}{mv} = \frac{6.63 \times 10^{-34}}{10 \times 10} \text{ m} = 6.63 \times 10^{-36}$  m which is not measurable.

1.13 Determine the de-Broglie wavelength of an electron having kinetic energy 2.0 eV

[Given : mass of electron =  $9.1 \times 10^{-31}$  kg,  $h = 6.63 \times 10^{-34}$  J.s]

[GGSIPIU, Feb. 2010 (2 marks)]

Hint:  $\lambda = \frac{hc}{\sqrt{K(K + 2m_0 c^2)}}$

$$= \frac{6.63 \times 10^{-34} \times 3 \times 10^8}{[2.0 \times 1.6 \times 10^{-19} (2.0 \times 1.6 \times 10^{-19} + 2 \times 9.1 \times 10^{-31} \times (3 \times 10^8)^2)]^{1/2}}$$

1.16 Find the phase and group velocities of an electron whose de-Broglie wavelength is  $1.2 \text{ \AA}$ .

[GGSIPU, Feb. 2012 (5 marks)]

Hint :  $m_0 c^2 = 511 \text{ keV}$ ,  $m_0 = 9.1 \times 10^{-31} \text{ kg}$

$$\frac{p^2}{2m} = \frac{h^2}{2m\lambda^2} \Rightarrow \frac{1}{2}mv^2 = \frac{h^2}{2m\lambda^2}$$

$$\Rightarrow v = \frac{h}{m\lambda} = 6.07 \times 10^6 \text{ m/s}$$

$$v = v_g = 6.07 \times 10^6 \text{ m/s}$$

But  $v_p v_g = c^2$

$$\Rightarrow v_p = \frac{c^2}{v_g} = \frac{(3 \times 10^8)^2}{6.07 \times 10^6} = 1.48 \times 10^{10} \text{ m/s.}$$

1.17 Calculate the de-Broglie wavelength of 40 keV electrons used in certain electron microscope.

[GGSIPU, May 2011 (2.5 marks)]

Hint :  $E = 40 \text{ keV} = 4.0 \times 10^3 \times 1.6 \times 10^{-16} \text{ J} = 6.4 \times 10^{-15} \text{ J}$

$$E = \frac{1}{2}mv^2 \quad \text{or} \quad v = \sqrt{2E/m}$$

$$\text{de-Broglie wavelength, } \lambda = \frac{h}{mv} = \frac{h}{\sqrt{2mE}} = 6.15 \times 10^{-15} \text{ m}$$

1.18 Calculate the de-Broglie wavelength of an electron accelerated through a potential difference 100 V.

[GGSIPU, April 2011 (2 marks)]

$$\text{Hint : } \lambda = \frac{h}{mv} = \frac{h}{\sqrt{2mE}} = \frac{1.228 \times 10^{-10}}{\sqrt{V}} = 1.228 \times 10^{-10} \text{ m} = 1.228 \text{ \AA}$$

1.19 A nuclear particle is confined to a nucleus of diameter  $5 \times 10^{-14} \text{ m}$ . Calculate the minimum uncertainty in the momentum of the nucleon. Also calculate the minimum kinetic energy of the nucleon.

[GGSIPU, Feb. 2009 (2 marks)]

Hint : The diameter of nucleus  $(\Delta x) = 5 \times 10^{-14} \text{ m}$

$$\therefore \Delta p \Delta x = \hbar$$

$$\Rightarrow \Delta p = \frac{\hbar}{\Delta x} = \frac{1.05 \times 10^{-34}}{5 \times 10^{-14}} \text{ kg m s}^{-1} = 2.1 \times 10^{-21} \text{ kg m s}^{-1}$$

$$\text{and the minimum K.E. of the nucleon} = \frac{p^2}{2m_0} = \frac{2.1 \times 10^{-21}}{2 \times m_0} \text{ here } m_0 = \text{mass of nucleon.}$$

# Group velocity & Phase velocity:-

Phase velocity:- Velocity of individual wave forming wave packet is called phase velocity.

According to de Broglie Hypothesis



If a particle is in motion a wave is associated with it. called matter wave

wavelength of wave  $\lambda = \frac{h}{mv}$  (1)

(where  $v$  is the velocity of particle)

$h$  - Planck's constant

$m$  - mass of the particle)

We have to calculate the velocity of wave

velocity of wave  ~~$v$~~

$v = v\lambda$  (2) ( $v$  - frequency)

Note

$E = hv \Rightarrow v = \frac{E}{h}$

But  $E = mc^2$

$v = \frac{mc^2}{h}$  (3)

Note Extra

$\Rightarrow$  using (1), (3) in (2)

$v = v\lambda = \frac{mc^2}{h} \times \frac{h}{mv} = \frac{c^2}{v} \Rightarrow v = \frac{c^2}{v}$  (4)

velocity of wave

A plane wave travelling along x-axis is given by -

$$\psi = A \sin(\omega t - kx)$$

$$u = v\lambda$$

(where  $\omega$  is Angular velocity  
 $k$  - Propagation constant)

$$\left( \begin{array}{l} \omega = 2\pi\nu \\ k = \frac{2\pi}{\lambda} \end{array} \right)$$

Now  $\psi = A \sin(\omega t - kx)$   $u = \frac{2\pi \nu \lambda}{2\pi}$

$$= \frac{\omega \lambda}{2\pi} = \frac{\omega}{k}$$

$$u = \frac{\omega}{k}$$

$$\boxed{v_p = \frac{\omega}{k}} \text{ --- phase velocity.}$$

# Group Velocity:-

~~Group velocity~~ is The velocity of wave packet formed by the individual waves is called group velocity.

Superposition of wave is called wave packet or wave group.



Let us consider a wave group which consist of two component of equal amplitude  $a$ , but slightly different angular frequencies  $\omega_1, \omega_2$  and propagation constant  $k_1, k_2$ . Displacement is given by.

$$y_1 = a \sin(\omega_1 t - k_1 x)$$

$$y_2 = a \sin(\omega_2 t - k_2 x)$$

Then Superposition gives -

$$y = y_1 + y_2$$

$$= a [\sin(\omega_1 t - k_1 x) + \sin(\omega_2 t - k_2 x)]$$

Using  $\left( \sin A + \sin B = 2 \sin \frac{A+B}{2} \cos \frac{A-B}{2} \right)$  we get -

$$y = 2a \sin \left( \frac{(\omega_1 t - k_1 x) + (\omega_2 t - k_2 x)}{2} \right)$$

$$\cos \left( \frac{\omega_1 t - k_1 x - \omega_2 t + k_2 x}{2} \right)$$

$$= 2a \sin \left( \frac{(\omega_1 + \omega_2)t}{2} - \frac{(k_1 + k_2)x}{2} \right)$$

$$\cos \left( \frac{(\omega_1 - \omega_2)t}{2} - \frac{(k_1 - k_2)x}{2} \right)$$

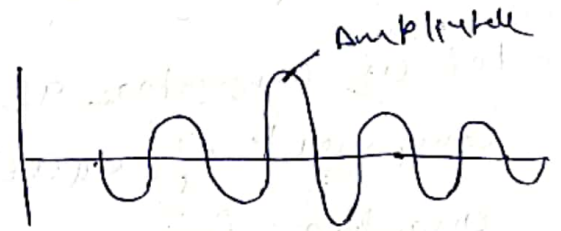
(16)

$$y = 2a \cos \left( \frac{(\omega_1 - \omega_2)t}{2} - \frac{(k_1 - k_2)x}{2} \right) \sin \left[ \frac{(\omega_1 + \omega_2)t}{2} - \frac{(k_1 + k_2)x}{2} \right]$$

Amplitude Angle

Note (  $\phi = \text{ASIN}(\omega t - kx)$  )

For group velocity we use Amplitude velocity



where Amplitude A

$$A = 2a \cos \left( \frac{(\omega_1 - \omega_2)t}{2} - \frac{(k_1 - k_2)x}{2} \right)$$

∴ Group velocity  $v_g = \frac{\omega_1 - \omega_2}{\frac{k_1 - k_2}{2}} = \frac{\omega_1 - \omega_2}{k_1 - k_2}$

$$\boxed{v_g = \frac{\Delta \omega}{\Delta k}} \leftarrow \text{Group velocity}$$

Relation between group velocity & Phase velocity!

We have  $v_p = \frac{\omega}{k} \Rightarrow \omega = (k v_p)$

$$v_g = \frac{d\omega}{dk} = \frac{d(k v_p)}{dk} = v_p + k \frac{dv_p}{dk}$$

(Using  $\boxed{k = \frac{2\pi}{\lambda}}$ )  $v_g = v_p + \frac{2\pi}{\lambda} \frac{dv_p}{d(\frac{2\pi}{\lambda})}$

$$v_g = v_p + \frac{1}{\lambda} \frac{dv_p}{d\left(\frac{1}{\lambda}\right)}$$

(17)

But  $d\left(\frac{1}{\lambda}\right) = -\frac{1}{\lambda^2} d\lambda$ , therefore

$$v_g = v_p - \lambda \frac{dv_p}{d\lambda}$$

1.20 An electron has a de Broglie wavelength 2 pm. Find its kinetic energy, phase velocity and group velocity of its de Broglie wave. Rest mass energy of electron is 511 keV.

[GGSIPU, Feb. 2017 (3 marks)]

Hint:  $m_0 c^2 = 511 \text{ keV} \Rightarrow m_0 = 9.1 \times 10^{-31} \text{ kg}$

$$\text{K.E.} = \frac{p^2}{2m} = \frac{h^2}{2m\lambda^2} = \frac{1}{2} m v_g^2 \Rightarrow v_g = ? \Rightarrow v_p v_g = c^2 = ?$$

1.1 Find the phase velocity and group velocity of the de-Broglie wave of an electron whose speed is  $0.9c$ .

[GGSIPU, Feb. 2009 (2 marks), May 2009 (2.5 marks); May 2019 (2.5 marks)]

Hint:  $v_p v_g = c^2$  and  $v_g = 0.9c$ . Then  $v_p = 3.33 \times 10^8 \text{ m/s}$ .

Expectation values:- (Average value)

(25)

The average value of a large number of measurements of physical quantity is represented by ~~large~~ Expectation value of a function  $f(r)$ .

$$\langle f(r) \rangle = \text{Average } f(r) = \overline{f(r)}$$

$$\langle f(r) \rangle = \int \psi^*(r,t), f(r) \psi(r,t) d\tau$$

Average value of position vector -

$$\langle \vec{r} \rangle = \int \psi^*(\vec{r},t) \vec{r} \psi(\vec{r},t) d\tau$$

Average value of momentum.

-(1)

(Expectation)

$$\langle P \rangle = \int \psi^* P \psi d\tau$$

$$\hat{P} = -i\hbar \nabla$$

$$\langle P \rangle = \int \psi^* (-i\hbar \nabla \psi) d\tau$$

Expectation value of potential Energy  $V$  can be written as -

-(2)

$$\langle V \rangle = \int \psi^*(r,t) V(r) \psi(r,t) d\tau$$

-(3)

Eqn<sup>n</sup> (1), (2), (3) are valid for Normalized wave function. If wave function are not Normalized then average value function can be written as -

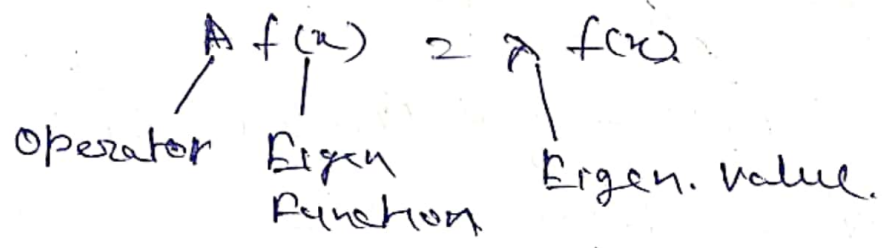
$$\langle f(r) \rangle = \frac{\int \psi^* f(r) \psi d\tau}{\int \psi^* \psi d\tau}$$

# Eigen Values & Eigen Function

If there is some function  $f(x)$  which when operated on operator  $A$  gives the Equation

$$\boxed{\hat{A}f(x) = \lambda f(x)}$$

$\lambda$  is constant and various possible value of  $\lambda$  are called eigen values of the operator and these functions are called Eigen Function.



Note  
Eigen values for matrices  
 $|A - \lambda I| = 0$

Suppose  $\psi$  be any wave function



Then  $A\psi$  (operator of  $\psi$ ) =



where  $A$  is operator.

(Note: An operator is a mathematical operation on wave function that produce another wave function)

Thus "Eigen value Equation is that Equation in which operator operates on a function give the function multiplied by some constant"

$$\boxed{A f(x) = \lambda f(x)}$$

**Example 2.2** A particle limited to the  $x$ -axis has the wave function  $\psi = ax$  between  $x = 0$  and  $x = 1$ ,  $\psi = 0$  elsewhere. Find (a) the probability that particle can found between  $x = 0.45$  and  $x = 0.55$ . (b) The expectation value  $\langle x \rangle$  of the particle's position. [GGSIPU, May 2014 (3 marks), Feb. 2008, April 2007, Feb. 2012 (2 marks)]

**Solution.** (a) The probability is

$$\int_{x_1}^{x_2} |\psi|^2 dx = a^2 \int_{0.45}^{0.55} x^2 dx = a^2 \left[ \frac{x^3}{3} \right]_{0.45}^{0.55} = 0.0251 a^2$$

(b) The expectation value of the particle's position is

$$\langle x \rangle = \int_0^1 x |\psi|^2 dx = a^2 \int_0^1 x^3 dx = a^2 \left[ \frac{x^4}{4} \right]_0^1 = \frac{a^2}{4}$$

**2.13** The eigenfunctions for a particle in a 1-D box are given by  $\psi(x) = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$ . Find the expectation values of (i) position and (ii) momentum in the  $n$ th quantum state.

[GGSIPU, May 2016 (4 marks)]

**Hint:** (i)  $\langle x \rangle = \int_{-\infty}^{\infty} \psi_n^* x \psi_n dx = \frac{2}{a} \int_0^a x \sin^2 \frac{n\pi x}{a} dx$

Integrating by parts, we get:  $\langle x \rangle = \frac{a}{2}$ .

This result is as expected; the probability density  $\psi^* \psi$  is symmetric about  $\langle x \rangle = \frac{a}{2}$ , indicating that the particle spends as much time to the left of the center as to the right.

$$\begin{aligned} \text{(ii)} \quad \langle p_x \rangle &= \int_{-\infty}^{\infty} \psi_n^* \left( i\hbar \frac{d}{dx} \right) \psi_n dx = \frac{2i\hbar}{a} \int_0^a \sin \frac{d}{dx} \left( \sin \frac{n\pi x}{a} \right) dx \\ &= -\frac{2i\hbar n\pi}{a} \int_0^a \sin \frac{n\pi x}{a} \cos \frac{n\pi x}{a} dx = -\frac{i\hbar n\pi}{a^2} \int_0^a \sin \frac{2n\pi x}{a} dx \end{aligned}$$

Solving the integral, we get  $\langle p_x \rangle = 0$ .

Again the result is expected. The particle moves back and forth, spending half its time moving towards the left and half its time moving towards the right. Thus, the average momentum must be zero.

## Schrodinger's Time dependent Equation in one dimension

Schrodinger presented the wave Equation describing the properties of electron in the form of equation called Schrodinger wave Equation to calculate the probabilities of electron at different places around nucleus in the atom OR.

It is a differential Equation for de-Broglie waves associated with particle and describing motion of particle.

$$i\hbar \frac{d\psi}{dt} = -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V\psi$$

$m$  - Mass of electron

$V$  - Potential Energy with reference to nucleus

$\psi$  - wave function.

Derivation:- Equation of a particle moving freely in positive  $x$  direction is.

$$\psi = A e^{-i(\omega t - kx)} \quad \text{--- (1)}$$

Note

$\omega = 2\pi\nu$ , $E = h\nu$	$k = \frac{2\pi}{\lambda}$ , $\lambda = \frac{h}{p}$
$\omega = \frac{2\pi E}{h}$ , $\nu = E/h$	

$k = \frac{2\pi p}{h}$

where  $A$  - Amplitude,  $k$  - Propagation constant  
 $\omega$  - Constant Angular Frequency.

Now substituting value of  $\omega$  &  $k$  in Equation (1)

$$\psi = A e^{-i(\omega t - kx)}$$

$$\psi = A e^{-i\left(\frac{2\pi E t}{h} - \frac{2\pi p x}{h}\right)}$$

$$\psi = A e^{-i \left( \frac{2\pi}{h} (Et - Px) \right)}$$

Using  $\hbar = \frac{h}{2\pi}$  here  $\hbar$  is modified form of Planck's constant.

$$\Rightarrow \psi = A e^{-\frac{i}{\hbar} (Et - Px)}$$

$$= A e^{\frac{i}{\hbar} (Px - Et)} \quad \text{--- (2)}$$

Now Differentiating (2) wrt  $x$

$$\frac{d\psi}{dx} = A e^{\frac{i}{\hbar} (Px - Et)} \frac{d}{dx} \left( \frac{i}{\hbar} (Px - Et) \right)$$

Note  
using  
 $\left( \frac{d}{dx} e^y = e^y \frac{dy}{dx} \right)$

$$= \psi \frac{iP}{\hbar} \quad \text{from (2)}$$

$$\Rightarrow P\psi = \frac{\hbar}{i} \frac{d\psi}{dx} \quad \text{--- (3)}$$

Now Differentiating wrt  $t$

$$\frac{d\psi}{dt} = A \left( -\frac{iE}{\hbar} \right) e^{\frac{i}{\hbar} (Px - Et)}$$

$$= -\frac{iE}{\hbar} \psi \quad \text{from (2)}$$

$$\Rightarrow E\psi = -\frac{\hbar}{i} \frac{d\psi}{dt} = \hbar i \frac{d\psi}{dt} \quad \text{--- (4)}$$

Now Total Energy

$E = \text{Kinetic Energy} + \text{Potential Energy}$

$$E = \frac{1}{2} m v^2 + V$$

$$\boxed{E = \frac{p^2}{2m} + V}$$

$$E(\psi) = \frac{p^2}{2m} (\psi) + V(\psi) \quad \text{--- (5)}$$

Now Substituting value of (3) (4) in (5)

we get

$$\hbar i \frac{d\psi}{dt} = \left( \frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} \right) + V(\psi).$$

$$\Rightarrow \boxed{i\hbar \frac{d\psi}{dt} = \left( -\frac{\hbar^2}{2m} \frac{d^2\psi}{dx^2} + V(\psi) \right)} \quad \text{--- (6)}$$

Time dependent Schrodinger wave Equation.

Time independent ~~Schrodinger~~ Schrodinger Equation

We have Equation of a Particle moving freely in positive  $x$  direction.

$$\psi = A e^{-i(\omega t - kx)}$$

$$\left( \text{Using } \omega = \frac{2\pi E}{h} \quad \& \quad k = \frac{2\pi p}{h} \right)$$

$$\psi(x,t) = A e^{-i \frac{2\pi}{h} (Et - px)}$$

$$= A e^{i \frac{2\pi}{h} (px - Et)}$$

$$= A \left( e^{\frac{2\pi i px}{h}} \right) \left( e^{-\frac{2\pi i Et}{h}} \right)$$

Again we have. --- (1)

$$\psi = A e^{-i(\omega t - kx)}$$

For time independent  $t=0$

$$\psi(x) = A e^{-i(-kx)}$$

$$= A e^{-i \left( -\frac{2\pi p}{h} x \right)}$$

$$\psi(x) = A e^{\frac{2\pi i px}{h}}$$

--- (2)

Substituting (2) in (1) we get.

$$\Psi(x,t) = \Psi(x) e^{-2\pi i E t / h} \quad \text{--- (A)}$$

Differentiating wrt t

$$\frac{d\Psi}{dt} = \Psi(x) e^{-\frac{2\pi i E t}{h}} \left( -\frac{2\pi i E}{h} \right) \quad \text{--- (B)}$$

Differentiating wrt x

$$\frac{d\Psi}{dx} = e^{-\frac{2\pi i E t}{h}} \frac{d\Psi}{dx}$$

$$\Rightarrow \frac{d^2\Psi}{dx^2} = e^{-\frac{2\pi i E t}{h}} \frac{d^2\Psi}{dx^2} \quad \text{--- (C)}$$

Substituting value of A, B, C in (5)

From (5) we have

$$-\frac{\hbar^2}{2m} \frac{d^2\Psi}{dx^2} + V\Psi = i\hbar \frac{d\Psi}{dt}$$

$$\Rightarrow -\frac{\hbar^2}{2\pi^2 m} \frac{d^2\Psi}{dx^2} + V\Psi = \frac{i\hbar}{2\pi} \frac{d\Psi}{dt}$$

$$\Rightarrow \frac{-\hbar^2}{8\pi^2 m} \left( e^{-\frac{2\pi i E t}{h}} \frac{d^2\Psi}{dx^2} \right) + V \left( \Psi(x) e^{-\frac{2\pi i E t}{h}} \right) = \frac{i\hbar}{2\pi} \Psi(x) e^{-\frac{2\pi i E t}{h}} \left( -\frac{2\pi i E}{h} \right)$$

$$\Rightarrow \frac{e^{-\frac{2\pi i E t}{h}}}{h} \left( -\frac{\hbar^2}{8\pi^2 m} \frac{d^2\Psi}{dx^2} + V\Psi \right) = \frac{e^{-\frac{2\pi i E t}{h}}}{h} E\Psi$$

$$\Rightarrow \boxed{\frac{d^2\Psi}{dx^2} + \frac{8\pi^2 m}{\hbar^2} (E - V)\Psi = 0}$$

Time independent Schrodinger Eqn

## The Free Particle!-

- The particle is said to be a free particle when it is moving in space without being subjected to any external force in any region of space and its potential energy is constant ( $V = \text{constant}$ ).

Then time-independent Schrodinger wave Equation for free particle is given as -

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2}(E - V)\psi = 0$$

$$V=0$$

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2}E\psi = 0$$

Note

$\text{Kinetic Energy} = \frac{1}{2}mv^2$ $= \frac{1}{2} \frac{m}{m} mv^2$ $E = \frac{p^2}{2m}$ $p = \sqrt{2mE}$	$\text{Propagation constant}$ $k = \frac{2\pi}{\lambda}$ $= \frac{2\pi p}{h}$ $k = \frac{2\pi\sqrt{2mE}}{h}$ <p>on squaring</p> $k^2 = \frac{8\pi^2mE}{h^2}$
--	--

Hence  $\frac{d^2\psi}{dx^2} + k^2\psi = 0$  where  $k^2 = \frac{8\pi^2mE}{h^2}$  (1)

a) Wave Function! → General solution of Equation (1)

is given by

$$\psi(x) = Ae^{ikx} + Be^{-ikx}$$

b) Energy! → Energy of particle is given by -

we have  $k^2 = \frac{8\pi^2mE}{h^2}$

(Using  $h = \frac{h}{2\pi}$ )

$$E = \frac{h^2 k^2}{8\pi^2 m} = \frac{\hbar^2 k^2}{2m}$$

c) Momentum!  $\psi = A e^{-i(\omega t - kx)}$

$$\left( \begin{array}{l} \omega = 2\pi\nu \\ \quad = \frac{2\pi E}{h} \end{array} \quad \begin{array}{l} k = \frac{2\pi}{\lambda} \\ \quad = \frac{2\pi p}{h} \end{array} \right)$$

$$\psi = A e^{-i \left( \frac{2\pi}{h} (Et - px) \right)}$$

using  
 $\left( \hbar = \frac{h}{2\pi} \right)$

$$= A e^{\frac{i}{\hbar} (px - Et)}$$

$$= A e^{\frac{i}{\hbar} (px - Et)}$$

$$\frac{d\psi}{dx} = A \frac{i p}{\hbar} e^{\frac{i}{\hbar} (px - Et)}$$

$$= \frac{i p}{\hbar} \psi$$

$$\Rightarrow \boxed{p\psi = \hbar \frac{d\psi}{dx}}$$

d) Position of particle  $\Rightarrow$  Probability of finding the particle is given by -

$$P dx = \psi^*(x,t) \psi(x,t)$$

The Probability density  $P$  for the position of particle with definite value of momentum is constant over  $x$ -axis.

# Particle in a Rigid one Dimensional box (Infinite Square well potential):

Let us consider a particle restricted to move along x-axis between  $x=0$  and  $x=L$ .

Suppose potential Energy  $V$  of the particle is zero inside the box, but rises to infinity outside.

that is  $V=0$  for  $(0 \leq x \leq L)$  - inside box  
 $V=\infty$  for  $(x < 0 + x > L)$  - outside the box

In such case the particle is said to be moving in an infinitely deep square well potential

Schrodinger Equation for the particle within box is -

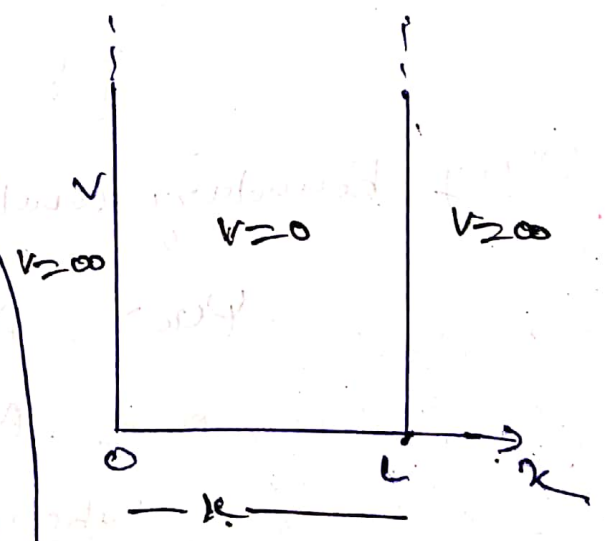
$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2mE}{h^2} \psi = 0 \quad \text{--- (1)}$$

$V=0$  (inside the box).

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m(E-V)}{h^2} \psi = 0$$

Substitute  $V=0$

$$\Rightarrow \frac{d^2\psi}{dx^2} + \frac{8\pi^2mE}{h^2} \psi = 0$$



Suppose  $\frac{8\pi^2mE}{h^2} = k^2$  --- (2)

$$\frac{d^2\psi}{dx^2} + k^2\psi = 0 \quad \text{--- (3)}$$

General Solution of differential Equation of the form -

$$\psi(x) = A \sin kx + B \cos kx \quad \text{--- (4)}$$

where the constant A & B are to be determined by boundary conditions. Therefore, the wavefunction  $\psi$  must be zero outside the box.

$\psi$  must also be zero at the walls ( $x=0, x=L$  at walls)

Using the boundary condition  $\psi=0$  at  $x=0$  in (4)

$$\psi(x) = A \sin kx + B \cos kx$$

$$0 = A \sin k(0) + B \cos k(0)$$

$$0 = 0 + B(1)$$

$$\boxed{B=0} \Rightarrow \boxed{\psi(x) = A \sin kx} \quad \text{--- (5)}$$

Using boundary condition  $\psi(x)=0$  at  $x=L$  in (4)

$$\psi(x) = A \sin kx + B \cos kx$$

$$0 = A \sin k(L) + B \cos k(L)$$

Substituting  $B=0$

$$0 = A \sin kL + 0$$

$$A \sin kL = 0$$

$$\sin kL = 0$$

$$\sin kL = 0 \quad (n = 1, 2, 3, \dots)$$

$$\sin kL = \sin n\pi$$

$$kL = n\pi$$

$$\boxed{k = \frac{n\pi}{L}}$$

Substituting  $k = \frac{n\pi}{L}$  in (5)

(28)

$$\Rightarrow \boxed{\psi(x) = A \sin\left(\frac{n\pi x}{L}\right)} \quad (6)$$

Now for Energy we have

$$k^2 = \frac{8\pi^2 m E}{h^2} \quad (\text{by (2)})$$

$$\Rightarrow E = \frac{k^2 h^2}{8\pi^2 m}$$

Substituting  $k = \frac{n\pi}{L}$

$$\boxed{E = \frac{n^2 h^2}{8mL^2}} \quad (n = 1, 2, 3, \dots)$$

Thus we see that in a potential well the particle can not have an arbitrary energy but can have only certain energy values corresponding to  $n = 1, 2, 3, \dots$  these are Eigen values of the particle in the well and correspond to energy level of the system.

$$E_n = \frac{n^2 h^2}{8mL^2}$$

Ground state Energy ( $n=1$ )

$$E_1 = \frac{h^2}{8mL^2}$$

First Excited State Energy ( $n=2$ )

$$E_2 = \frac{4h^2}{8mL^2}$$

We have calculated the values of Energy (24)

Now we will find the solution of wave function in Equation (6) and value of A

From (6) we have

$$\psi(x) = A \sin\left(\frac{n\pi x}{L}\right)$$

where A  $\rightarrow$  Normalization factor or Normalization constant

To find the value of A we will use Normalization condition.

$$\int_0^L \psi \psi^* dx = 1$$

Normalization condition

$$\Rightarrow \int_0^L \left( A \sin\left(\frac{n\pi x}{L}\right) \right) \left( A \sin\left(\frac{n\pi x}{L}\right) \right) dx = 1$$

$$\Rightarrow A^2 \int_0^L \sin^2\left(\frac{n\pi x}{L}\right) dx = 1$$

$$\Rightarrow A^2 \int_0^L \left( \frac{1 - \cos\left(\frac{2n\pi x}{L}\right)}{2} \right) dx = 1$$

$$\left( \begin{array}{l} \text{using } \cos 2\theta = 1 - 2\sin^2\theta \\ 2\sin^2\theta = 1 - \cos 2\theta \\ \sin^2\theta = \frac{1 - \cos 2\theta}{2} \end{array} \right)$$

$$\Rightarrow \frac{A^2}{2} \int_0^L \left( 1 - \cos\left(\frac{2n\pi x}{L}\right) \right) dx = 1$$

$$\Rightarrow \frac{A^2}{2} \left[ (x)_0^L - \left[ \sin\left(\frac{2n\pi x}{L}\right) \right]_0^L \cdot \frac{L}{2n\pi} \right] = 1$$

$$\Rightarrow \frac{A^2}{2} [L - 0] = 1$$

$$A^2 = \frac{2}{L}$$

$$A = \sqrt{\frac{2}{L}}$$

in (6)

$$\psi(x) = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$$

Solution of wave function

Plotting of wave function:-

First three Eigen function  $\psi_1, \psi_2, \psi_3$  together with the Probability densities  $|\psi_1|^2, |\psi_2|^2, |\psi_3|^2$  are shown in Figure

we have  $\psi_n = \sqrt{\frac{2}{L}} \sin\left(\frac{n\pi x}{L}\right)$

1)  $n=1$  for ground state

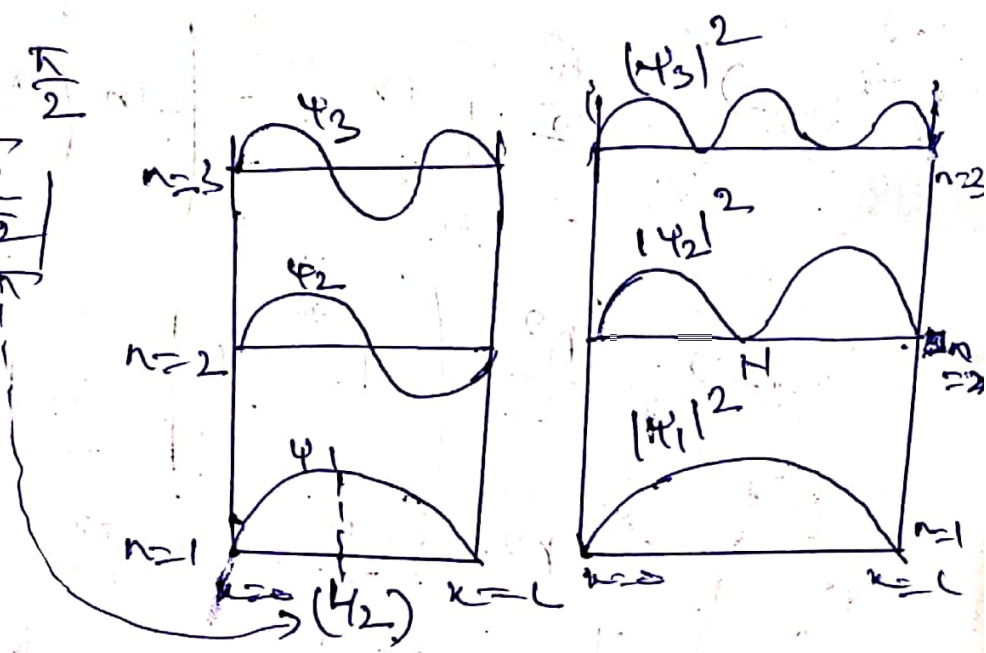
$$\psi_1 = \sqrt{\frac{2}{L}} \sin\left(\frac{\pi x}{L}\right)$$

Maximum value of  $\psi = \sin 90^\circ = \sin\left(\frac{\pi}{2}\right)$

$\Rightarrow \sin\left(\frac{n\pi x}{L}\right) = \sin\left(\frac{\pi}{2}\right) = 1$

$\frac{n\pi x}{L} = \frac{\pi}{2}$

$x = \frac{L}{2}$



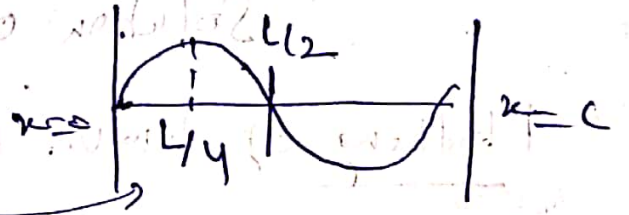
ii)  $n=2$

$\psi_2 = \sqrt{\frac{2}{L}} \sin\left(\frac{2\pi x}{L}\right)$

$\sin\left(\frac{2\pi x}{L}\right) = \sin\frac{\pi}{2}$

$\frac{2\pi x}{L} = \frac{\pi}{2}$

$x = \frac{L}{4}$



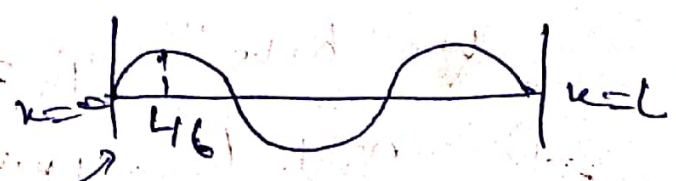
iii)  $n=3$

$\psi_3 = \sqrt{\frac{2}{L}} \sin\left(\frac{3\pi x}{L}\right)$

$\sin\left(\frac{3\pi x}{L}\right) = \sin\frac{\pi}{2}$

$\frac{3\pi x}{L} = \frac{\pi}{2}$

$x = \frac{L}{6}$



- 2.9 Consider a particle confined in one-dimensional box of width  $l$ . Find the probability that the particle is found between  $x = 0$  and  $x = l/n$  when it is in  $n$ th state. [IGGSIPU, April 2015 (2 marks)]

Hint :  $\psi_n(x) = \sqrt{\frac{2}{l}} \sin \frac{n\pi x}{l}$

$$p = \int_{x_1}^{x_2} |\psi_n(x)|^2 dx = \frac{2}{l} \int_{x_1}^{x_2} \sin^2 \frac{n\pi x}{l} dx$$

$$= \left[ \frac{x}{l} - \frac{1}{2n\pi} \sin \frac{2n\pi x}{l} \right]_{x_1=0}^{x_2=l/n} = \left[ \frac{1}{n} - \frac{1}{2n\pi} \sin \frac{2n\pi l}{nl} \right] = \frac{1}{n}$$

- 2.10 An electron is constrained to move in a one dimensional box of length 0.1 nm. Find the first three energy eigen values and the corresponding de Broglie wavelengths. [IGGSIPU, May 2015 (4 marks)]

Hint :  $E_n = \frac{n^2 h^2}{8ml^2} \Rightarrow E_n = \frac{n^2 \times (6.623 \times 10^{-34})^2}{8 \times (9.1 \times 10^{-31}) \times (0.1 \times 10^{-9})^2}$

$$E_n = \frac{n^2 \times (6.623 \times 10^{-34})^2}{8 \times (9.1 \times 10^{-31}) \times (0.1 \times 10^{-9})^2 \times (1.6 \times 10^{-19})} \text{ eV} = 37.5 n^2 \text{ eV}$$

Then  $E_1 = 37.5 \text{ eV}$ ,  $E_2 = 150 \text{ eV}$ ,  $E_3 = 337.5 \text{ eV}$

For de Broglie wavelength  $\lambda_n = \frac{h}{\sqrt{2mE_n}}$

- 2.11 Find the probability that a particle trapped in a box ' $L$ ' wide can be found between  $0.45L$  and  $0.55L$  for the first excited state. [IGGSIPU, May 2015 (4 marks)]

Hint :  $\psi_n = \sqrt{\frac{2}{l}} \sin \frac{n\pi x}{l}$  ;  $P = \int_{x_1}^{x_2} |\psi_n(x)|^2 dx = \frac{2}{l} \int_{x_1}^{x_2} \sin^2 \frac{2n\pi x}{l} dx$

$$= \left[ \frac{x}{l} - \frac{1}{2n\pi} \sin \frac{2n\pi x}{l} \right]_{x_1}^{x_2}$$

Now  $x_1 = 0.45L$  and  $x_2 = 0.55L$  and  $n = 2$  for first excited state, on solving it we get  $P = 9.8\%$ .

- 2.12 The wave function for a particle in a 1-D box is given by  $\psi(x) = \sqrt{\frac{2}{a}} \sin \frac{m\pi x}{a}$ . Show that the wave function for two different states are orthonormal. [IGGSIPU, May 2016 (3 marks)]

Hint : For the eigenfunctions in the  $m$ th and  $n$ th states,  $m \neq n$ , to be orthogonal, we must have

$$\int_{-\infty}^{\infty} \psi_m^* \psi_n dx = 0$$

For the given wavefunction,  $\psi_m(x) = \sqrt{\frac{2}{a}} \sin \frac{m\pi x}{a}$ , then  $\psi_n(x) = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$

$$\int_{-\infty}^{\infty} \psi_m^* \psi_n dx = \frac{2}{a} \int_0^a \sin \frac{m\pi x}{a} \sin \frac{n\pi x}{a} dx$$

Therefore  $= \frac{2}{a} \left[ \int_0^a \cos \frac{(m-n)\pi x}{a} dx - \int_0^a \cos \frac{(m+n)\pi x}{a} dx \right]$

2.8 An electron is in a box of 0.01 nm. Find its permitted energy.

[GGSIU, Feb. 2013, (2 marks) ; May 2008 (2.5 marks)]

Hint :

$$E_n = \frac{n^2 h^2}{8ml^2} = \frac{n^2 \times (6.626 \times 10^{-34})^2}{8 \times 9.1 \times 10^{-31} (0.01 \times 10^{-9})^2}$$

$$= 6 \times 10^{-16} n^2 \text{ J} = \frac{6 \times 10^{-16}}{1.6 \times 10^{-19}} \text{ eV} = \frac{60000}{16} n^2 \text{ eV}$$

For  $n = 1, 2, 3, \dots$  the energy eigen value  $E_1, E_2, E_3, \dots$

$$E_1 = 3750 \text{ eV}, E_2 = 15000 \text{ eV}, E_3 = 33750 \text{ eV}, \dots$$

2.13 The eigenfunctions for a particle in a 1-D box are given by  $\psi(x) = \sqrt{\frac{2}{a}} \sin \frac{n\pi x}{a}$ . Find the expectation values of (i) position and (ii) momentum in the  $n$ th quantum state.

[GGSIU, May 2016 (4 marks)]

Hint : (i)  $\langle x \rangle = \int_{-\infty}^{\infty} \psi_n^* x \psi_n dx = \frac{2}{a} \int_0^a x \sin^2 \frac{n\pi x}{a} dx$

Integrating by parts, we get :  $\langle x \rangle = \frac{a}{2}$ .

This result is as expected ; the probability density  $\psi^* \psi$  is symmetric about  $\langle x \rangle = \frac{a}{2}$ , indicating that the particle spends as much time to the left of the center as to the right.

(ii)  $\langle p_x \rangle = \int_{-\infty}^{\infty} \psi_n^* \left( i\hbar \frac{d}{dx} \right) \psi_n dx = \frac{2i\hbar}{a} \int_0^a \sin \frac{d}{dx} \left( \sin \frac{n\pi x}{a} \right) dx$

$$= -\frac{2i\hbar}{a} \frac{n\pi}{a} \int_0^a \sin \frac{n\pi x}{a} \cos \frac{n\pi x}{a} dx = -\frac{i\hbar n\pi}{a^2} \int_0^a \sin \frac{2n\pi x}{a} dx$$

Solving the integral, we get  $\langle p_x \rangle = 0$ .

Again the result is expected. The particle moves back and forth, spending half its time moving towards the left and half its time moving towards the right. Thus, the average momentum must be zero.

- 2.4 For an electron in a one-dimensional box of width  $2 \text{ \AA}$ , calculate the separation between the lowest two levels in eV. [GGSIPU, June 2013 (4.5 marks)]

Hint :  $E_n = \frac{n^2 h^2}{8ml^2}$

$$\Rightarrow E_1 = \frac{1 \times (6.63 \times 10^{-34})^2}{8 \times 9.1 \times 10^{-31} \times (2 \times 10^{-10})^2} = 2.273 \times 10^{-19} \text{ J} = 1.42 \text{ eV}$$

- 2.17 An electron is confined to move between two rigid walls separate by  $2 \times 10^{-9} \text{ m}$ . Find the de Broglie wavelengths representing the first three allowed energy states of the electron. [GGSIPU, May 2017 (4 marks)]

Hint :  $E_n = \frac{n^2 h^2}{8ml^2} \Rightarrow \frac{hc}{\lambda_n} = \frac{n^2 h^2}{8ml^2} \Rightarrow \lambda_n = \frac{c}{n^2 h} \times 8ml^2$

$$\lambda_1 = \frac{3 \times 10^8}{(6.62 \times 10^{-34})} \times 8 \times (9.1 \times 10^{-31}) \times (2 \times 10^{-9})^2 = 1.32 \times 10^{-5} \text{ m}$$

$$\lambda_2 = \frac{3 \times 10^8}{4 \times (6.62 \times 10^{-34})} \times 8 \times (9.1 \times 10^{-31}) \times (2 \times 10^{-9})^2 = 3.3 \times 10^{-6} \text{ m}$$

$$\lambda_3 = \frac{3 \times 10^8}{9 \times (6.62 \times 10^{-34})} \times 8 \times (9.1 \times 10^{-31}) \times (2 \times 10^{-9})^2 = 1.47 \times 10^{-6} \text{ m}$$

**Problem 2.8** The wave function of a certain particle is  $\psi = A \cos^2 x$  for  $-\frac{\pi}{2} < x < \frac{\pi}{2}$

- (a) Find the value of A (b) Find the probability that the particle be found between  $x=0$  and  $x = \frac{\pi}{4}$ .

[GGSIPU, Feb. 2012, (5 marks) reappear]

**Solution.** As  $\psi = A \cos^2 x$   $-\frac{\pi}{2} < x < \frac{\pi}{2}$

(a)  $\int_{-\pi/2}^{\pi/2} |\psi|^2 dx = 1$

$$\Rightarrow 2A^2 \int_0^{\pi/2} \cos^4 x dx = 1 \quad \text{or} \quad 2A^2 \frac{3\pi}{16} = 1$$

$$\frac{3\pi}{8} A^2 = 1 \quad \therefore \quad A = \sqrt{\frac{8}{3\pi}}$$

(b)  $P = \int_0^{\pi/4} |\psi|^2 dx = A^2 \int_0^{\pi/4} \cos^4 x dx$  as  $A = \sqrt{\frac{8}{3\pi}}$

$$P = \frac{8}{3\pi} \int_0^{\pi/4} \cos^4 x dx = 0.462$$

## Schrodinger wave Equation for finite well:-

(Skip barrier / Potential barrier & tunnel Effect).

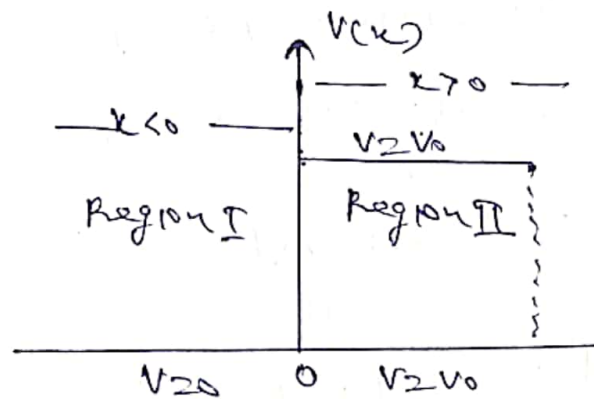
### 1) Step barrier / Potential barrier.

When a particle approaches a region in which entrance of particle is opposed by some force then that region is said to form a barrier for a particle.

Note:- Potential barrier is also a system used to identify incident, reflected & transmitted matter waves.

Let us consider a particle moving along  $x$ -axis from left to right and meeting a barrier of height  $V_0$ .

Let  $E$  be the total energy of particle and  $V$  be the value of constant potential.



We have two different regions supported by discontinuous change of potential -

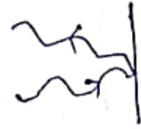
Region I :  $x < 0 \Rightarrow V = 0$

$x > 0 \Rightarrow V = V_0$

Behaviour of particle can be described in quantum mechanics as if energy of particle -

1) If  $(E < V_0)$

Then there is a probability that the particle will be reflected back from the barrier



2) If  $(E > V_0)$

Then there is a probability that it will penetrate through the barrier.



Such behaviour of particle is impossible from the classical viewpoint. It follows Schrodinger Wave Equation

a) Region I

Let us consider the region  $x < 0$  time independent Schrodinger wave Equation is given by -

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2} (E - V) \psi = 0$$

( $V=0$  in the region  $x < 0$ )

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2} E \psi = 0$$

$$\text{Let } k_1^2 = \frac{8\pi^2mE}{h^2}$$

$$\Rightarrow \frac{d^2\psi}{dx^2} + k_1^2 \psi = 0 \quad \text{--- (1)}$$

b) Region II

Schrodinger wave Equation is given by -

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2} (E - V_0) \psi = 0$$

$$\frac{d^2\psi}{dx^2} + k_2^2 \psi = 0$$


$$\text{where } k_2^2 = \frac{8\pi^2m}{h^2} (E - V_0) \quad \text{--- (2)}$$

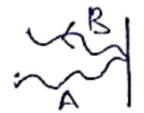
Solution of Differential Equation (1) & (2) are -

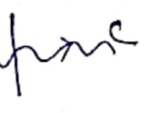
$$\psi_1 = A e^{ik_1 x} + B e^{-ik_1 x}$$

$$\psi_2 = C e^{ik_2 x} + D e^{-ik_2 x}$$

Where -

A - Amplitude of incident wave in Region I 

B - Amplitude of reflected wave in Region I 

C - Amplitude of wave penetrate the barrier in Region II 

D - No particle can flow to left hence  $D=0$

Now In order to be a proper wave function  $\psi$  must be finite & continuous everywhere on  $x$ -axis.

It means at boundary ( $x=0$ ) wave function inside & outside must have same values.

Hence  $\psi_1, \psi_2$  & its derivatives must be continuous at the boundary ( $x=0$ )

$$\psi_1 = \psi_2 \quad (\text{at } x=0)$$

$$\Rightarrow A e^{ik_1 x} + B e^{-ik_1 x} = C e^{ik_2 x}$$

substituting  $x=0$

$$\Rightarrow A + B = C \quad \text{--- (3)}$$

Also its derivatives must be continuous at  $x=0$

$$\Rightarrow \frac{d\psi_1}{dx} = \frac{d\psi_2}{dx}$$

$$\frac{d(A e^{ik_1 x} + B e^{-ik_1 x})}{dx} = \frac{d(C e^{ik_2 x})}{dx}$$

$$\Rightarrow A i k_1 e^{ik_1 x} - B i k_1 e^{-ik_1 x} = C i e^{ik_2 x}$$

Now at  $x=0$  we have

$$A i k_1 - B i k_1 = C i k_2$$

$$\Rightarrow (A-B) = \frac{k_2}{k_1} C \quad \text{--- (4)}$$

on solving (3) & (4) we get

$$2B = \left(1 - \frac{k_2}{k_1}\right) C$$

$$= \frac{k_1 - k_2}{k_1} \times \frac{2k_1}{k_1 + k_2} \times A$$

$$\Rightarrow \boxed{B = \left(\frac{k_1 + k_2}{k_1 - k_2}\right) A}$$

Note  
For solving  
Add (3) + (4)  
& Subtract (3) - (4)

Also  $2A = \left(1 + \frac{k_2}{k_1}\right) C$

$$\Rightarrow \boxed{A = \frac{k_1 + k_2}{2k_1} C}$$

Now substituting value of A in (3)

we get  $\boxed{C = \frac{2k_1}{k_1 + k_2} A}$

Now case I For  $E < V_0$

Probability density is given by

$$(J_x) = \frac{\hbar}{2im} \left( \psi_1^* \frac{d\psi_1}{dx} - \psi_1 \frac{d\psi_1^*}{dx} \right)$$

we have  $\psi_1 = A e^{ik_1 x} + B e^{-ik_1 x}$

$$\psi_1^* = A e^{-ik_1 x} + B e^{ik_1 x}$$

$$\Rightarrow (J_x) = \frac{\hbar}{2im} \left[ (A e^{-ik_1 x} + B e^{ik_1 x}) \times \frac{d}{dx} (A e^{ik_1 x} + B e^{-ik_1 x}) \right. \\ \left. - [(A e^{ik_1 x} + B e^{-ik_1 x}) \times \frac{d}{dx} (A e^{-ik_1 x} + B e^{ik_1 x})] \right]$$

$$= \frac{\hbar}{2im} \left( (Ae^{-ik_1x} + Be^{ik_1x}) \times (ik_1) (Ae^{ik_1x} - Be^{-ik_1x}) \right) \\ \rightarrow \left( (Ae^{ik_1x} + Be^{-ik_1x}) \times (-ik_1) (Ae^{-ik_1x} - Be^{ik_1x}) \right) \\ = \frac{k_1 \hbar}{m} (A^2 - B^2)$$

Here  $A^2$  represent magnitude of incident current  
 $B^2$  - magnitude of Reflected current

Reflection coefficient =  $\frac{\text{Magnitude of Reflected current}}{\text{Magnitude of incident current}}$

$$R = \frac{B^2}{A^2} \\ = \frac{\left( \frac{k_1 - k_2}{k_1 + k_2} \right)^2 A^2}{A^2} = \left( \frac{k_1 - k_2}{k_1 + k_2} \right)^2$$

Case - II Region II

$$\psi_2 = ce^{ik_2x} \quad \psi_2^* = ce^{-ik_2x}$$

$$J_x = \frac{\hbar}{2im} \left( \psi_2^* \frac{d\psi_2}{dx} - \psi_2 \frac{d\psi_2^*}{dx} \right)$$

$$= \frac{\hbar}{2im} \left( ce^{-ik_2x} \frac{d}{dx} (ce^{+ik_2x}) - (ce^{+ik_2x}) \frac{d}{dx} (ce^{-ik_2x}) \right) \\ = \frac{\hbar}{2im} (cxc(ik_2) - cxc(-ik_2))$$

$$\boxed{J = 0}$$

How Transmittance coefficient =

Magnitude of Transmittance current

Magnitude of incident current

$$= \frac{c^2}{A^2}$$

$$= \frac{\left(\frac{2k_1}{k_1+k_2}\right)^2 A^2}{A^2} \times \frac{k_2}{k_1}$$

$$T = \frac{4k_1 k_2}{(k_1+k_2)^2}$$

Also according to definition  $R+T=1$  can be verified.

## Tunneling Effect:

When a particle is able to cross the potential barrier even when ~~the~~ its energy is less than the barrier height ( $E < V_0$ ), then this phenomenon is called "Tunnel Effect".

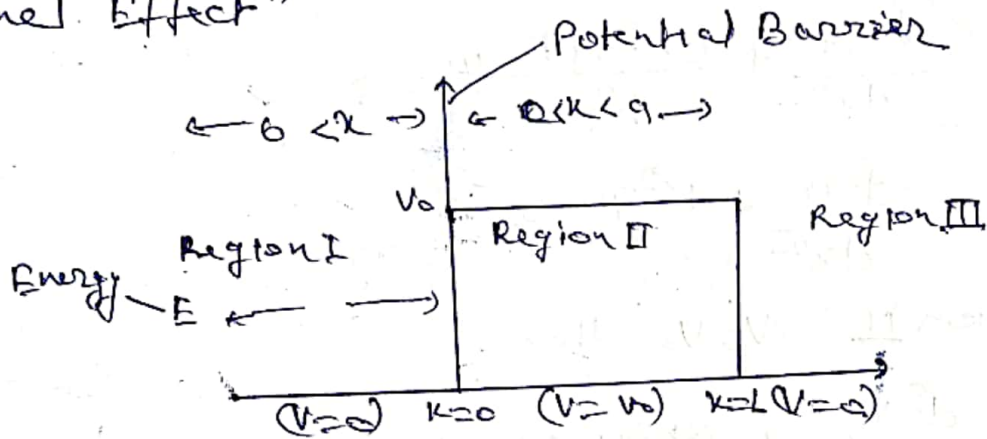


Figure represent a potential barrier of height  $V_0$  & thickness  $L$ . Let us consider a particle moving along  $x$ -axis from left to right and moving a rectangular potential barrier of height  $V_0$  and length  $L$ .

We have three different region separated by discontinuous change of potential energy.

- Potential energy is zero in the region I & III for  $x < 0$  &  $x > L$
- potential energy is  $V_0$  in the Region II.

If  $E < V_0$  according to quantum mechanics, still there is a finite chance for the electron to leak the other side of barrier. Hence the electron is tunneled through the potential barrier. This phenomenon called Tunneling.

The Schrodinger wave Equation for a particle along X-AXIS is given by -

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2} (E-V)\psi = 0$$

In Region I  $V=0$  then

$$\frac{d^2\psi_1}{dx^2} + \frac{8\pi^2mE}{h^2} \psi_1 = 0$$

$$\frac{d^2\psi_1}{dx^2} + \frac{k^2}{h^2} \psi_1 = 0 \quad \text{--- (1)} \quad \left( \begin{array}{l} k^2 = \frac{8\pi^2mE}{h^2} \\ k = \sqrt{\frac{2mE}{h}} \end{array} \right.$$

In Region II  $V=V_0$  then

$$\frac{d^2\psi_2}{dx^2} + \frac{8\pi^2m}{h^2} (E-V_0)\psi_2 = 0$$

$$\frac{d^2\psi_2}{dx^2} - \frac{8\pi^2m}{h^2} (V_0-E)\psi_2 = 0$$

$$\frac{d^2\psi_2}{dx^2} - \beta^2 \psi_2 = 0 \quad \text{--- (2)}$$

$$\beta = \sqrt{\frac{8\pi^2m}{h^2} (V_0-E)} = \sqrt{\frac{2m(V_0-E)}{h}}$$

(Note:- In Region II wave is decaying exponentially so Negative value of  $E$  in  $(V_0-E)$  is used Hence solution of this Eqn ( $\psi_2$ ) doesn't give i (imaginary) term)

Region -III  $V=0$  then

$$\frac{d^2\psi_3}{dx^2} + k^2 \psi_3 = 0 \quad \text{--- (3)}$$

Solution of Equation (1), (2) & (3) will be

$$\psi_1 = Ae^{ikx} + Be^{-ikx}$$

$$\psi_2 = Ce^{\beta x} + De^{-\beta x}$$

$$\begin{aligned}\psi_3 &= Ge^{ikx} + He^{-ikx} \\ &= Ge^{ikx}\end{aligned}$$

H=0 because no wave travel back from infinity

Applying boundary conditions to particle

1) At  $x=0$   $\psi_1 = \psi_2$

$$Ae^{ikx} + Be^{-ikx} = Ce^{\beta x} + De^{-\beta x}$$

At  $x=0$

$$A+B = C+D \quad \text{--- (4)}$$

We also have -

$$\frac{d\psi_1}{dx} = \frac{d\psi_2}{dx}$$

$$ikAe^{ikx} - ikBe^{-ikx} = \beta Ce^{\beta x} - \beta De^{-\beta x}$$

at  $x=0$   $ik(A-B) = \beta(C-D)$

$$A-B = \frac{\beta}{ik}(C-D) \quad \text{--- (5)}$$

From equation (4) & (5) we get

$$A = \frac{1}{2} \left[ \left(1 + \frac{\beta}{ik}\right)C + \left(1 - \frac{\beta}{ik}\right)D \right] \quad \text{--- (6)}$$

$$B = \frac{1}{2} \left[ \left(1 - \frac{\beta}{ik}\right)C + \left(1 + \frac{\beta}{ik}\right)D \right] \quad \text{--- (7)}$$

At  $x=0$  we have

$$\psi_2 = \psi_3$$
$$C e^{\beta x} + D e^{-\beta x} = G e^{ikx}$$

At  $x=a$

$$\Rightarrow C e^{\beta a} + D e^{-\beta a} = G e^{ika} \quad \text{--- (8)}$$

Also we have -

$$\frac{d\psi_2}{dx} = \frac{d\psi_3}{dx}$$

$$C \beta e^{\beta x} - D \beta e^{-\beta x} = G i k e^{ikx}$$

At  
 $x=a$

$$C \beta e^{\beta a} - D \beta e^{-\beta a} = \frac{ik}{\beta} G e^{ika} \quad \text{--- (9)}$$

From eqn (8) and (9)

$$C = \frac{1}{2} \left( 1 + \frac{ik}{\beta} \right) e^{-\beta a} G e^{ika}$$

$$D = \frac{1}{2} \left( 1 - \frac{ik}{\beta} \right) e^{\beta a} G e^{ika}$$

Put the value of C + D in eqn (6)

From (6) we have -

$$A = \frac{1}{2} \left( \left( 1 + \frac{\beta}{ik} \right) C + \left( 1 - \frac{\beta}{ik} \right) D \right)$$

$$A = \frac{1}{4} \left( \left( 1 + \frac{\beta}{ik} \right) \left( 1 + \frac{ik}{\beta} \right) e^{-\beta a} G e^{ika} \right)$$

$$+ \frac{1}{4} \left( \left( 1 - \frac{\beta}{ik} \right) \left( 1 - \frac{ik}{\beta} \right) e^{\beta a} G e^{ika} \right)$$

$$A = \frac{G}{4} e^{ika} \left[ \left( 1 + \frac{ik}{\beta} + \frac{\beta}{ik} + \frac{\beta ik}{\beta ik} \right) e^{-\beta a} + \left( 1 - \frac{ik}{\beta} - \frac{\beta}{ik} - \frac{ik \beta}{ik \beta} \right) e^{\beta a} \right]$$
$$= \frac{G}{4} e^{ika} \left( 2 + \frac{ik}{\beta} + \frac{\beta}{ik} \right) e^{-\beta a} + \left( 2 - \frac{ik}{\beta} - \frac{\beta}{ik} \right) e^{\beta a}$$

$$\begin{aligned}
 A &= \frac{4}{4} e^{ika} \left( 2C e^{-\beta a} + e^{\beta a} \right) + \left( \frac{\beta}{ik} + \frac{ik}{\beta} \right) e^{-\beta a} - \left( \frac{\beta}{ik} + \frac{ik}{\beta} \right) e^{\beta a} \\
 &= \frac{4}{4} e^{ika} \left( 2 \frac{(e^{-\beta a} + e^{\beta a})}{2} + \left( \frac{\beta}{ik} + \frac{ik}{\beta} \right) \cdot (e^{-\beta a} - e^{\beta a}) \right) \\
 &= \frac{4}{4} e^{ika} \left( 2 \frac{(e^{-\beta a} + e^{\beta a})}{2} + 2 \left( \frac{\beta}{ik} + \frac{ik}{\beta} \right) \left( \frac{e^{\beta a} - e^{-\beta a}}{2} \right) \right) \\
 &= \left[ \cosh \beta a + \frac{i}{2} \left( \frac{\beta}{k} - \frac{k}{\beta} \right) \sinh \beta a \right] 4e^{ika}
 \end{aligned}$$

$$\Rightarrow \frac{A}{4} = \left[ \cosh \beta a + \frac{i}{2} \left( \frac{\beta}{k} - \frac{k}{\beta} \right) \sinh \beta a \right] e^{ika} \quad \text{--- (10)}$$

We have Transmittance coefficient (T)

$$(T) = \frac{\text{Magnitude of Transmitted wave}}{\text{Magnitude of Incident wave}} = \left( \frac{4^2}{A^2} \right)$$

From (10)

$$\left| \frac{A^*A}{4^2 4} \right| = \left| \frac{A}{4} \right|^2 = \frac{1}{4} = \cosh^2 \beta a + \frac{1}{4} \left( \frac{\beta}{k} - \frac{k}{\beta} \right)^2 \sinh^2 \beta a$$

$$= 1 + \left( 1 + \frac{1}{4} \left( \frac{\beta}{k} - \frac{k}{\beta} \right)^2 \right) \sinh^2 \beta a$$

$$= 1 + \left[ \frac{1}{4} \left( 4 + \left( \frac{\beta}{k} - \frac{k}{\beta} \right)^2 \right) \sinh^2 \beta a \right] \quad \text{--- (11)}$$

We have  $k = \frac{\sqrt{2mE}}{\hbar}$  &  $\beta = \frac{\sqrt{2m(V_0 - E)}}{\hbar}$

$$\Rightarrow \frac{\beta}{k} - \frac{k}{\beta} = \sqrt{\frac{V_0 - E}{E}} - \sqrt{\frac{E}{V_0 - E}} = \frac{V_0 - E - E}{\sqrt{E(V_0 - E)}} = \frac{V_0 - 2E}{\sqrt{E(V_0 - E)}}$$

Now  $\left( \frac{\beta}{k} - \frac{k}{\beta} \right)^2 + 4 = \frac{(V_0 - 2E)^2}{E(V_0 - E)} + 4 = \frac{V_0^2 + 4E^2 - 4EV_0 + 4EV_0 - 4E^2}{E(V_0 - E)} = \frac{V_0^2}{E(V_0 - E)}$

$$2) \quad \frac{1}{T} = 1 + \frac{1}{4} \left[ \frac{V_0^2}{E(V_0 - E)} \right] \sinh^2 \beta a$$

$$T = \frac{1}{1 + \frac{V_0^2}{4E(V_0 - E)}} \sinh^2 \beta a$$

Also  $\sinh^2 \beta a = \left( \frac{e^{\beta a} - e^{-\beta a}}{2} \right)^2 = e^{\frac{2\beta a}{4}}$

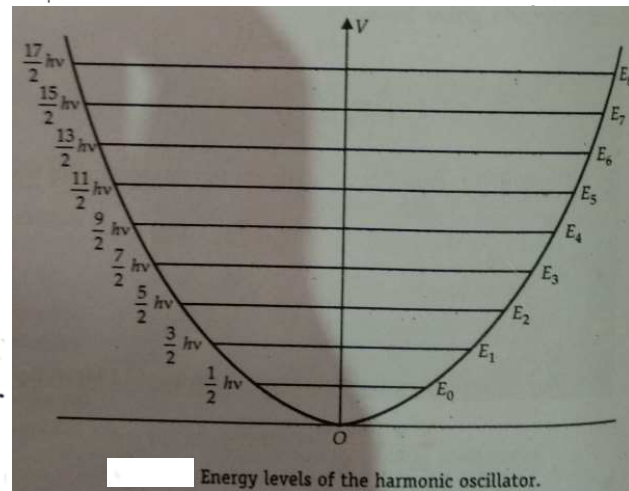
$$2) \quad T = \frac{1}{1 + \frac{V_0^2}{4E(V_0 - E)}} \times e^{\frac{2\beta a}{4}}$$

↓  
Transmission coefficient.

## Simple Harmonic oscillator & Zero Point Energy!

A simple harmonic oscillator is a particle performing one dimensional motion under restoring force  $F$  which is proportional to particle displacement  $x$  from equilibrium position.

The plot of  $V$  against  $x$  is a parabola and we may describe that the particle being a periodic potential well. If  $E$  be the total energy of particle. It oscillate back & forth. from center.



Now force is proportional to particle displacement  $x$  from equilibrium position

$$F \propto -x$$

$$F = -kx$$

Potential Energy of harmonic oscillator is,

$$V(x) = - \int_0^x F(x) dx$$

$$= \int_0^x kx dx$$

$$V(x) = \frac{1}{2} kx^2$$

$$k = m\omega^2$$

$$V(x) = \frac{1}{2} m\omega^2 x^2 = 2\pi^2 m\nu^2 x^2 \quad \text{--- (1)}$$

Schrodinger wave Equation for particle motion is given by

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m(E-V)}{h^2} \psi = 0$$

Substituting value of ①

$$\frac{d^2\psi}{dx^2} + \frac{8\pi^2m}{h^2} (E - 2\pi^2mv^2x^2) \psi = 0$$

$$\frac{d^2\psi}{dx^2} + \left( \frac{8\pi^2mE}{h^2} - \frac{16\pi^4m^2v^2x^2}{h^2} \right) \psi = 0$$

$$\text{Let } \frac{8\pi^2mE}{h^2} = \alpha \quad \text{and} \quad \frac{4\pi^2mv}{h} = \beta$$

$$\beta^2 = \frac{16\pi^4m^2v^2}{h^2}$$

$$\Rightarrow \frac{d^2\psi}{dx^2} + (\alpha - \beta^2x^2)\psi = 0 \quad \text{--- (2)}$$

$$\text{Let } y = \sqrt{\beta}x$$

$$\text{on sq. } y^2 = \beta x^2$$

$$x^2 = \frac{y^2}{\beta}$$

$$\text{in (2)} \quad \frac{d^2\psi}{dx^2} + \left( \alpha - \beta^2 \frac{y^2}{\beta} \right) \psi = 0$$

$$\frac{d^2\psi}{dx^2} + (\alpha - \beta y^2) \psi = 0 \quad \text{--- (3)}$$

Now we have  $y = \sqrt{\beta}x$

$$\frac{d\psi}{dx} = \frac{d\psi}{dy} \times \frac{dy}{dx} = \frac{d\psi}{dy} \times \sqrt{\beta}$$

$$\Rightarrow \frac{d^2\psi}{dx^2} = \beta \frac{d^2\psi}{dy^2}$$

Substituting in (3)

$$\Rightarrow \beta \frac{d^2\psi}{dy^2} + (\alpha - \beta y^2) \psi = 0$$

$$\frac{d^2\psi}{dy^2} + \left( \frac{\alpha}{\beta} - y^2 \right) \psi = 0$$

--- (4)

The Solution of Equation (4)

will satisfy the condition  $\frac{\alpha}{\beta} = (2n+1)$

$$\alpha = \frac{8\pi^2 m E}{h^2} \quad \beta = \frac{4\pi^2 m \nu}{h}$$

$$\frac{\frac{8\pi^2 m E}{h^2}}{\frac{4\pi^2 m \nu}{h}} = (2n+1)$$

$$\Rightarrow E = \frac{h\nu}{2} (2n+1)$$

$$E_n = h\nu \left( n + \frac{1}{2} \right)$$

The lowest value of Energy which oscillation can take for  $n=0$  is

$$E_0 = \frac{1}{2} h\nu$$

- \* This level is ground state. The Energy  $E_0$  is called zero point Energy.
- \* A very significant result that  $E_0$  is never zero but has a lowest or minimum value.
- \* A harmonic oscillator approaches this value as the temperature approaches 0K. This energy called zero point energy.
- \* The existence of zero point Energy is a consequence of Uncertainty Principle. At ~~T=0K~~ T=0K position & momentum can be determined precisely.

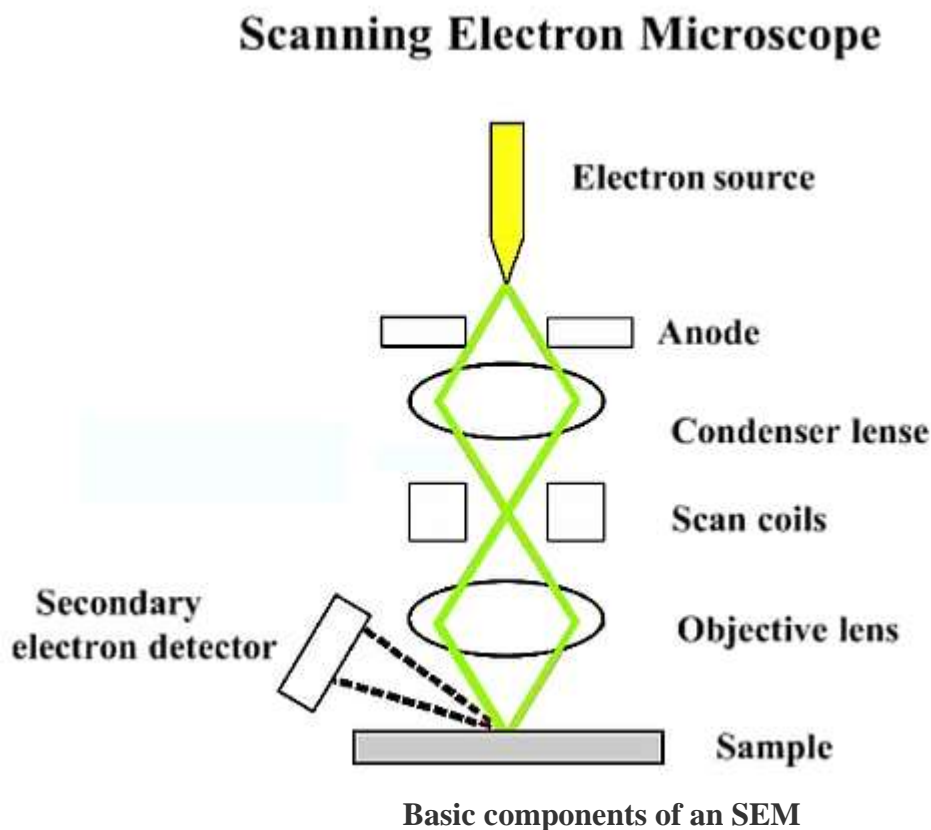
## Scanning electron microscopes

Scanning electron microscopes (SEMs) have become powerful and versatile tools for material characterization, especially in recent years, as the size of materials used in various applications.

Electron microscopes use electrons for imaging in a similar way that light microscopes use visible light. SEMs use the electrons that are reflected or knocked off the near-surface region of a sample to create an image. Since the wavelength of electrons is much smaller than that of light, the resolution of SEMs is superior to that of a light microscope.

### How SEMs work

In scanning electron microscopy, the electron beam scans the sample in a raster pattern. First, electrons are generated at the top of the column by the **electron source**.



Due to very narrow electron beam SEM Micrograph have large depth of field giving structure of sample.

The scan coil produces a magnetic field which deflect the electron beam in a controlled way. Varying voltage is applied to the coil of the cathode ray tube which produce pattern of light.

When the electron beam hit the sample, it produces secondary electron.

The entire electron column needs to be under vacuum. Like all components of an electron microscope, the electron source is sealed inside a special chamber to preserve vacuum and protect it against contamination, vibrations, and noise.

Advantages: 1. Bulk/powder samples can be used

2. Provide much greater depth of view so it can produce image that can represent 3D structure of sample.

## UNIT-2 QUANTUM STATISTICS

Statistics:- Collection and analysis of data is known as statistics.

Quantum Statistics:- A statistical description of a system of particles that obey rules of quantum mechanics.

Mechanics:- Science concerned with the motion of bodies under action of force.

Quantum Mechanics:- Deals with smallest thing in universe.

Eg. Electron, Proton, Neutron etc.

### Classical Statistics & Quantum Statistics:-

Classical Statistics:- In classical statistics all particles in a system are considered distinguishable. Means individual particles in a system can be tracked. As a result changing the position of any two particles leads to completely different configuration.



( A, B, C are distinguishable particles inside a system )

Quantum Statistics:- In quantum statistics, all particles in a system are considered indistinguishable. i.e. we can't label the particle. Interchanging the particle doesn't change the system.

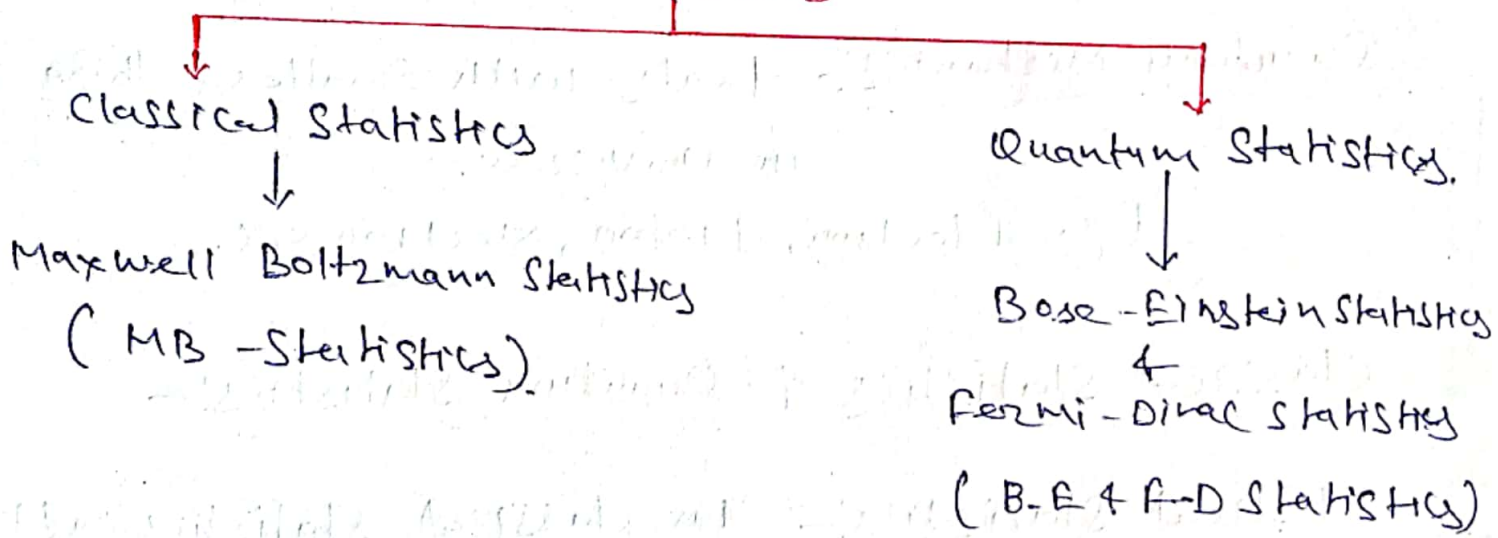
Note:- Apply on Non-observable.

Eg. -



(A, A, A are indistinguishable particles in a system)

### STATISTICS



### Examples

I) M-B Statistics:- Gas molecule

II) B-E Statistics:- photons or Bosons.

III) F-D Statistics:- Electrons

311 COMPARISON OF MAXWELL-BOLTZMANN, BOSE-EINSTEIN AND FERMI-DIRAC STATISTICS

	<i>Particular</i>	<i>M.B. Statistics</i>	<i>B.E. Statistics</i>	<i>F.D. Statistics</i>
1.	<i>Nature of particles</i>	Particles are identical but distinguishable.	Particles are identical but indistinguishable.	Particles are identical but indistinguishable.
2.	<i>Size of the phase cell</i>	The available volume of the phase space cell can be small as we like and can even approach zero.	Minimum size of the available phase space volume is of the order of $h^3$ , where $h$ is the Planck's constant.	Minimum size of the available phase space volume is of the order of $h^3$ .
3.	<i>Number of phase cells</i>	The phase space can be divided into any large number of cells of our choice.	The phase space can be divided into number of cells whose number is limited as the phase space volume of each cell cannot be less than $h^3$ .	The phase space can be divided into the number of cells, whose number is limited as the phase space volume of each cell cannot be less than $h^3$ .
4.	<i>Number of particles in a cell</i>	Any number of particles can occupy a single phase space cell.	Any number of particles can occupy a single phase cell.	A phase space cell cannot accommodate more than one particle.
5.	<i>Macrostates and microstates</i>	A macrostate can have more than one microstates.	Each macrostate has one microstate.	Each macrostate has one microstate.
6.	<i>Energy</i>	Continuous distribution of energy is assumed.	Energy is quantized.	Energy is quantized.
7.	<i>Spin</i>	The particles are spinless.	The particles have integral spin.	The particles have half integral spin.
8.	<i>Occupation index</i>	$\frac{1}{e^{\alpha+\beta E_i}}$	$\frac{1}{e^{\alpha+\beta E_i} - 1}$	$\frac{1}{e^{\alpha+\beta E_i} + 1}$
9.	<i>Total energy</i>	$\frac{3}{2} Nk_B T$	$\frac{3}{2} Nk_B T \left( 1 - \frac{1}{2^{5/2} e^{-\alpha}} \right)$	$\frac{3}{2} Nk_B T \left( 1 + \frac{1}{2^{5/2} e^{-\alpha}} \right)$
10.	<i>Wave function</i>		Symmetric under interchange of the coordinates of any two <i>Bosons</i> .	Antisymmetric on interchange of the coordinates of any two <i>Fermions</i> .

## Emissive Power:-

Emissive power of a body at a given temperature is the quantity of radiant energy emitted by the body per unit time per unit area of the body at that temperature.

## Emissivity ( $\epsilon$ )

Emissivity is defined as the ratio of energy radiated from a material surface to the radiated energy of perfectly black body at that temperature.

## Stefan's Boltzmann law:-

It states that total radiations emitted from a black body per unit area is proportional to fourth power of its absolute temperature.

(Absolute Temperature - Temperature of Black body)

$$E \propto T^4$$

$$E = \sigma T^4$$

where  $\sigma$  - Stefan's constant

$$\sigma = 5.67 \times 10^{-8} \text{ W/m}^2\text{-K}^4$$

$$E = \frac{Q}{At}$$

$$Q = At \times E$$

$$Q = At \times (\sigma T^4)$$

$Q$  - Amount of heat emitted by a black body.

## Black body Radiation Spectrum:-

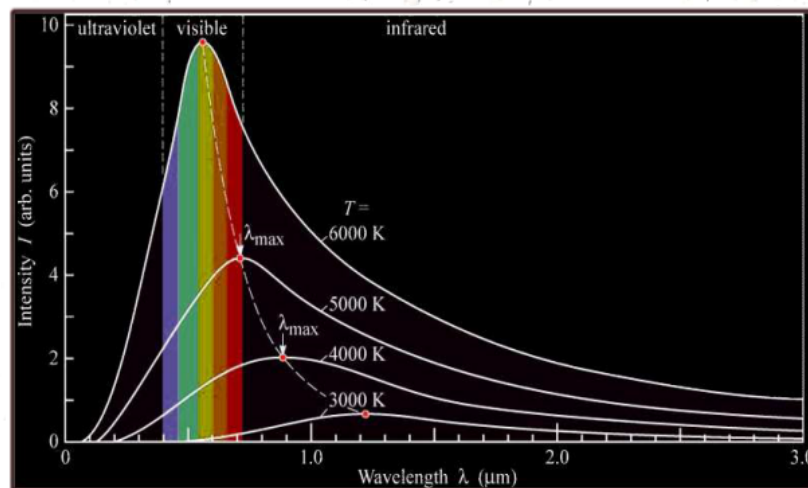
In practice there is not perfect black bodies, Scientist made a blackbody in laboratory by taking hollow sphere & drilling a small hole in it, which can absorb 99% of Radiations.

Inner surface was coated with carbon black when Radiations entering inside hit the wall of sphere and get absorbed 99%.

Remaining radiations reflected back and collide with wall again.

Due to Absorption of heat its temperature inside increase and it starts emitting Radiations from narrow hole. This is called blackbody Radiations.

When radiations from black body traced on a graph classical mechanics fails to explain graph.



- I) The blackbody radiation spectrum is not having uniform distribution of Energy wrt wavelength.
- ii) At a given temperature different wavelength radiations have different energies.
- iii) Total Energy of Radiations at any temperature is given by area under curve. The area is

is directly proportional to fourth power of its absolute temperature  $E \propto T^4$  (Stefan's law)

iv) The wavelength corresponding to maximum energy shift towards shorter wavelength with increase in temperature.  $\lambda_m T = \text{constant}$

(Wien's Displacement law)

$\lambda_m$  - wavelength corresponds to maximum Energy.

## Wien's Displacement Law:-

Wien's displacement law is the relationship between temperature of a blackbody and wavelength at which it emits light.

Wien's Displacement law states that the blackbody Radiation curve for different temperature peaks at a wavelength is inversely proportional to temperature.

$$\lambda_m \propto \frac{1}{T}$$

$$\lambda_m = \frac{b}{T}$$

b - Wien's constant

$$b = 2.898 \times 10^{-3} \text{ meter-kelvin (MK)}$$

## Drawbacks:-

This law is applicable only for short wavelength region & for high temperature of the source of Radiations.

It fails to explain decrease in Energy for longer wavelength.

## Rayleigh Jeans Law:-

Rayleigh Jeans law Express the Energy distribution in the spectrum as a function of Temperature.

$$E_\nu d\nu = \frac{8\pi \nu^2 KT}{c^3} d\nu$$

According to this law Radiation emitted by a particle in a blackbody travel as a wave.

Average Energy per unit vibration =  $KT$

$K$  - Boltzmann constant

$T$  - Temperature.

In term of wavelength Energy density within the range  $\lambda$  to  $\lambda + d\lambda$  is -

Note

$$\left[ \begin{array}{l} \nu = \frac{c}{\lambda} \\ d\nu = -\frac{c}{\lambda^2} d\lambda \\ \text{Since increase in frequency correspond to} \\ \text{decrease in wavelength} \\ E_\nu d\nu = -E_\lambda d\lambda \end{array} \right]$$

inc 1

$$E_\lambda d\lambda = \frac{8\pi \left(\frac{c}{\lambda}\right)^2 KT}{c^3} \times \left(\frac{c}{\lambda^2}\right) d\lambda$$

$$E_\lambda d\lambda = \frac{8\pi KT d\lambda}{\lambda^4}$$

↓

In term of wavelength.

Draw backs:-

Rayleigh Jeans law agree Experimental Results at large wave length (low frequencies). But strongly disagree at shorter wavelength (high frequencies).

## Planck's law of Radiations

After failure of Wien's law & Rayleigh law  
"Planck's introduced an entirely new idea to explain the distribution of Energy among various wavelength of blackbody Radiation."

Following Assumptions were made -

- I) Blackbody is made up of large number of oscillating particles. These particles vibrate in possible frequencies.
- II) Atoms of the walls of uniform temperature behaves as oscillator each with characteristic frequency of oscillation.
- III) An oscillator can have discrete Energy given by -  
$$E = nh\nu$$

$\nu$  - frequency of oscillator.  
 $h$  - Planck's constant  
 $n$  - an Integer quantum number.
- IV) An oscillator do not emit or absorb energy continuously but only in jump.

Let  $N$  be the total number of oscillators &  $E$  be the total Energy of these oscillators, then Average Energy is given by

$$\bar{E} = \frac{E}{N}$$

Now let  $N_0, N_1, N_2, \dots, N_n$  be the number of oscillator having energies  $0, h\nu, 2h\nu, \dots, nh\nu$ .

By Max distribution formula we have

$$N_n = N_0 e^{-nh\nu/KT}$$

According to Maxwell Boltzmann Distribution law -  
 Number of oscillator ( $N_n$ ) with Energy ( $E_n$ ) is given by.

$$N_n = N e^{-E_n/KT}$$

Average Energy of oscillator is given by -

$$\langle E \rangle = \frac{\sum_{n=0}^{\infty} E_n N_n}{\sum_{n=0}^{\infty} N_n}$$

$$= \frac{\sum_{n=0}^{\infty} E_n N e^{-E_n/KT}}{N \sum_{n=0}^{\infty} e^{-E_n/KT}}$$

$$= \frac{\sum_{n=0}^{\infty} nh\nu e^{-nh\nu/KT}}{e^{-nh\nu/KT}}$$

$$= \frac{0 + h\nu e^{-h\nu/KT} + 2h\nu e^{-2h\nu/KT} + \dots}{1 + e^{-h\nu/KT} + e^{-2h\nu/KT} + \dots}$$

Let  $e^{-h\nu/KT} = x$

$$\Rightarrow \langle E \rangle = \frac{h\nu x (1 + 2x + 3x^2 + \dots)}{1 + x + x^2 + \dots}$$

$$= \frac{h\nu x (1-x)^{-2}}{(1-x)^{-1}}$$

$$= \frac{h\nu x}{(1-x)}$$

$$= \frac{h\nu e^{-h\nu/KT}}{(1 - e^{-h\nu/KT})}$$

$$E = \frac{h\nu}{e^{h\nu/KT} - 1}$$

ⓐ

Average Energy of Planck's oscillation

## Planck's Radiation Formula:-

Number of oscillation in the frequency range  $\nu$  to  $\nu+d\nu$  are estimated to be -

$$\frac{8\pi\nu^2 d\nu}{c^3}$$

The Energy density of Radiations  $E_\nu$  in the frequency Range  $\nu$  to  $\nu+d\nu$  is related to average Energy of oscillator as given by.

$$E_\nu d\nu = \frac{8\pi\nu^2 d\nu}{c^3} \times E$$

$$= \frac{8\pi\nu^2 d\nu}{c^3} \times \frac{h\nu}{e^{h\nu/kT} - 1}$$

(Answer 1)

$$E_\nu d\nu = \frac{8\pi h \nu^3}{c^3} \cdot \frac{1}{e^{h\nu/kT} - 1} d\nu$$

↓

Planck's Radiation Law.

In term of wavelength:-

Note  $\nu = c/\lambda$   
 $d\nu = -\frac{c}{\lambda^2} d\lambda$   
Since increase in frequency correspond to decrease in wavelength  
 $E_\lambda d\lambda = -E_\nu d\nu$

$$E_\lambda d\lambda = \frac{8\pi h \left(\frac{c}{\lambda}\right)^3 \left(\frac{c}{\lambda^2}\right) d\lambda}{c^3 \times e^{hc/\lambda kT} - 1}$$

$$E_\lambda d\lambda = \frac{8\pi h c}{\lambda^5} \cdot \frac{d\lambda}{e^{hc/\lambda kT} - 1}$$

\* Ink  
Wien's law and Rayleigh law from Planck's law!

Wien's law and Rayleigh can be derived from Planck's law:-

From Planck's law we have (by ②)

$$u_{\lambda} d\lambda = \frac{8\pi hc}{\lambda^5} \cdot \frac{d\lambda}{e^{hc/\lambda kT} - 1}$$

i) When wave length  $\lambda$  is very small then  $e^{hc/\lambda kT} \gg 1$  so Planck's formula can be (Equation ②) written as - (Equation ③)

$$u_{\lambda} d\lambda = \frac{8\pi hc}{\lambda^5} e^{-hc/\lambda kT} d\lambda$$

Putting  $8\pi hc = A$  &  $\frac{hc}{k} = B$  we have

$$u_{\lambda} d\lambda = \frac{A}{\lambda^5} e^{-B/\lambda T} d\lambda$$

Wien's law.

ii) When  $\lambda$  is very large then -

$$e^{hc/\lambda kT} = 1 + \frac{hc}{\lambda kT}$$

So Equation ② can be written as -

$$u_{\lambda} d\lambda = \frac{8\pi hc}{\lambda^5 \left(1 + \frac{hc}{\lambda kT} - 1\right)} d\lambda$$

$$u_{\lambda} d\lambda = \frac{8\pi kT}{\lambda^4} d\lambda$$

Rayleigh's law.

## Photo electric Effect

"When a beam of light of suitable frequency falls on a metal plate slow moving electrons are emitted from the metal surface. This phenomenon is called photoelectric effect"

Some of the basic features of photoelectric effect are -

- I) The lowest frequency of light which can emit photoelectron from a material is called threshold frequency of that material. ( $\nu_0$ )
- II) Photoelectric current or number of photoelectrons emitted per second is directly proportional to the intensity of incident light.
- III) Velocity of photoelectron is completely independent on intensity of incident light but depend on its frequency.
- IV) The Negative potential  $V_0$  at which the photoelectric current is reduced to zero called stopping potential or cut off potential.

When light of frequency  $\nu$  falls on metallic surface a part of its energy is used to liberate the electron from the surface known as work-function and rest energy is carried by emitted electron as kinetic energy.

$$h\nu = \frac{1}{2}mv^2 + \phi_0$$

Kinetic Energy                      work Function.

$$h\nu = \frac{1}{2}mv^2 + h\nu_0$$

( $\nu_0$  - Threshold Frequency)

$$\boxed{\frac{1}{2}mv^2 = h(\nu - \nu_0)}$$
 - Einstein photoelectric Equation.

## Concept of Statistical Mechanics:

Microstate: A state of a system where all the parameters of a system particle are specified.

Many microstates exist for each of the systems, and there are many parameters for each state.

Macrostate: Any state of a system that describes the actual observation of its microscopic statistical properties.

Phase Space: Phase space refers to the plotting of both particle momentum & position on a two-dimensional graph.

If we want to define the state of a particle in phase space coordinate system, then we require position (3D) & momentum (3D). Thus phase space is a six-dimensional coordinate consisting of 3 dimensions of position and 3 dimensions of momentum.

Position Space: Three-dimensional space in which position (location) of a particle is specified. Eg. small volume element ( $dV$ ) in position space is given by -  $dV = dx dy dz$

Momentum Space: Three-dimensional space in which momentum of particle is specified by momentum coordinates  $P_x, P_y, P_z$ .

Cells in Phase Space: - cell in phase space corresponding to a state of a system, can be labeled with some number which.  
A cell is a sub-compartment. Each compartment is divided into large number of cells, in such a way that each cell is of same size.

Has Molecules! Identical particles of any spins & can be distinguished.

Bosons! - Identical particles of 0 or integral spins & can't be distinguished.

Fermions! - Identical particles of odd half integral spins & can't be distinguished.

## Failure of Classical Statistics (M-B Statistics)

classical statistics fails to explain -

- I) observed energy distribution of Electron in a metal.
- II) observed Energy distribution of photons inside an Enclosure.

That's why we need Quantum Statistics

## Classical Statistics (Maxwell - Boltzmann Statistics)

"Maxwell Boltzmann law gives distribution of total Energy of a system having identical but distinguishable particles." It is mostly used in kinetic theory of gas to find various aspects such as molecular velocity, probable speed of molecule of gas etc.

### Assumption of M-B Statistics.

- I) All the particles which obey M-B statistics are known as Maxons & are identical & distinguishable
- II) chosen volume has large number of cells and volume of each phase space is extremely small.
- III) There is no limit on number of particles in a cell.
- IV) Total number of particles and total energy remains constant in the system.
- V) The state of each particle is specified by position and momentum coordinates.
- VI) Energy is continuous.

probability  $f(E)$  that a boson occupies a state of energy  $E$  is given by -

$$f(E) = A \exp^{-E/KT}$$

## Maxwell Boltzmann Distribution Law:-

"Maxwell Boltzmann distribution law gives how the total energy of a system is distributed among various member of a system having identical distinguishable particle"

Let us consider an assembly of  $N$  molecule in phase space. Let us divide phase space into number of cells  $1, 2, \dots, k$ , which contain number of particle  $n_1, n_2, \dots, n_k$  respectively in equilibrium. Also suppose energy of  $N$  molecule are limited to  $E_1, E_2, \dots, E_k$  within sequence of energy level.

Note (  $N$  - Molecule  
 $1, 2, 3, \dots, k$  - Number of cells.  
 $n_1, n_2, \dots, n_k$  - Number of particle in cell )

As the molecule are moving continuously in many different ways and obey fundamental postulates of Statistical Mechanics as -

① Total Number of Molecule is constant

$$N = n_1 + n_2 + \dots + n_k = \text{const.}$$

$$\delta N = \delta n_1 + \delta n_2 + \dots + \delta n_k = 0$$

$$\sum \delta n_i = 0 \quad \text{--- ①}$$

② Total Energy of System is constant.

$$E = E_1 n_1 + E_2 n_2 + \dots + E_k n_k = \text{const.}$$

# ~~Maxwell-Boltzmann Statistics~~

$$\delta E = E_1 \delta n_1 + E_2 \delta n_2 \dots \delta E_k \delta n_k = 0$$

$$\sum_i E_i \delta n_i = 0 \quad \text{--- (2)}$$

(3) When a gas is in Equilibrium probability is Maximum.

$$\delta P = 0$$

$$\delta \log P = 0 \quad \text{--- (3)}$$

Now considering

$j_i$  - Number of Energy level each having Energy  $E_i$

$n_i$  - Number of particles having Energy  $E_i$

Then Probability  $P$  for distribution is given by -

$$P = \frac{N!}{n_1! n_2! n_3! \dots n_k!} (g_1)^{n_1} (g_2)^{n_2} \dots (g_k)^{n_k}$$

$$P = \frac{N!}{n_1! n_2! n_3! \dots n_k!} (g_1)^{n_1} (g_2)^{n_2} \dots (g_k)^{n_k}$$

$$P = \frac{\sum_{i=1}^k (g_i)^{n_i} N!}{\sum_{i=1}^k n_i!}$$

Taking log on both side

$$\log P = \log N! + \sum_i \log (g_i)^{n_i} - \sum_i \log n_i!$$

$$\log P = \log N! + \sum_i n_i \log g_i - \sum_i \log n_i!$$

(Using Stirling Approximation  
 $\log N! = N \log N - N$ )

$$\Rightarrow \log P = N \log N - N + \sum_i n_i \log g_i - \sum_i n_i \log n_i - \sum_i n_i$$

(Using  $\sum_i n_i = N$  (total number of particle)  
 $\sum_i n_i^2 = N$ )

$$= N \log N - \cancel{N} + \sum_i n_i \log g_i - \sum_i n_i \log n_i + \cancel{N}$$

$$\log P = N \log N - \sum_i n_i \log n_i + \sum_i n_i \log g_i$$

(Now Differentiating both side  
 and considering  $N$  &  $g_i$  as constant)

$$\delta \log P = 0 - \left( \sum_i n_i \delta(\log n_i) + \sum_i \log n_i \delta n_i \right) + \sum_i \log g_i \delta n_i$$

$$\delta \log P = 0 - \sum_i \frac{n_i}{g_i} \delta n_i - \sum_i \log n_i \delta n_i + \sum_i \log g_i \delta n_i = 0$$

$$= 0 - 0 - \sum_i \log n_i \delta n_i + \sum_i \log g_i \delta n_i$$

(Using  $\sum_i \delta n_i = \text{Total Number of molecules} = 0$ )

$$= - \left( \sum_i \log \frac{n_i}{g_i} \delta n_i \right)$$

(Now substituting  $\delta \log P = \text{Total probability} = 0$ )

$$\Rightarrow \sum_i \log \frac{n_i}{g_i} \delta n_i = 0 \quad \text{--- (4)}$$

Multiplying (1) by  $-\alpha$  (2) by  $-\beta$  and adding (1), (2), (4)

$$\sum_i \log \frac{n_i}{g_i} \delta n_i + \sum_i \alpha \delta n_i + \sum_i \beta \epsilon_i \delta n_i = 0$$

$$\sum_i \left( \log \frac{n_i}{g_i} + \alpha + \beta \epsilon_i \right) \delta n_i = 0$$

$$\Rightarrow \log \frac{n_i}{g_i} + \alpha + \beta \epsilon_i = 0$$

$$\log \frac{n_i}{g_i} = -\alpha - \beta \epsilon_i$$

$$\frac{n_i}{g_i} = e^{-\alpha - \beta \epsilon_i}$$

$$n_i = \frac{g_i}{e^{\alpha + \beta \epsilon_i}} = g_i e^{-\alpha} e^{-\beta \epsilon_i}$$

Letting  $A = e^{-\alpha}$

$$\boxed{n_i = g_i A e^{-\beta \epsilon_i}}$$

## Bose-Einstein Statistics:-

"Bose-Einstein Statistics (B.E Statistics) determine the statistical distribution of identical indistinguishable particles (Boson) over the Energy state"

### BASIC Postulates:-

- I) All the particles which obey B.E statistics are known as Bosons. Bosons are identical & indistinguishable particles.
- II) Bosons obey Uncertainty Principle
- III) Size of cell can't be less than  $h^3$ , where  $h$  is the Planck's constant
- IV) The wave function representing boson do overlap slightly and are symmetric.
- V) The bosons don't obey Pauli Exclusion Principle
- VI) Energy states are discrete.
- VII) Number of phase cells is comparable with number of particles.
- VIII) Bose-Einstein statistics is applicable to number of particle with integral spin angular momentum in unit of  $(\frac{h}{2\pi})$

Let us consider a system consisting of  $N$  independent and indistinguishable particles. Let there be  $n_1, n_2, n_3, \dots, n_k$  particles, having mean energy  $E_1, E_2, \dots, E_k$  respectively, in compartment numbered as  $1, 2, 3, \dots, k$ , containing  $g_1, g_2, \dots, g_i, \dots, g_k$  cells respectively.

Note  $\left\{ \begin{array}{l} N - \text{Total Number of particles} \\ n_1, n_2, n_3, \dots, n_k - \text{Number of particles in a cell} \\ g_1, g_2, g_3, \dots, g_i, \dots, g_k - \text{cells} \end{array} \right\}$

Now these particles obey fundamental postulates of quantum statistics as-

① Total Number of particles is constant

$$N = n_1 + n_2 + \dots + n_k = \text{constant}$$

$$\delta N = \delta n_1 + \delta n_2 + \dots + \delta n_k = \text{constant} = 0$$

$$\Rightarrow \boxed{\sum \delta n_i = 0} \quad \text{--- (1)}$$

② Total Energy of system is constant

$$E = E_1 n_1 + E_2 n_2 + \dots + E_i n_i + \dots + E_k n_k = \text{constant}$$

$$\delta E = E_1 \delta n_1 + E_2 \delta n_2 + \dots + E_i \delta n_i + \dots + E_k \delta n_k = 0$$

$$\Rightarrow \boxed{\sum E_i \delta n_i = 0} \quad \text{--- (2)}$$

③ Distribution of particles is constant. (Thermodynamic Probability)

$$\frac{1}{W} \frac{\delta W}{\delta n_i} = 0 \quad \text{--- (3)}$$

Hence the Number of ways in which  $n_i$  Particles are to be distributed in  $g_i$  cells in  $i$ th compartment is given by -

$$w_i = \frac{(n_i + g_i - 1)!}{n_i! (g_i - 1)!}$$

Total Number of Arrangements is -

$$W = \sum_{i=1}^K \frac{(n_i + g_i - 1)!}{n_i! (g_i - 1)!} = \sum_{i=1}^K \frac{(n_i + g_i)!}{n_i! g_i!}$$

Taking log on both sides.

$$\log W = \sum_{i=1}^K (\log (n_i + g_i)! - \log n_i! - \log g_i!)$$

( Using Stirling Approximation.  
 $\log N! = N \log N - N$  )

$$\log W \approx \sum_{i=1}^K ( (n_i + g_i) \log (n_i + g_i) - (n_i + g_i) - n_i \log n_i - g_i \log g_i + g_i )$$

$$\log W \approx \sum_{i=1}^K ( n_i \log (n_i + g_i) + g_i \log (n_i + g_i) - n_i \log n_i - g_i \log g_i )$$

Now <sup>Differentiating</sup> ~~taking log~~ on both side & considering  $g_i$  as constant

$$\delta(\log W) \approx \sum_{i=1}^K ( n_i \delta \log (n_i + g_i) + \log (n_i + g_i) \delta n_i + g_i \delta \log (n_i + g_i) - (n_i \delta \log n_i + \log n_i \delta n_i) + 0 )$$

$$\Rightarrow 0 = \sum_{i=1}^K ( n_i \delta \log (n_i + g_i) + g_i \delta \log (n_i + g_i) + \log (n_i + g_i) \delta n_i - n_i \delta \log n_i + \log n_i \delta n_i )$$

$$\Rightarrow 0 = \sum_{i=1}^k \left( (n_i + g_i) \int \log(n_i + g_i) + \log(n_i + g_i) S_{ni} - n_i \int \log n_i - \log n_i S_{ni} \right)$$

$$\Rightarrow 0 = \sum_{i=1}^k \left( \cancel{(n_i + g_i)} \times \frac{1}{\cancel{(n_i + g_i)}} S_{ni} + S_{ni} \log(n_i + g_i) - \cancel{n_i} \times \frac{1}{\cancel{n_i}} S_{ni} - S_{ni} \log n_i \right)$$

$$\Rightarrow 0 = \sum_{i=1}^k \left( S_{ni} \log(n_i + g_i) - S_{ni} \log n_i \right)$$

$$\Rightarrow - \sum_{i=1}^k \left( \log n_i - \log(n_i + g_i) \right) S_{ni} = 0$$

$$\Rightarrow \sum_{i=1}^k \log \left( \frac{n_i}{n_i + g_i} \right) S_{ni} = 0 \quad \text{--- (4)}$$

Now multiplying (1) by  $\alpha$  & (2) by  $\beta$   
Then adding (1), (2) & (4)

$$\Rightarrow \sum_{i=1}^k \left( \log \left( \frac{n_i}{n_i + g_i} \right) + \alpha S_{ni} + \beta E_i S_{ni} \right) = 0$$

$$\Rightarrow \log \left( \frac{n_i}{n_i + g_i} \right) + \alpha + \beta E_i = 0$$

$$\log \left( \frac{n_i}{n_i + g_i} \right) = -\alpha - \beta E_i$$

$$\left( \frac{n_i}{n_i + g_i} \right) = e^{-\alpha - \beta E_i} = e^{-(\alpha + \beta E_i)}$$

$$\frac{n_i + g_i}{n_i} = e^{(\alpha + \beta E_i)}$$

$$1 + \frac{g_i}{n_i} = e^{(\alpha + \beta E_i)}$$

$$\Rightarrow \boxed{n_i = \frac{g_i}{e^{\alpha + \beta E_i} - 1}}$$

## Fermi Dirac Statistics:

Fermi Dirac describe distribution of particles in a system. (identical indistinguishable particles) that obey Pauli Exclusion Principle.

### Basic Postulates:-

- (i) All the particles which obey Fermi-Dirac statistics known as fermions.
- ii) Fermions obey Uncertainty Principle.
- iii) Fermions obey Pauli Exclusion Principle.
- iv) Fermi Dirac distribution law is represented by -

$$n_i = \frac{g_i}{e^{\alpha} e^{E_i/kT} + 1}$$

Let us consider a system of  $n$  independent and indistinguishable particles. Let the compartment be marked  $1, 2, \dots, k$  and their mean energy values be represented by  $E_1, E_2, \dots, E_k$ . It contains  $g_1, g_2, \dots, g_k$  cells. Total number of particles in the system is  $n = n_1 + n_2 + \dots + n_k$ .

Now these particles obey fundamental postulates of quantum statistics. as -

① Total number of particles is constant

$$N = n_1 + n_2 + \dots + n_k = \text{constant}$$

$$\delta N = \delta n_1 + \delta n_2 + \dots + \delta n_k = \text{constant} = 0$$

$$\Rightarrow \left| \sum \delta n_i = 0 \right| \quad \text{--- (1)}$$

② Total Energy of system is constant

$$E = E_1 n_1 + E_2 n_2 + \dots + E_k n_k = \text{constant}$$

$$\delta E = E_1 \delta n_1 + E_2 \delta n_2 + \dots + E_k \delta n_k = 0$$

$$\Rightarrow \left| \sum E_i \delta n_i = 0 \right| \quad \text{--- (2)}$$

③ Distribution of particle (Thermodynamic Probability) is constant

$$\text{i.e. } \frac{\delta W_2}{\int \log W_2} \quad \text{--- (3)}$$

∴ Hence the Number of ways in which  $n_i$  particles, can be distributed in  $g_i$  cells in  $i$ th compartment is -

$$W = \sum_{i=1}^k \frac{g_i!}{n_i! (g_i - n_i)!}$$

Taking log on both sides -

$$\log W = \sum_{i=1}^k \left( \log g_i! - \log n_i! - \log (g_i - n_i)! \right)$$

( Using Stirling Approximation,  $\log N! = N \log N - N$  )

$$\log W = \sum_{i=1}^k \left( g_i \log g_i - g_i - (n_i \log n_i - n_i) - ( (g_i - n_i) \log (g_i - n_i) - (g_i - n_i) ) \right)$$

$$\Rightarrow \log W = \sum_{i=1}^k \left( g_i \log g_i - g_i - n_i \log n_i + n_i - (g_i - n_i) \log (g_i - n_i) + g_i - n_i \right)$$

$$\Rightarrow \log W = \sum_{i=1}^k \left( g_i \log g_i - n_i \log n_i - (g_i - n_i) \log (g_i - n_i) \right)$$

( Now Differentiating both sides and considering  $g_i$  as constant as  $g_i$  is Number of cells which is constant )

$$\Rightarrow \frac{d}{dn_i} \log W = \sum_{i=1}^k \left( 0 - (n_i \log n_i + \log n_i \cdot n_i) - ( (g_i - n_i) \log (g_i - n_i) + \log (g_i - n_i) \cdot (-n_i) ) \right)$$

$$0 = \sum_{i=1}^k \left( 0 - \left( n_i \times \frac{1}{n_i} \cdot n_i + \log n_i \cdot n_i \right) - \left( (g_i - n_i) \times \frac{1}{(g_i - n_i)} \cdot (-n_i) + \log (g_i - n_i) \cdot (-n_i) \right) \right)$$

$$0 = \sum_{i=1}^k \left( 0 - 0 - \log n_i \cdot n_i - 0 + \log (g_i - n_i) \cdot n_i \right)$$

$$0 = \sum_{i=1}^k \left( -\log (g_i - n_i) \cdot n_i + \log n_i \cdot n_i \right) \quad \text{--- (4)}$$

∴ Multiplying (1) by  $\alpha$  (2) by  $\beta$  and Adding (1) + (2) + (4)

$$\Rightarrow \sum_{i=1}^k \left( \log \left( \frac{n_i}{g_i - n_i} \right) + \alpha \delta n_i + \beta \delta E_i \right) = 0$$

$$\Rightarrow \sum_{i=1}^k \left( \log \left( \frac{n_i}{g_i - n_i} \right) + \alpha \delta n_i + \beta \delta E_i \right) \delta n_i = 0$$

$$\Rightarrow \log \frac{n_i}{g_i - n_i} = -\alpha - \beta E_i$$

$$\frac{n_i}{g_i - n_i} = e^{-\alpha - \beta E_i} = e^{-(\alpha + \beta E_i)}$$

$$\frac{g_i - n_i}{n_i} = e^{\alpha + \beta E_i}$$

$$\frac{g_i}{n_i} - 1 = e^{\alpha + \beta E_i}$$

$$\frac{g_i}{n_i} = e^{\alpha + \beta E_i} + 1$$

$$n_i = \frac{g_i}{e^{\alpha + \beta E_i} + 1}$$

## Applications of Quantum Statistics:-

i) Maxwell Boltzmann (M-B) Statistics:-

Application:- Molecular speed and Energies of an ideal gas.

ii) Bose-Einstein (B-E) Statistics:-

Application - Black body spectrum.

iii) Fermi-Dirac (F-D) Statistics).

Application:- Free Electron Theory.

## Application of Maxwell's Boltzmann Statistics:-

### (a) Molecular Speed and Energy of an ideal gas.

Let us consider an ideal gas in a container of volume  $V$   
then according to Maxwell Boltzmann Distribution law

The number of molecule in the  $i$ th cell of Energy  $E_i$  will be

$$n_i = g_i e^{-(\alpha + \beta E_i)}$$

$$n_i = \frac{g_i}{e^{(\alpha + \beta E_i)}}$$

Substitute  $\beta = \frac{1}{k_B T}$

$$\Rightarrow n_i = \frac{g_i}{e^{(\alpha + E_i/k_B T)}} = g_i e^{-\alpha} e^{-E_i/k_B T}$$

Considering continuous Distribution of Molecular energy

$$n(E) d(E) = g(E) e^{-\alpha} e^{-(E/k_B T)} dE \quad \text{--- (1)}$$

Now a molecule of Energy  $E$  has momentum  $p$   
whose magnitude is

$$p = \sqrt{2mE}$$

Let us consider a volume of spherical cell in a momentum space of radius  $p$  and thickness  $dp$ . Volume of this spherical cell is

$$4\pi p^2 dp$$

Now  $g(p) dp =$  Number of cell in momentum range  $p$  to  $p+dp$

$$= 4\pi p^2 dp$$

$$= B p^2 dp$$

where  $B$  is some constant.

Since each magnitude of momentum  $p$  correspond to Energy  $E$

Now  $g(E) d(E) =$  Number of Energy State

$$= g(p) dp$$

$$= B p^2 dp$$

But  $p^2 = 2mE$  and  $dp = \frac{2m dE}{\sqrt{2mE}}$

we have

$$\begin{aligned} g(E) d(E) &= B p^2 dp \\ &= B \times \frac{2mE}{\sqrt{2mE}} \times \frac{2m dE}{\sqrt{2mE}} \\ &= B \times \sqrt{2mE} \times 2m dE \\ &= B (2m)^{3/2} \times E^{1/2} dE \end{aligned} \quad (2)$$

Substituting (2) in (1)

From (1)  $n(E) d(E) = g(E) e^{-\alpha} e^{-(E/k_B T)} dE$

$$= B (2m)^{3/2} E^{1/2} \times e^{-\alpha} e^{-(E/k_B T)} dE$$

Let  $B (2m)^{3/2} e^{-\alpha} = \text{constant} = C$

$$\Rightarrow n(E) d(E) = C \sqrt{E} e^{-(E/k_B T)} dE \quad (3)$$

Now to find C make the condition that total No. of molecule is N.

$$N = \int_0^{\infty} n(E) dE = C \int_0^{\infty} \sqrt{E} e^{-(E/k_B T)} dE$$

Using  $\int_0^{\infty} \sqrt{x} e^{-ax} dx = \frac{1}{2a} \sqrt{\frac{\pi}{a}}$

$$\Rightarrow N = C \frac{1}{2(1/k_B T)} \sqrt{\frac{\pi}{(1/k_B T)}} = C \times \frac{\sqrt{\pi}}{2} \times (k_B T)^{3/2}$$

$$\Rightarrow C = \frac{2\pi N}{(\pi k_B T)^{3/2}}$$

Put value of C in (3)

$$\Rightarrow n(E) d(E) = \frac{2\pi N}{(\pi k_B T)^{3/2}} \times \sqrt{E} e^{-(E/k_B T)} dE \quad (4)$$

\* Now total Energy of the system is -

$$E = \int_0^{\infty} E n(E) dE = \frac{2\pi N}{(\pi k_B T)^{3/2}} \int_0^{\infty} E^{3/2} e^{-E/k_B T} dE$$
$$= \frac{2\pi N}{(\pi k_B T)^{3/2}} \times \frac{3}{4} (k_B T)^2 \sqrt{\pi k_B T}$$

( Using  $\int_0^{\infty} \sqrt{x} e^{-ax} dx = \frac{1}{2a} \sqrt{\frac{\pi}{a}}$  )

$$E = \frac{3}{2} k_B T$$

b) Maxwell Boltzmann velocity distribution.

Now Number of Molecule with velocity  $v$  &  $v + dv$  is,

Substituting  $E = \frac{1}{2} m v^2$ ,  $dE = m v dv$ . (4)

We get

$$n(v) dv = \frac{(\sqrt{2} \pi N m^{3/2})}{(\pi k_B T)^{3/2}} v^2 e^{-mv^2/2k_B T} dv$$

## Some Molecular Speeds!:-

- 1) Mean or Average Speed!:- Average speed of all the molecules.

$$V_{avg} = \sqrt{\frac{8KT}{\pi m}}$$

- 2) Root Mean Square Speed!:- Mean speed of square of all the molecule.

$$V_{rms} = \sqrt{\frac{3KT}{m}}$$

- 3) Most Probable Speed!:- Speed possessed by maximum number of molecule.

$$V_{mp} = \sqrt{\frac{2KT}{m}}$$

Relation between  $V_{rms}$ ,  $V_{mp}$  &  $V_{avg}$ .

$$\begin{aligned} V_{rms} : V_{avg} : V_{mp} &= \sqrt{\frac{3KT}{m}} : \sqrt{\frac{8KT}{\pi m}} : \sqrt{\frac{2KT}{m}} \\ &= 1.73 \sqrt{\frac{KT}{m}} : 1.59 \sqrt{\frac{KT}{m}} : 1.414 \sqrt{\frac{KT}{m}} \\ &= 1 : 0.92 : 0.817 \end{aligned}$$

$$V_{rms} > V_{avg} > V_{mp}$$

## Applications of Fermi Dirac Statistics:-

### Free Electron Theory:-

It is of 2 type

- I) Classical (Lorentz-Drude Theory)
- II) Quantum (Sommerfeld Theory).

### Classical Free Electron Theory:

According to classical free electron theory electrons are assumed completely free. A single free electron is surrounded by large number of free electron and positive ion core.

#### ASSUMPTIONS:

- \* In this theory it is assumed that when a free electron moves in a solid its potential energy is zero. i.e. electron can move freely throughout the crystal
- \* Electron can move freely inside the solid, but it can't leave the surface. (outside no force on electron but inside a constant force on electron). This is an infinite barrier so no electron can cross the barrier



Features:- Free electron gas obey Maxwell Boltzmann statistics.

- i) It successfully defines Ohm's law.
- ii) These free electrons can move freely so metals have high electrical & thermal conductivity.

## Failure of Classical Free electron theory.

- i) It can't explain temperature dependence of Electrical Resistivity.
- ii) It fails to explain Electronic heat capacity and Paramagnetic susceptibility of Free electron.

## Quantum Free Electron theory:

Classical theory assumed that electron in Free Electron gas follow M-B Statistics. but in Quantum Free Electron theory electron follow Fermi Dirac Statistics as Pauli Exclusion Principle applicable on Electron.

→ All Assumption of Quantum Free Electron theory are same as classical Free Electron theory except now Electron follow Fermi Dirac Statistics. (Pauli Exclusion Principle).

Note:- (Pauli Exclusion Principle:- No two Electron can have same Quantum State).

Assumptions -

- i) Energy of Free electron are quantized
- ii) Free  $e^-$  obey Pauli Exclusion Principle
- iii) Distribution of Free  $e^-$  obey Fermi-Dirac quantum Statistics.
- iv) Free  $e^-$  have same potential Energy everywhere within metal
- v) Electrons are treated free

## Energy Distribution among free electron!

We have to find distribution of Electron Energy in a metal

Number of Electron with Energy  $E$  is given by -

$$n(E) d(E) = \frac{g(E) d(E)}{e^{\lambda} e^{E/KT} + 1} \quad \text{--- (1)}$$

Electron Energy State is given by -

$$g(E) d(E) = \frac{8\sqrt{2} \pi V m^{3/2}}{h^3} \frac{E^{1/2} dE}{e^{E/KT}}$$

Substituting  $g(E) d(E)$  in (1)

$$n(E) d(E) = \frac{8\sqrt{2} \pi V m^{3/2}}{h^3} \frac{E^{1/2} dE}{e^{\lambda} e^{E/KT} + 1}$$

$$\lambda = \frac{E_F}{KT}$$

$$n(E) d(E) = \frac{8\sqrt{2} \pi V m^{3/2}}{h^3} \frac{E^{1/2} dE}{e^{(E-E_F)/KT} + 1}$$

Fermi Energy: Fermi Energy is the Energy difference between highest & lowest occupied state.

Suppose a metal contains  $N$  free Electrons

Total Number of Free Electrons is equal to the total number of quantum states upto Energy  $E_F$ .

$$N = \int_0^{E_F} g(E) dE$$

$$N = \frac{8\sqrt{2}\pi V m^{3/2}}{h^3} \int_0^{E_F} E^{1/2} dE$$

$$= \frac{16\sqrt{2}\pi V m^{3/2}}{3h^3} E_F^{3/2}$$

$$= \frac{8\pi V}{3} \left(\frac{2m}{h^2}\right)^{3/2} E_F^{3/2}$$

$$E_F = \frac{h^2}{2m} \left(\frac{3N}{8\pi V}\right)^{2/3} \quad \text{--- Fermi Energy}$$

using  
( $k = \frac{h}{2\pi}$ )

OR

$$E_F = \frac{h^2}{2m} \left(\frac{3\pi^2 N}{V}\right)^{2/3}$$

## Electronic Specific Heat:-

Electronic Specific Heat of Electron is defined as the Ratio of portion of Heat used by Electron to Rise the temperature of the System.

we have

Internal Energy Per mole of Metal is.

$$E_e = \frac{3}{2} NkT = \frac{3}{2} RT \quad \left( k = \frac{R}{N} \right)$$

R - Gas Constant. ~~N = Avogadro Number.~~

N - Avogadro Number.

Then molar Specific Heat of Electron should be.

$$C_{ve} = \left( \frac{dE_e}{dT} \right)_v = \frac{3}{2} R$$

Total Specific Heat of Metal should be

$$C_v = 3R + \frac{3}{2} R = \frac{9}{2} R$$

$$\left| C_{ve} = \frac{3}{2} R \right|$$

- Electronic Specific Heat.

## Fermi Distribution Function:

$$f(E) = \frac{1}{e^{\left(\frac{E-E_F}{kT}\right)} + 1}$$

This describes the probability of occupancy of  $N e^-$ .  
OR

At temperature  $T$  the probability of occupation of an electron state of energy  $E$  is given by Fermi Distribution Function.

### Behaviour at Different temperature:

i)  $T=0K$ ,  $E < E_F$

$$f(E) = \frac{1}{e^{-\infty} + 1} = \frac{1}{e^{-\infty} + 1} = 1$$

$$\Rightarrow \boxed{f(E) = 1}$$

Probability of state to be filled is 1 i.e. state is 100% filled.

ii)  $T=0K$ ,  $E > E_F$

$$f(E) = \frac{1}{e^{+\infty} + 1} = \frac{1}{e^{+\infty} + 1} = 0$$

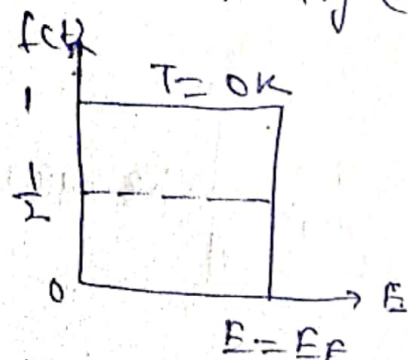
$$\Rightarrow \boxed{f(E) = 0}$$

It means no state is filled i.e. completely empty.

iii)  $T=0$ ,  $E = E_F$

$$\boxed{f(E) = \frac{1}{2}}$$

Fermi level is the level of 50% occupancy. Fermi level lying between completely filled and completely empty level.



Q. Use Fermi distribution function to obtain value  $F(E)$  for  $(E - E_F) = 0.01 \text{ eV}$  at  $200 \text{ K}$ .

Sol. Difference between Energy level  $(E - E_F) = 0.01 \text{ eV}$   
 $= 0.01 \times 1.6 \times 10^{-19} \text{ J} = 1.6 \times 10^{-21} \text{ J}$

We have Fermi Distribution function -

$$F(E) = \frac{1}{1 + e^{(E - E_F)/kT}}$$
$$= \frac{1}{1 + e^{(1.6 \times 10^{-21}) / (1.38 \times 10^{-23} \times 200)}}$$
$$= \frac{1}{1 + e^{0.5797}} = 0.3589$$

Q. Calculate Fermi Energy & Fermi temperature in a metal. Fermi velocity of Electron in Metal is  $0.86 \times 10^6 \text{ ms}^{-1}$

Sol. Given velocity of electron  $v = 0.86 \times 10^6 \text{ ms}^{-1}$   
Mass of Electron  $m = 9.1 \times 10^{-31} \text{ kg}$   
Electronic charge  $e = 1.6 \times 10^{-19} \text{ C}$   
Boltzmann's constant  $k = 1.38 \times 10^{-23} \text{ J K}^{-1}$

Now Fermi Energy  $E_F = \frac{1}{2} m v^2$

$$= \frac{1}{2} \times 9.1 \times 10^{-31} \times (0.86 \times 10^6)^2$$
$$= 3.368 \times 10^{-19} \text{ J}$$

Fermi Temperature  $T_F = \frac{E_F}{k}$

$$= \frac{3.368 \times 10^{-19}}{1.38 \times 10^{-23}} = 2.43 \times 10^4 \text{ K}$$

Fermi Temperature  $T_F$  is  $2.43 \times 10^4 \text{ K}$

## Numericals! -

Q Calculate the different number of ways of arranging 6 fermions in 6 elementary phase cells.

Sol. According to Fermi-Dirac statistics different ways of arranging  $n_i$  fermions in  $g_i$  phase cell of  $i$ th energy level is given by -

$$(W)_{F-D} = \frac{g_i!}{(g_i - n_i)! n_i!}$$

$$\text{Hence } n_i = 6 \quad \& \quad g_i = 10$$

$$= \frac{10!}{(10-6)! 6!}$$

$$= \frac{10 \times 9 \times 8 \times 7 \times 6!}{4 \times 3 \times 2 \times 1 \times 6!}$$

$$= 210$$

Q Calculate the number of different arrangement of 8 indistinguishable particles among seven cells, there being no restriction on number of particles in obeying B-E statistics.

Sol. According to B-E statistics no. of different arrangement of  $n_i$  indistinguishable particle in  $g_i$  cell is given by -

$$(W)_{B-E} = \frac{(n_i + g_i - 1)!}{n_i! (g_i - 1)!}$$

$$n_i = 8, \quad g_i = 7$$

$$\text{Hence } (W)_{B-E} = \frac{(8+7-1)!}{8! \times (7-1)!} = \frac{14!}{8! 6!}$$

$$= \frac{14 \times 13 \times 12 \times 11 \times 10 \times 9 \times 8!}{8! \times (6 \times 5 \times 4 \times 3 \times 2 \times 1)}$$

$$= 3003$$

**Example 3.1** An electron gas obeys the Maxwell-Boltzman statistics. Calculate average thermal energy (in eV) of an electron in the system at 300 K. [GGSIPU, March 2015 (2 marks)]

**Solution.** 
$$E = \frac{3}{2} k_B T = \frac{3}{2} \times 1.38 \times 10^{-23} \times 300 \text{ J} = \frac{3 \times (1.38 \times 10^{-23}) \times 300}{2 \pi (1.6 \times 10^{-19})} \text{ eV} = 0.039 \text{ eV}$$

**Example 3.2** At what temperature would one in a thousand of atom in a gas of atom hydrogen be in  $n=2$  energy level ?

**Solution.** For hydrogen 
$$n(E) = -\frac{13.6}{n^2} \text{ eV}$$

$$n(E_1) = n_1 = -13.6 \text{ eV}$$

$$n(E_2) = n_2 = -3.4 \text{ eV}$$

$$g(E) = \text{no. of states formed} = 2n^2$$

$$g(E_1) = g_1 = 2 \text{ and } g(E_2) = g_2 = 8$$

For Maxwell-Boltzmann distribution is

$$\frac{n_2}{n_1} = \frac{g_2}{g_1} \frac{e^{-E_2/k_B T}}{e^{-E_1/k_B T}} = \frac{8}{2} e^{-(E_2 - E_1)/k_B T}$$

or 
$$\frac{1}{10^3} = 4e^{-(E_2 - E_1)/k_B T} \quad \text{or} \quad e^{(E_2 - E_1)/k_B T} = \frac{10^4}{2.5} = 4000$$

Taking logarithms both the sides

$$\frac{E_2 - E_1}{k_B T} = \ln 4000$$

$$k_B T = \frac{(E_2 - E_1)}{\ln 4000} = \frac{10.2 \text{ eV}}{8.29}$$

$$T = \frac{10.2 \times 1.6 \times 10^{-19}}{8.29 \times 1.38 \times 10^{-23}} \text{ K} = 1.43 \times 10^4 \text{ K}$$

$$T = 14300 \text{ K} = 14300 - 273 = 14027^\circ \text{C}$$

## Dying Stars:-

### 1) White / Black Dwarfs.

A star like our Sun will become a white dwarf when it will exhaust its Nuclear Fuel; which then settle down to become a white dwarf.

White dwarfs have no way to keep themselves hot, Eventually such stars cool completely & become black dwarfs.

Typical Data of white Dwarfs -

Content - Mostly Helium

Density -  $10^7 \text{ g/cc}$

Mass -  $10^{33} \text{ g} = \text{Mass of Sun}$

Central Temperature =  $10^7 \text{ K} = \text{Sun's Temperature}$ .

At its central temperature Helium atoms are expected to be completely ionised, we regard the gas of electrons as ideal Fermi gas.

$$E_F = \frac{\hbar^2}{2m} \left( \frac{3N}{8\pi V} \right)^{3/2} = 20 \text{ MeV}$$

$$T_F = 10^{11} \text{ K}$$

Also No white ~~star~~ dwarf star can have mass larger than.

$$M_0 = 10^{33} \text{ g} = \text{Mass of Sun}$$

If,  $M_0 = 1.4 \times \text{Mass of Sun}$

This mass is known as Chandrasekhar limit.

## ii) Neutron Stars:-

Neutron stars are typically about 10 miles in diameter have about 1.4 time the mass of a sun and spin very rapidly. Neutron stars are fascinating because they are densest object known. Due to small size and high density, a neutron star posses a surface gravitational field of about 300,000 times that of Earth.

Magnetic field about 1,000,000,000,000 times stronger than Earth.

## iii) Black holes:-

Black holes are object so dense that not even light can escape their gravity and nothing can escape from inside a black hole.

~~The~~ Existence of two types of black holes:

i) Mass with 4-15 times of mass of sun.

ii) with mass of a typical galaxy.

This evidence comes not from seeing the black holes directly, but by observing the behaviour of stars & other material near them.

## Formulae at a Glance

### 3.1 Volume of each phase cell

$$\delta V = \delta x \delta y \delta z \delta p_x \delta p_y \delta p_z$$

$$\delta x \delta p_x \geq \hbar, \delta y \delta p_y \geq \hbar \text{ and } \delta z \delta p_z \geq \hbar$$

$$\delta V = \hbar^3 \quad \left[ \because \hbar = \frac{h}{2\pi} \right]$$

### 3.2 Maxwell-Boltzmann Statistics

(i) The distribution law  $n_i = \frac{g_i}{e^{(\alpha + E_i/k_B T)}}$

(ii) Molecular energy in an ideal gas

$$E = \frac{2\pi N}{(\pi k_B T)^{3/2}} \frac{3}{4} (k_B T)^2 \sqrt{\pi k_B T} = \frac{3}{2} N k_B T$$

and  $\bar{E} = \frac{3}{2} k_B T$

(iii) Maxwell-Boltzmann velocity distribution law

$$n(v)dv = \frac{\sqrt{2} N m^{3/2}}{(\pi k_B T)^{3/2}} v^2 e^{-\frac{mv^2}{2k_B T}} dv$$

(a) RMS velocity  $v_{rms} = \sqrt{\frac{3k_B T}{m}}$

(b) Average velocity  $= \bar{v} = \sqrt{\frac{8k_B T}{\pi m}}$

(c) Most probable velocity ( $v_p$ )  $= \sqrt{\frac{2k_B T}{m}}$

3.3 (i) Symmetric wave function  $\psi_s(1,2) = +\psi(2,1)$

(ii) Antisymmetric wave function

$$\psi_A(1,2) = -\psi(2,1)$$

### 3.4 Bose-Einstein statistics

(i) Distribution law  $n_i = \frac{g_i}{(e^{\alpha + E_i/k_B T} - 1)}$

(ii) Planck's radiation law

$$dE = E(v)dv = \frac{8\pi h v^3}{c^3} \frac{1}{(e^{hv/k_B T} - 1)} dv$$

(iii) Rayleigh-Jeans law  $E(v)dv = \frac{8\pi v^2 k_B T}{c^3} dv$

(iv) Wein's displacement law

$$\lambda_{max} T = \text{Constant}$$

(v) Bose-Einstein Condensation

(a) Critical temperature

$$T_C = \frac{h^3}{2\pi m k_B T} \left( \frac{1}{2612} \frac{n}{V} \right)^{2/3}$$

(b) Rest of particles

$$n_0 = n \left[ 1 - \Sigma \left( \frac{T}{T_C} \right)^{3/2} \right] \text{ for } T < T_C$$

### 3.5 Fermi Dirac Statistics

(i) Distribution law  $n_i = \frac{g_i}{(e^{\alpha + E_i/k_B T} + 1)}$

(ii) Fermi momentum  $p_F = \left( \frac{3N}{8\pi V} \right)^{1/3} h$

(iii) Fermi Energy  $E_F = \frac{p_F^2}{2m} = \frac{h^2}{2m} \left( \frac{3N}{8\pi V} \right)^{2/3}$

(iv) Ground state energy

$$E_0 = \frac{8\pi V}{h^3} \frac{p_F^5}{10m} = \frac{3}{5} N E_F$$

(v) Mean energy of Fermion  $\bar{E} = \frac{E_0}{N} = \frac{3}{5} E_F$

(vi) Electronic heat capacity

$$(C_v)_{El} = \frac{\pi^2}{2} N k_B \left( \frac{T}{T_F} \right) = aT$$

(vii) Sommerfield constant  $a = \frac{N k_B^2 \pi^2}{2 E_F}$

## Miscellaneous Solved Numerical Problems

**Problem 3.1** A gas has only two particles A and B Show that with the help of diagrams, how these two particles can be arranged in three series 1, 2, 3 using (i) Maxwell-Boltzmann, (ii) Bose-Einstein and (iii) Fermi-Dirac statistics.

**Solution. (i) Maxwell-Boltzmann statistics :**

The two particles are distinguishable.

There is no limit to the number of particles in any one state.

The total number of ways =  $3^2 = 9$ .

States	Possible distribution in various states								
1	A	B	-	-	A	B	AB	-	-
2	B	A	A	B	-	-	-	AB	-
3	-	-	B	A	B	A	-	-	AB

**(ii) Bose-Einstein statistics :**

If A and B are quantum particles, they are indistinguishable. Thus they have to be given the same name, say A.

There is no limit to the number of particles in any one state.

The total number of ways = 6

States	Possible distribution in various states					
1	A	A	-	AA	-	-
2	A	-	A	-	AA	-
3	-	A	A	-	-	AA

**(iii) Fermi-Dirac statistics :**

The particles are indistinguishable and not more than one particle can be in any one state.

The total number of ways = 3

States	Possible distribution in various states		
1	A	A	-
2	A	-	A
3	-	A	A

**Problem 3.2** Show that on increasing temperature, the number of atoms in excited state increases.

**Solution.**  $n(E) = g(E) \cdot f(E)$

For Maxwell-Boltzmann distribution

$$n(E_i) = \frac{g(E_i)}{e^{(\alpha + E_i/k_B T)}} = g(E_i) A e^{-E_i/k_B T}$$

Then  $n(E_1) = g(E_1) A e^{-E_1/k_B T}$

and  $n(E_2) = g(E_2) A e^{-E_2/k_B T}$

$\Rightarrow \frac{n(E_2)}{n(E_1)} = \frac{g(E_2)}{g(E_1)} e^{-(E_2 - E_1)/k_B T}$

On increasing temperature  $\Rightarrow n(E_2) > n(E_1)$

**Problem 3.3** (a) Show that at  $E = E_F$  and  $T \neq 0$  K. The occupancy is equal to 0.5. (b) Show that at  $E < E_F$  and  $T = 0$  K. The occupancy of state is perpendicular and at  $E > E_F$  and  $T = 0$  K, the occupancy of state is zero.

**Solution.** From Fermi-Dirac statistics

$$n_i = \frac{g_i}{e^{\alpha + E_i/k_B T} + 1}$$

This is written as 
$$n(E) = \frac{g(E)}{e^{(E-E_F)/k_B T} + 1} = g(E) f(E)$$

Then occupancy 
$$f(E) = \frac{1}{(e^{(E-E_F)/k_B T} + 1)}$$

(a) When  $T \neq 0$  and  $E = E_F$  then 
$$f(E) = \frac{1}{e^0 + 1} = \frac{1}{1+1} = \frac{1}{2} = 0.5$$

(b) When  $T = 0$  and  $E < E_F$  then 
$$f(E) = \frac{1}{e^{-\infty} + 1} = \frac{1}{0+1} = 1$$

**Problem 3.7** At what temperature can we expect a 10% probability that electrons in a metal will have an energy which is 1% above  $E_F$  ? The Fermi energy of the metal is 5.5 eV. [GGSIPU, May 2014 (4.5 marks)]

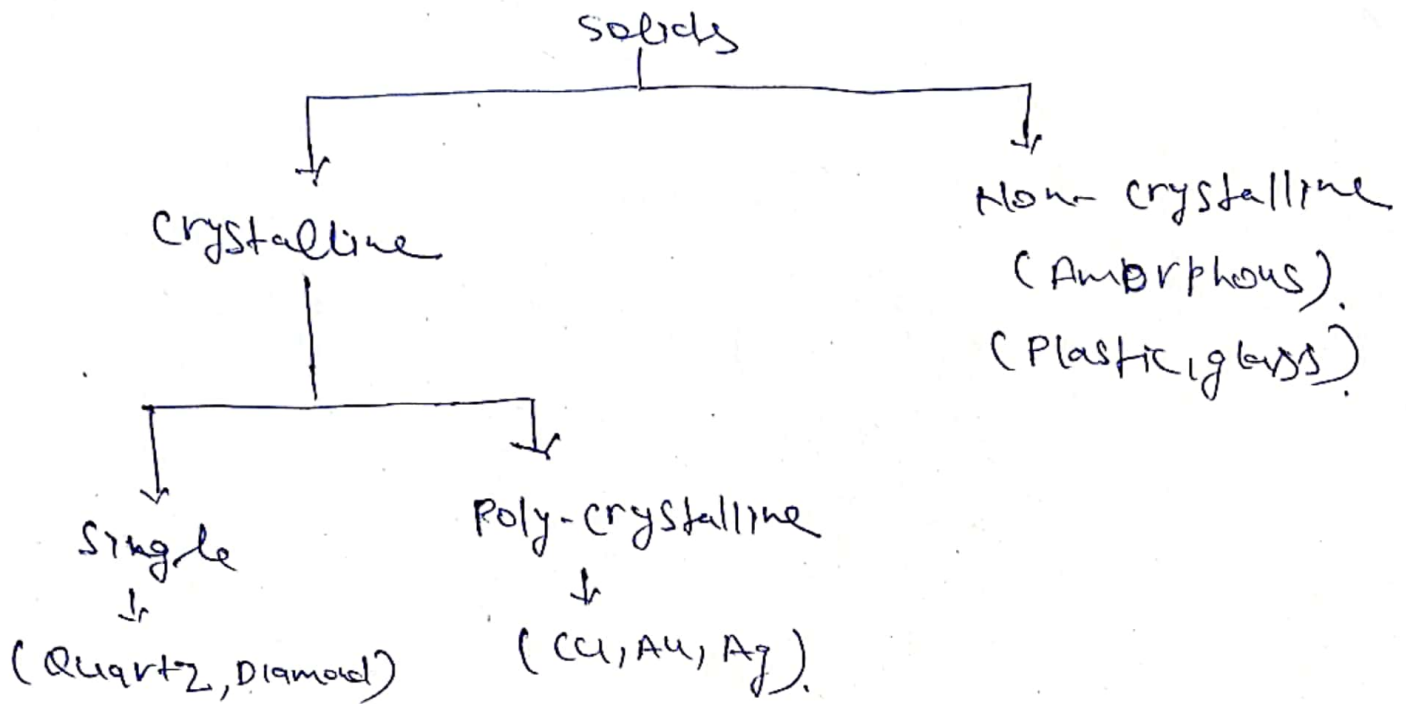
**Solution.** Given :  $f(E) = 10\%$ ,  $E = E_F + 1\%$  of  $E_F$ ,  $E_F = 5.5$  eV,  $T = ?$

$$f(E) = \frac{1}{e^{(E-E_F)/k_B T} + 1}$$

$$E = 5.5 + \frac{5.5}{100} = 5.5 + 0.555; \quad E - E_F = 0.555.$$

$$0.1 = \frac{1}{\left( \exp. \frac{0.555 \times 1.6 \times 10^{-19}}{1.38 \times 10^{-23} T} \right) + 1} = \frac{1}{\exp. \left( \frac{637.7}{T} \right) + 1} \quad \text{or} \quad T = 290.2 \text{ K}$$

## UNIT-3 CRYSTAL STRUCTURE



Crystalline Solids! - Crystalline solids are characterized by regular or periodic arrangement of atoms or molecules.

- \* Every solid is made up of atoms or molecules.
- \* Atoms are repeating themselves after certain distance.



Non Crystalline Solids! - (Amorphous)

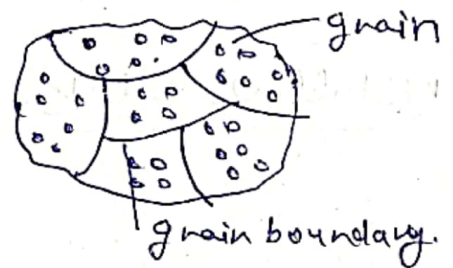
- are characterized by random arrangement of atoms or molecules.
- Arrangement of atoms will describe their properties.
- Plastic is non-crystalline.
- Copper is crystalline.
- Diamond & copper both are crystalline, but diamond is a single crystal & copper is a poly-crystalline.

Single crystal:- Any solid object in which an ordered three dimensional arrangement of atom ion or molecule is repeated throughout the entire volume.



Poly crystalline crystals:- Poly crystalline is the aggregate of large number of single crystals.

Small single crystals are known as grain. & boundaries separating these crystals are grain boundary.



Note:- In Nature most of the solids are in the form of crystalline solids because crystalline solid is more stable as compare to Amorphous solids.

## Difference between Crystalline Solids & Amorphous Solids <sup>3</sup>

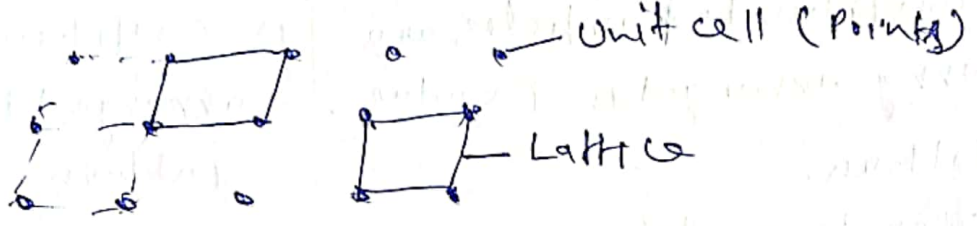
### Crystalline Solids

- I) Constituents particles are arranged in regular pattern.
- II) They have definite shape
- III) They have sharp melting point.
- IV) They have definite Enthalpy of fusion
- V) They have long order arrangement
- VI) They are more hard
- VII) They are true solids  
Cu, Fe Salt etc
- VIII) They are Anisotropic

### Amorphous Solids

- I) Constituent particles are arranged in irregular pattern.
- II) They have indefinite shape
- III) They melt over a range of temperature (in  $^{\circ}\text{C}$ )
- IV) Don't have definite Enthalpy of fusion
- V) Have short order arrangement.
- VI) They are less hard
- VII) Supercooled liquid.  
Glass, Rubber (Plastic)
- VIII) They are isotropic

Space lattice! Geometrical regular arrangements of ~~points~~ Repeated points in space



Bravais Lattice! - Three dimensional array of points such that each point is surrounded by neighbouring points in an identical way.

Note! - Lattice will be a Bravais lattice or space lattice if each lattice point has identical surroundings.

Basis! - Basis is an atom or group of atoms. By attaching the basis to each lattice point we get crystal structure

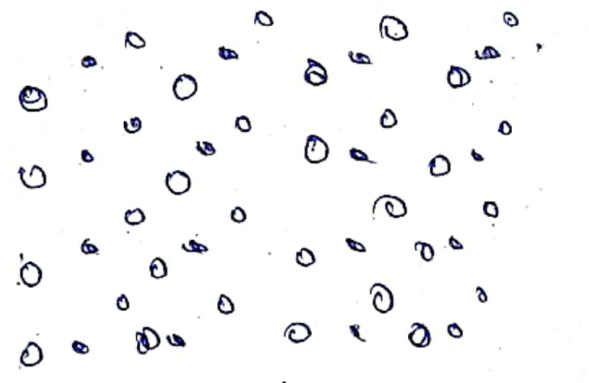
Lattice + Basis = crystal



space lattice



Basis

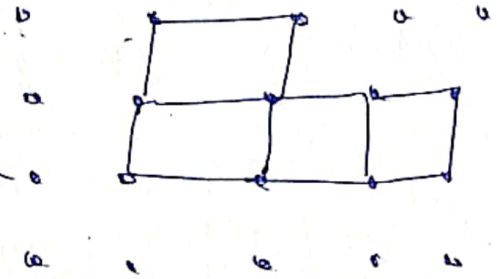


crystal structure

Unit cell:- ("Smallest Repeating Unit of the ... crystal or lattice")

Note:- Unit cell on continuous Repetition can generate whole lattice or crystal.

~~Unit cell~~ If we want to construct a crystal structure we have to arrange unit cell side by side.



Unit cells are of two type:-

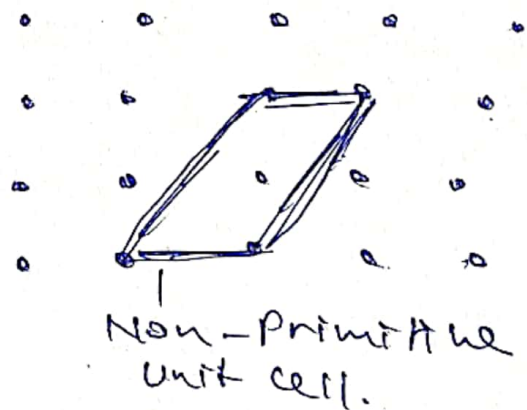
- I) Primitive
- II) Non-Primitive.

Primitive:- Lattice points belonging only to the corner of unit cells.

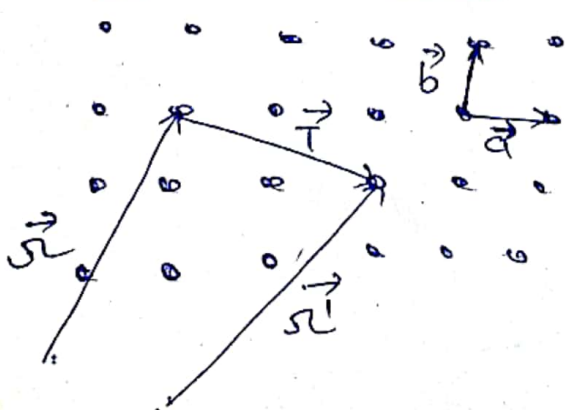
Non-Primitive:- Lattice points are belonging to the corner as well as other position of unit cell.

Note:- Effective number of lattice point for primitive cell per unit cell is 1

→ For Non-Primitive unit cell Effective number of lattice points are more than one.



Translation Vector:-



Mathematically a lattice is defined by three fundamental vector  $\vec{a}$ ,  $\vec{b}$ ,  $\vec{c}$  such that atomic arrangement look exactly identical when viewed from any two points  $\vec{r}$  &  $\vec{r}'$

$$\vec{r}' = \vec{r} + u\vec{a} + v\vec{b} + w\vec{c}$$

$u, v, w$  are integers

Types of Unit cell:- 3 types mainly.

(6)

a) Simple Cubic cell (S.C.C):- One atom is present at each corner of cube and  $\frac{1}{8}$  part is contributed by each atom inside the cube (unit cell). So total number of atom per unit cell in SCC is  $\frac{1}{8} \times 8 = 1$

b) Body centered cubic (B.C.C):

\* One atom is present at each corner of the cube contribute  $\frac{1}{8}$

\* So one atom is contributed by 8 atoms at 8 corner  $\frac{1}{8} \times 8 = 1$ .

\* One single (complete) atom is present at the center of cube contribute completely to cube.

So total number of atom =  $\frac{1}{8} \times 8 + 1 = 1 + 1 = 2$

(c) Face centered cubic (F.C.C):-

Particles or atoms are arranged at each corner as well as each faces

So contribution.

Corner

$$\frac{1}{8} \times 8$$

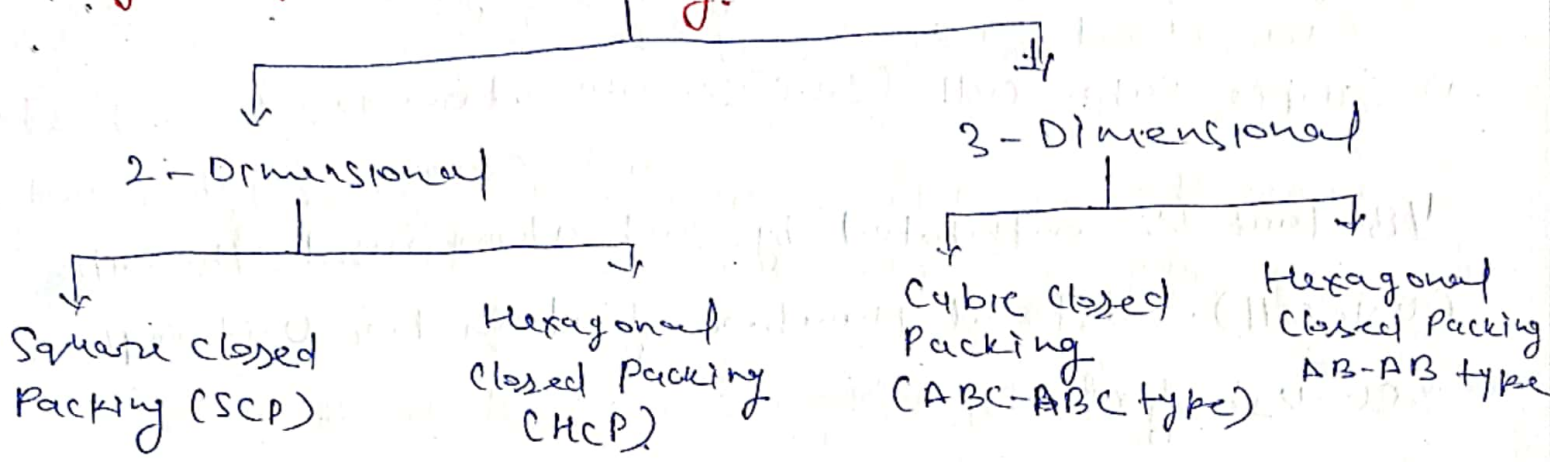
Faces

$$\frac{1}{2} \times 6$$

Total number of atom =  $3 + 1 = 4$

Co-ordination Number:- Number of Nearest touching particles in a closed packing structure.

# Types of Closed Packing:-



## I) Square closed Packing:-

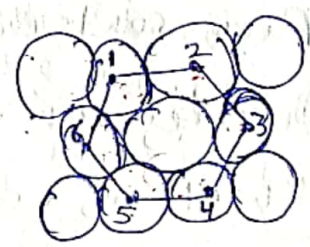
The central Atom is touched by four Neighbouring atom.



% occupied by atom in crystal = 52.4 %  
 Coordination Number = 4.

## II) Hexagonal closed Packing (2-D)

When one layer is normal and atom of second layer are put in the depression of 1st layer then it give rise to HCP in 2-Dimension.

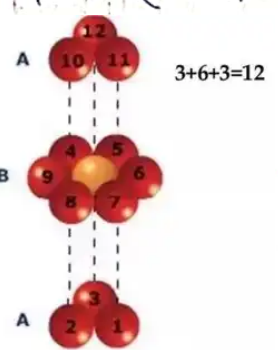


% of space occupied = 62.4 %  
 Coordination Number = 6.

## III) Hexagonal closed Packing (HCP) in 3D:-

When atom of third layer are put in the depression of tetrahedral voids then it give rise to HCP.

% of space occupied = 74 %  
 Coordination Number = 12



# Cubic Close Packing (CCP) in 3D:-

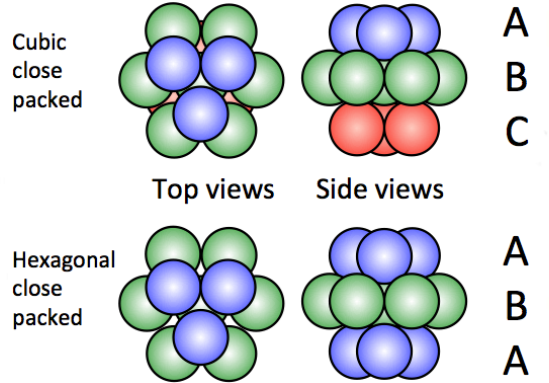
When the second layer arranged on the first layer to give rise to octahedral and tetrahedral voids in the second layer:

When atom of third layer are put in the depression of these voids then give rise to

CCP type packing

% Space occupied - 74%

Coordination Number - 12



voids or holes in crystals: These are unoccupied or blank space are left between atom when they are closely packed.

They are generally two types.

- 1) Tetrahedral void!- The void or hole which is created by four sphere in contact.
  - 2) Octahedral void!- The void or hole created by six sphere in contact.
- void or hole also called interstices.

## Packing Fraction or Efficiency of Packing:-

9

It is defined as the % of space occupied by the atom in a unit cell. It is different for different arrangement.

3)

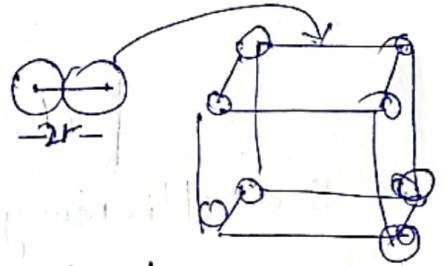
$$\% \text{ efficiency} = \frac{\text{Vol. of Sphere} \times \text{No. of Sphere}}{\text{Total Vol. of Cube}} \times 100$$

1) SCC :-

Relation between edge length (a) and radius of atom in SCC is.

$$a = 2r$$

$$\Rightarrow \boxed{r = a/2}$$



In SCC there ~~are~~ only 1 atom which contribute inside cube

$$\text{Vol} = \frac{4}{3} \pi r^3$$

$$\Rightarrow \% \text{ efficiency} = \frac{\frac{4}{3} \pi r^3 \times 1}{(2r)^3} \times 100$$

$$\boxed{\% \text{ efficiency} = 54\%}$$

11) FCC :- Relation bet edge length a and radius of atom r in FCC is

$$r = \frac{\sqrt{2}a}{4} \quad \text{or} \quad \boxed{a = \frac{4r}{\sqrt{2}}}$$

No. of sphere is = 4

$$\% \text{ Efficiency} = \frac{\frac{4}{3} \pi r^3 \times 4}{(2\sqrt{2}r)^3} \times 100$$

$$\boxed{\% \text{ Efficiency} = 74\%}$$

iii) BCC:- Relation between edge length and radius of atom (r) is

$$4r = \sqrt{3} a$$

$$\boxed{a = \frac{4r}{\sqrt{3}}}$$

No. of sphere contributing in cube in case of BCC is 2

$$\text{So \% efficiency} = \frac{\frac{4}{3} \pi r^3 \times 2}{\left(\frac{4r}{\sqrt{3}}\right)^3} \times 100$$

$$\boxed{\% \text{ efficiency} = 68\%}$$

So efficiency of packing is 68% in BCC.

Note

In BCC

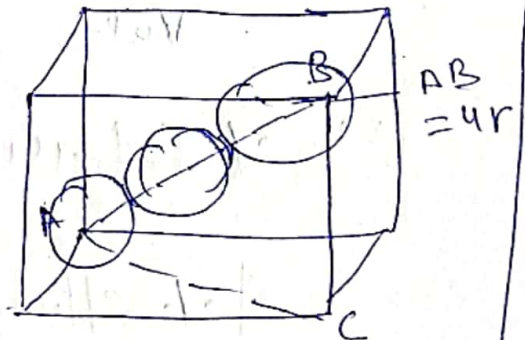
$$AC = \sqrt{2} a$$

$$AB^2 = AC^2 + BC^2$$

$$= (\sqrt{2} a)^2 + a^2$$

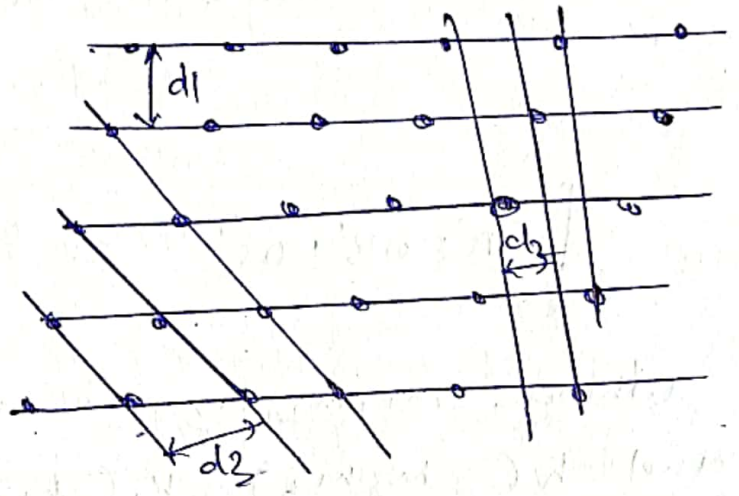
$$\boxed{AB = \sqrt{3} a}$$

$$\text{or } \boxed{4r = \sqrt{3} a}$$



Lattice Plane of a crystal :- A lattice plane may be considered as an aggregate of set of parallel equally spaced plane passing through the lattice points.

The Equally spaced planes are called lattice planes and perpendicular distance between adjacent plane is called interplanar spacing ( $d$ ).



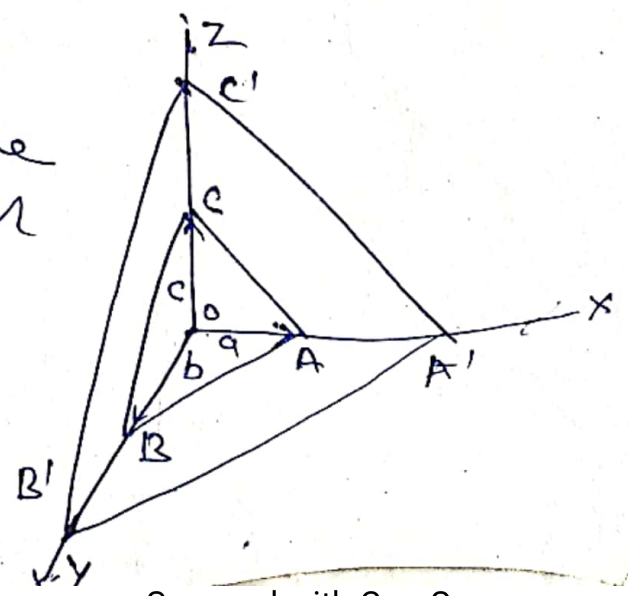
( $d_1, d_2, d_3$  are interplanar spacing)

Miller indices :- The position & orientation of lattice plane in a crystal are determined by three whole numbers, which have same ratio with one another as the reciprocals of the intercepts of the plane on the three crystal axes. These numbers are denoted as  $h, k, l$  known as Miller indices.

Let  $OX, OY, OZ$  be the three axes parallel to the crystal axes the plane  $XOY, YOZ, ZOY$  are parallel to the face of crystal

Let  $ABC$  be a standard plane cutting all three axes at points  $A, B, C$  with

$OA = a, OB = b, OC = c$   
"be the three intercept.



This law state that " the face which is parallel to the plane whose intercept on the three axes are  $m_1a$ ,  $m_2b$  &  $m_3c$  is a possible face of crystal. (Orientation)

Then

$$OA' : OB' : OC' = m_1a : m_2b : m_3c$$

$$= \frac{a}{m_2m_3} : \frac{b}{m_1m_3} : \frac{c}{m_1m_2}$$

$$OA' : OB' : OC' = \frac{a}{h} : \frac{b}{k} : \frac{c}{l}$$

where  $m_1, m_2, m_3$  are small whole number and  $h (= m_2m_3)$ ,  $k (= m_1m_3)$ ,  $l (= m_1m_2)$  Hence  $(h, k, l)$  are the Miller indices of the plane  $A', B', C'$  with respect to standard plane.

\* Note Miller indices of standard plane are always  $(h, k, l)$ .

Q. In a crystal lattice plane cut the intercept of a, 2b, 3c along three axes where a, b, c are primitive vector of unit cell. Determine Miller indices of given plane.

Sol.

$$a : 2b : 3c = \frac{a}{h} : \frac{b}{k} : \frac{c}{l}$$

$$\frac{1}{h} : \frac{1}{k} : \frac{1}{l} = 1 : 2 : 3$$

$$h : k : l = 1 : \frac{1}{2} : \frac{1}{3}$$

$$= 6 : 3 : 2$$

$$h=6, k=3, l=2$$

Q. Deduce the Miller indices of a plane which cut off intercept in the ratio 1a : 3b : -2c, along three axes, where a, b, c are primitive.

Sol.

$$1a : 3b : -2c = \frac{a}{h} : \frac{b}{k} : \frac{c}{l}$$

$$\frac{1}{h} : \frac{1}{k} : \frac{1}{l} = 1 : 3 : -2$$

$$h : k : l = 1 : \frac{1}{3} : -\frac{1}{2} = 6 : 2 : -3$$

$$h=6, k=2, l=-3$$

Q. An orthorhombic crystal has axial vector in ratio a:b:c :: 0.424 : 1 : 0.367. Find the Miller indices of those crystal plane whose Miller indices are in the ratio

- i) 0.212 : 1 : 0.183
- ii) 0.848 : 1 : 0.732

Sol.

$$i) P : q : r = \frac{a}{h} : \frac{b}{k} : \frac{c}{l}$$

$$0.212 : 1 : 0.183 = \frac{0.424}{h} : \frac{1}{k} : \frac{0.367}{l}$$

$$M.I = 2 \quad h : k : l = 2 : 1 : 2$$

(11)  $0.848 : 1 : 0.732 = \frac{0.424}{h} : \frac{1}{k} : \frac{0.367}{l}$

$$\frac{1}{h} : \frac{1}{k} : \frac{1}{l} = 2 : 1 : 2$$

Q Find the Miller indices of an orthorhombic crystal which cut intercept of  $3a, 4b$  along  $x$  &  $y$  Axes.

Sol.

$$3a : 4b : \infty = \frac{a}{h} : \frac{b}{k} : \frac{c}{l}$$

$$\frac{1}{h} : \frac{1}{k} : \frac{1}{l} = 3 : 4 : \infty$$

$$h : k : l = \frac{1}{3} : \frac{1}{4} : \frac{1}{\infty}$$

$$h : k : l = 4 : 3 : 0$$

Q Find the ratio of intercept of three axes by  $(1\bar{3}2)$  plane in a simple cubic lattice

Sol.

$P, Q, R$  be the intercept along  $x, y, z$  axes resp.

$$P : Q : R = \frac{a}{h} : \frac{b}{k} : \frac{c}{l}$$

(given  $h=1, k=-3, l=2$ )

$$P : Q : R = \frac{a}{1} : \frac{a}{-3} : \frac{a}{2} \quad (\text{For simple cubic})$$

$$P : Q : R = 6 : -2 : 3$$

Q Lattice constant for a cubic lattice is  $a$ . Deduce the spacing between  $(011)$ ,  $(101)$  &  $(112)$  plane.

Sol.

$$d = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$d(011) = \frac{a}{\sqrt{0^2 + 1^2 + 1^2}} = \left( \frac{a}{\sqrt{2}} \right)$$

## Bragg's X-Ray Diffraction:-

It is the relation between spacing of Atomic planes in crystal and angle of incidence at which these planes produce the most intense reflection.

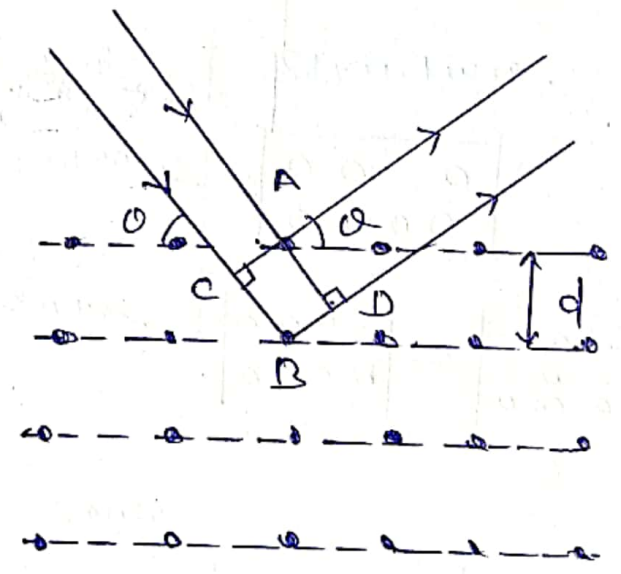
OR

When an X-Ray incident into a crystal, its angle of incidence,  $\theta$  will reflect with same angle of scattering  $\theta$ . Then the path difference  $d$  is equal to whole number  $n$  of wavelength  $\lambda$ .

OR

Diffraction beams are formed when the reflection from various plane of Atom interfere constructively.

Figure shows a particular set of Atomic planes in a crystal,  $d$  being the interplanar spacing. Suppose a monochromatic X-Ray beam is incident at an angle  $\theta$ . It is scattered by the Atom. line A & B.



The path difference for the waves reflected from adjacent planes is,

$$CB + BD = d \sin \theta + d \sin \theta$$

$$\boxed{2d \sin \theta = n \lambda}$$

↓  
Bragg's law.

Q In Bragg's spectrometer using set of Atomic plane of a crystal of interplanar spacing  $2.8 \text{ \AA}$ . The glancing angle for 1st order maxima was  $10^\circ$ . Find the wavelength of x-rays employed. At what angle would the second order maxima occur.

sol From Bragg's Eqn

$$2d \sin \theta = n\lambda$$

$$\lambda = \frac{2d \sin \theta}{n} \quad (n=1)$$

$$d = 2.8 \text{ \AA} \quad \sin \theta = \sin 10^\circ \approx 0.1736$$

$$n=1$$

$$\lambda = \frac{2 \times 2.8 \times 0.1736}{1} \approx 0.97 \text{ \AA}$$

For second order maxima  $n=2$ , Then-

$$\sin \theta = \frac{n\lambda}{2d} = \frac{2 \times 0.97 \text{ \AA}}{2 \times 2.8 \text{ \AA}}$$

$$\frac{2 \times 0.97 \text{ \AA}}{2 \times 2.8 \text{ \AA}} = 0.346$$

$$\theta = \sin^{-1}(0.346) \approx 20.3^\circ$$

(17)

Interplanar spacing! - The separation between successive lattice planes known as Interplanar Spacing.

$$d = \frac{1}{\sqrt{\frac{h^2}{a^2} + \frac{k^2}{b^2} + \frac{l^2}{c^2}}}$$

$$d = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

Q Lattice constant for a cubic lattice is  $a$  for a cubic lattice. Deduce the spacing between (111), (101) & (112) planes.

Sol:

$$d_{111} = \frac{a}{\sqrt{0^2 + 1^2 + 1^2}} = \frac{a}{\sqrt{2}}$$

$$d_{101} = \frac{a}{\sqrt{1^2 + 0^2 + 1^2}} = \frac{a}{\sqrt{2}}$$

$$d_{112} = \frac{a}{\sqrt{1^2 + 1^2 + 1^2}} = \frac{a}{\sqrt{3}}$$

Q Calculate the interplanar spacing for (321) plane in a simple cubic lattice whose lattice constant is  $4.2 \times 10^{-8}$  cm.

Sol:

$$d = \frac{a}{\sqrt{h^2 + k^2 + l^2}}$$

$$d = \frac{4.2 \times 10^{-8}}{\sqrt{3^2 + 2^2 + 1^2}}$$

$$= \frac{4.2 \times 10^{-8}}{\sqrt{14}} = \frac{4.2 \times 10^{-8}}{3.74}$$

$$= 1.12 \times 10^{-8}$$

$$d = 1.12 \text{ \AA}$$

Q2 An x-ray Analysis of cubic crystal is carried, using x-rays of wavelength  $0.58 \text{ \AA}$ . Bragg reflections are observed at glancing angle  $6.45^\circ$ ,  $9.15^\circ$  and  $11.1^\circ$ . Find interplanar spacing of crystals.

sol: Bragg's Eqn for 1st order is

$$2d \sin \theta = \lambda$$

$d_1, d_2, d_3$  be the int planar spacing at glancing angle  $6.45^\circ, 9.15^\circ, 11.1^\circ$  resp. Then

$$d_1 = \frac{\lambda}{2 \sin 6.45^\circ} = \frac{0.58}{2 \times 0.1123} = 2.58 \text{ \AA}$$

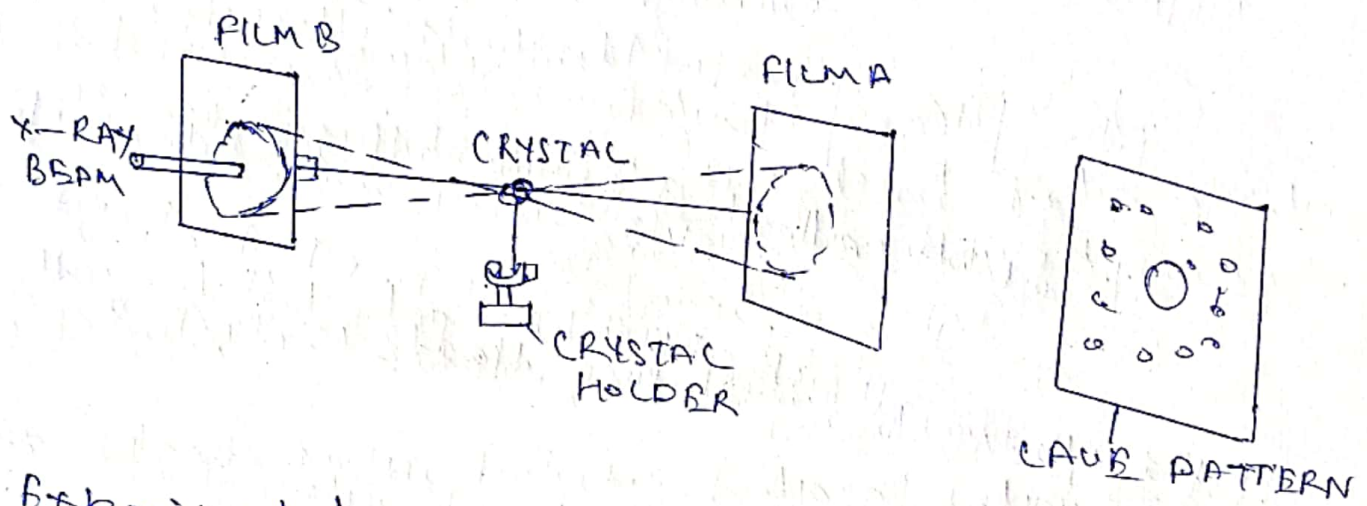
$$d_2 = \frac{\lambda}{2 \sin 9.15^\circ} = \frac{0.58}{2 \times 0.1590} = 1.82 \text{ \AA}$$

$$d_3 = \frac{\lambda}{2 \sin 11.1^\circ} = \frac{0.58}{2 \times 0.1928} = 1.51 \text{ \AA}$$

$$d_1 : d_2 : d_3 = 2.58 \text{ \AA} : 1.82 \text{ \AA} : 1.51 \text{ \AA}$$

$$= 1 : 0.705 : 0.585 = 1 : \frac{1}{\sqrt{2}} : \frac{1}{\sqrt{3}}$$

Laue Method:- In Laue's method, a single crystal is held stationary in a continuous x-ray beam. The crystal diffracts the discrete value of  $\lambda$  for which the crystal planes of spacing  $d$  and the incidence angle  $\theta$  satisfy the Bragg's equation.



Experimental arrangement is shown in figure. A continuous ( $0.2 \text{ \AA}$  to  $2 \text{ \AA}$ ) x-ray beam, well collimated by pinhole arrangement, is allowed to fall on a crystal. A film A is placed to receive the transmitted diffracted beam.

- \* Each spot in Laue pattern correspond to an interference maximum for a set of crystal plane satisfying the Bragg's equation  $2d \sin \theta = n\lambda$  for a wavelength selected from continuous beam.
- \* The distribution of spot depends on the symmetry of crystal and its orientation wrt x-ray beam.

Applications:- Used for rapid determination of crystal orientation and symmetry.

II) Used to study crystalline imperfection

Seven crystal systems:-

When a unit cell has lattice points only at corner called simple or primitive unit cell.

In all these are seven type of simple or primitive unit cell.

	Axial Distance	Axial Angle
i) Cubic	$a = b = c$	$\alpha = \beta = \gamma = 90^\circ$
ii) Tetrahedral	$a = b \neq c$	$\alpha = \beta = \gamma = 90^\circ$
iii) Orthorhombic	$a \neq b \neq c$	$\alpha = \beta = \gamma = 90^\circ$
iv) Monoclinic	$a \neq b \neq c$	$\alpha = \gamma = 90^\circ; \beta \neq 90^\circ$
v) Rhombohedral	$a = b = c$	$\alpha = \beta = \gamma \neq 90^\circ$
vi) Triclinic	$a \neq b \neq c$	$\alpha \neq \beta \neq \gamma \neq 90^\circ$
vii) Hexagonal	$a = b \neq c$	$\alpha = \beta = 90^\circ, \gamma = 120^\circ$

Crystal structure of NaCl:- (FCC)

The crystal of NaCl contains same  $\text{Na}^+$  &  $\text{Cl}^-$  ions both have coordination number of 6. Radius Ratio of NaCl is  $\frac{\text{Na}^+}{\text{Cl}^-} = 0.524$  which suggest that ~~NaCl~~  $\text{Cl}^-$  forms cubic closed packed structure in which  $\text{Na}^+$  ions occupy octahedral void in the lattice hence both have coordination number of 6:6.

6:6 ::  $\text{Na}^+:\text{Cl}^-$

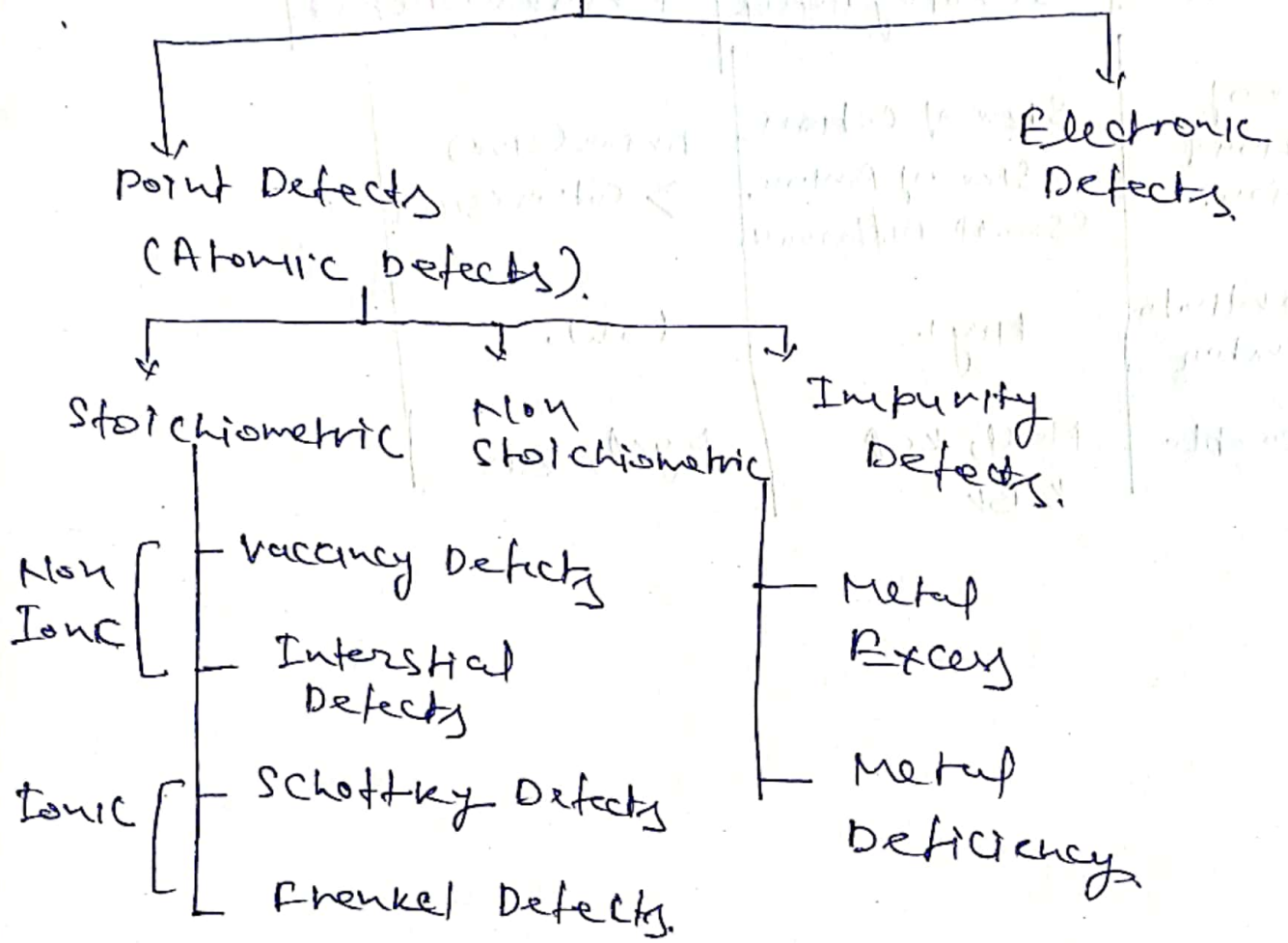
Crystal structure of Diamond:- (FCC)

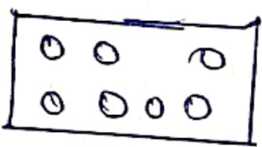
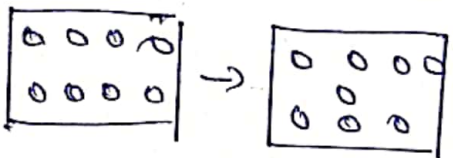
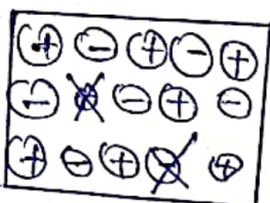
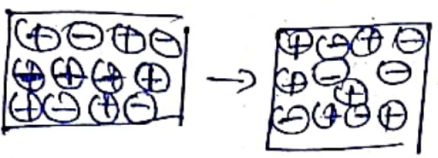
The space lattice of Diamond is face-centered cubic (FCC) with a basis of two carbon atoms associated with each lattice point.

In diamond each carbon atom forms covalent bond with four other carbon atoms, that occupy four corner of cube in tetrahedral structure.

length of each bond is  $1.54 \text{ \AA}$   
 Angle between bond is  $109.5^\circ$   
 Coordination number of crystal = 4  
 Number of Atom per unit cell = 8

Defects



Defects	Definition	Solids	Density	Structure
i) Vacancy	Some of lattice sites are vacant	Non-Ionic	Decrease	
ii) Interstitial	Some of the constituent atoms occupy interstitial site	Non-Ionic	Same	
iii) Schottky	Equal No. of cation & anion removed	Ionic	Decrease	
iv) Frenkel	one of the cation & anion occupy interstitial site	Ionic	Same	

S.No.	Condition	Schottky Defect	Frenkel Defect
I)	Size of Cation & Anion.	Size of Cation = Size of Anion. (Small difference)	Anion (Size) > Cation (Size)
II)	Coordination Number	High	Low.
III)	Example	NaCl, KCl, KBr	AgCl, Zn

Condition	Schottky Defect	Frenkel Defect	Condition	Schottky Defect	Frenkel Defect
Size of Cation & Anion	Size of Cation = Size of Anion. (Small difference)	Anion (Size) > Cation (Size)	Coordination Number	High	Low.
Example	NaCl, KCl, KBr	AgCl, Zn			

# Defects

Imperfection in Solids! At Absolute zero ( $0^{\circ}\text{K}$ ), there is no molecular motion, and therefore disorder is zero. As the temperature increase the random motion of particle takes place, due to which atoms leave their correct position & cause defects.

Point Defects! - mainly two types.

I) Stoichiometric Defects! - In this defect Ratio of cation and Anion ~~change~~ doesn't change.

Eg. Frenkel and Schottky Defects.

II) Non-Stoichiometric Defects! - Ratio of cation & Anion changes.

Eg. Metal Excess & Metal Deficiency Defects.

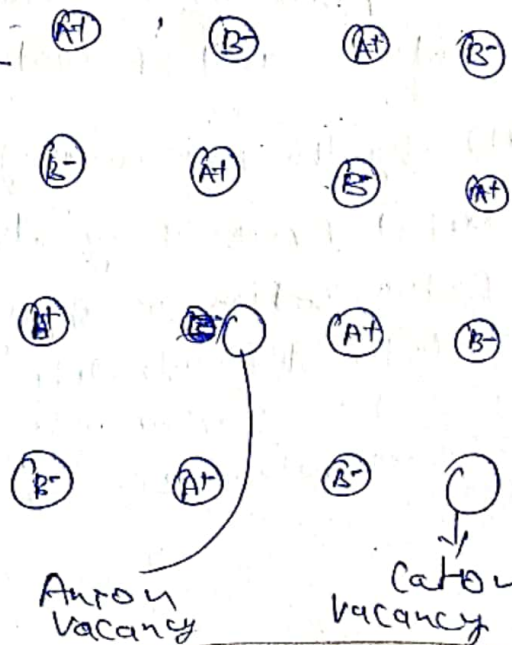
I) Schottky Defect! - when Equal Number of cation & Anion missed from lattice site

\* occur in those solids which has high co-ordination Number

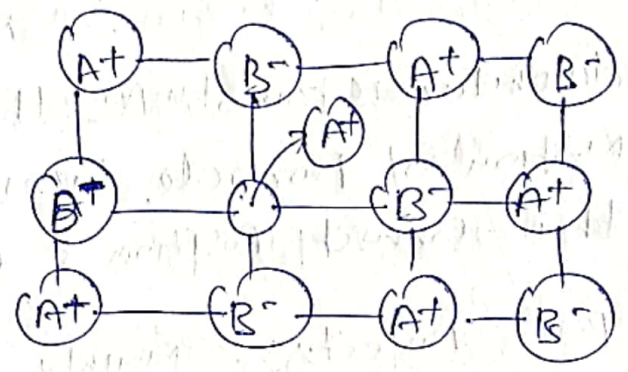
\* occur in those solids which has nearly similar size of cation & Anion.

Imp \* Density of crystal decrease during this defect.

Eg. NaCl, CsCl, KCl, KBr etc.



Frenkel defect!- This defect arises when an ion leaves its correct lattice position and occupies interstitial site.



→ Ionic compound remains electrically neutral & stoichiometry of compound also remains same.

→ occur in the crystal of low co-ordination number and large difference in size of ion.

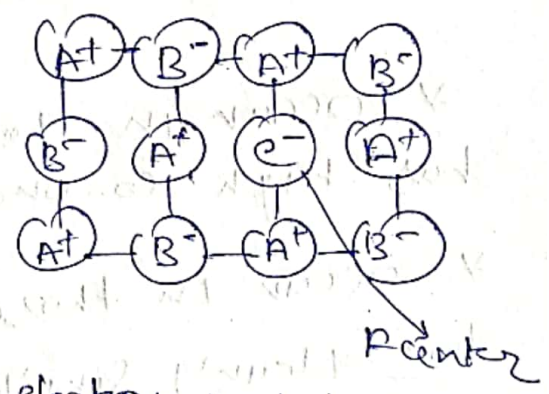
Eg. AgCl, AgBr, AgI, ZnS.

→ This defect don't change the density of crystal.  
\* AgBr shows both Schottky & Frenkel defect

Non-stoichiometric Defects:-

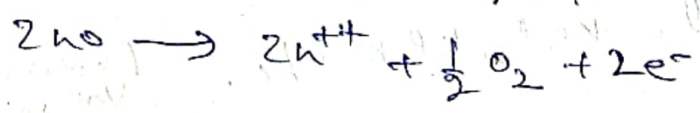
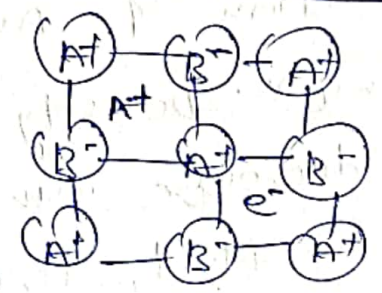
i) Metal Excess Defect (due to Anion vacancy)

In this defect the -ve ion is missed from lattice site. The hole caused by -ve ion is occupied by  $e^-$  to maintain the electrical neutrality. The holes occupied by  $e^-$  are called color centres and F centres. Eg. NaCl, KCl.

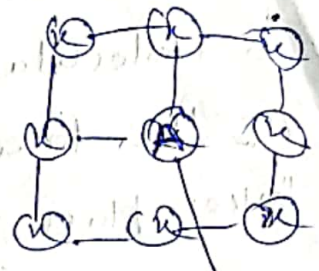


ii) By the presence of extra cation and electron in interstitial site

Metal excess may also be caused by an extra cation occupying the interstitial site. Electrical neutrality is maintained by an  $e^-$  present in another interstitial site. This defect occurs when Zn is heated in crystal.

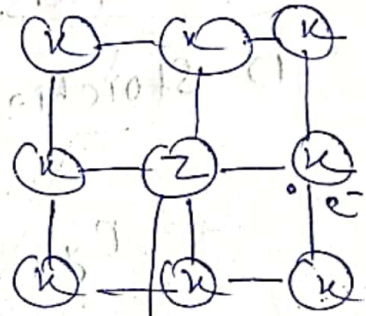


Impurity defect:- In covalent crystal when a crystal of 14 group is doped with impurity of 15 group then an extra  $e^-$  is present. This type of doping (impurity adding) give rise to a semiconductor of n-type.



Impurity Atom

ii) When a crystal of 14 group is doped with the impurity of 13 group then there is a deficiency of one  $e^-$  which breaks a hole in the crystal lattice. due to which current flow. This type of doping give rise to a p-type semiconductor.



hole

Note ( n type - Negative or  $e^-$  excess)  
 P type - Positive or  $e^-$  deficiency)

*(Faint, mostly illegible handwritten notes)*

- (1)
- (2)
- (3)
- (4)
- (5)
- (6)
- (7)
- (8)
- (9)
- (10)
- (11)
- (12)
- (13)
- (14)
- (15)
- (16)
- (17)
- (18)
- (19)
- (20)

## Band Theory of Solids!-

Band theory of Solids!- This is the theory of valence electron which are nearly free.

Free electron theory was successful in explaining the various Electronic & thermal properties of Metal like-

- I) Specific Heat
- II) Temperature Susceptibility
- III) Temperature dependence of Electrical Resistivity.

But it can't explain the distinction between conductor insulator & semiconductor.

Band theory explain that Each atom has a discrete set of Energy levels

1s, 2s, 2p, 3s, 3p -----

The splitted Energy levels are (becomes so numerous) so close together forming Energy bands.

Na- 1s<sup>2</sup> 2s<sup>2</sup> 2p<sup>6</sup> 3s<sup>1</sup>

3s - 3s' level is the first occupied level to be splitted into band.

2p - 2p level doesn't split until the interatomic distance becomes smaller.

OR Energy bands in solid correspond to Energy level in Atom.

\* Electrical properties in solid depend upon its Energy band structure.

Depending on the Nature of band occupation by electron and on the width of forbidden band, all the solids can be classified as conductors, insulator & semiconductor.

Valence band:- Filled level form a band known as valence band.

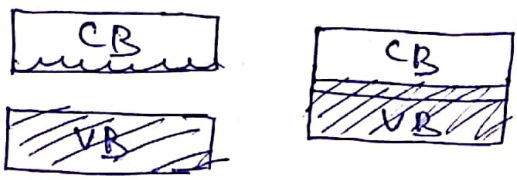
Conduction band:- Unfilled level form an another band known as conduction band.

Energy gap:- V.B & C.B separated by some energy gap known as forbidden Energy gap.

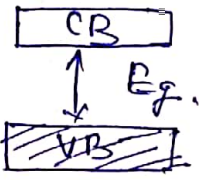
**Classification of Solids:-**

Solids can be classified into three types:-

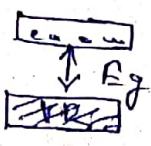
I) Metal or conductor:- C.B either partially filled or V.B & C.B partially overlap each other.



II) Insulator:- V.B is completely filled and C.B is completely empty & forbidden Energy gap is quite large.



III) Semiconductor:- At 0K in semiconductor V.B is completely & C.B is completely empty. Energy gap is smaller than insulator.



Semi conductors:-

A solid substance whose conductivity lies between an insulator and that of most metals.  
( Either due to addition of impurity or because of temperature effect).

Types of Semi conductors:-

- I) Intrinsic semiconductor II) Extrinsic Semiconductor
- I) Intrinsic Semiconductor:- A pure semiconductor is known as intrinsic semiconductor.

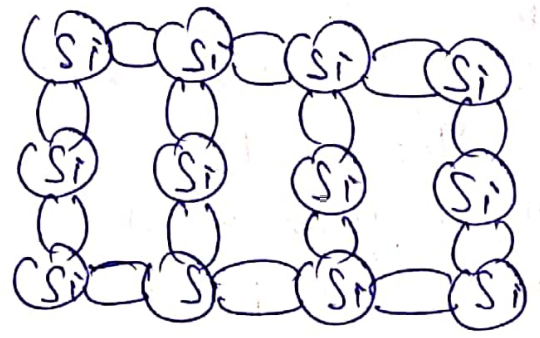
I) on the basis of Band theory:-

At 0K Semiconductors are insulators but at room temp. Some electron gain Energy due to thermal and moves from VB to CB.

- \* Vacancy is created in VB called Hole.
- \* Motion of hole in VB and electron in CB is responsible for Electronic conduction.

II) on the basis of crystal structure:- Atoms are strongly bonded by covalent bond.

→ on receiving Additional Energy some electron break their bond and produce free electron & equal No. of holes.



Thus effective current in semiconductor is the sum of Electronic current & Holic current.

## Extrinsic Semiconductor:-

Process of Addition of impurity to a Pure Semiconductor.

Doped Semiconductor called Extrinsic semiconductor.

### Types of Extrinsic Semiconductor:-

I) N-type Semiconductor      II) P-type Semiconductor.

I) N-type Semiconductor:- Intrinsic semiconductor doped with Pentavalent impurity atom.

1) On the basis of crystal structure

\* The impurity atom when mixed with Si or Ge atom, form covalent bond with neighbouring Si or Ge atom.

\* At room temperature some covalent bond break and produce equal number of holes and free electrons.

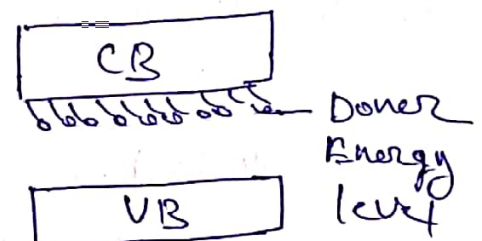
\* The number of electrons are large compare to holes.

\* Electrons are majority charge carrier and holes are minority charge carriers.

→ Conduction is mainly due to Electron + holes.

2) on the basis of Band theory

→ When impurity atom generate a new energy level just below the CB known as donor energy level.



→ Now it is easy for them to jump from VB to CB.

→ All electron of donor energy level jumps to CB at room temperature.

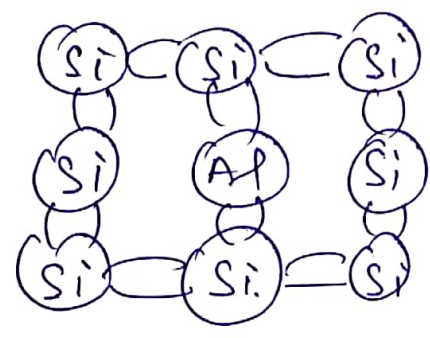
→ Some electron of VB also jumps to CB.

Hence the number of electron in CB is large as compare to VB.

# P-type Semiconductor:-

on the basis of crystal structure:-

→ When trivalent impurity atom is mixed with Si or Ge, they replace Si or Ge atom and form covalent bond with. Neighbouring Si or Ge.



→ Number of holes are very large as compare to Electrons. (due to trivalent impurity)

→ At room temperature covalent bond break and produce equal number of holes and electrons.

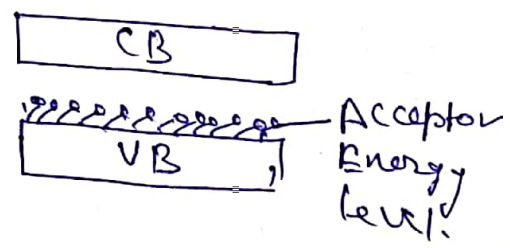
→ Holes are Majority and electron are Minority charge carrier.

⇒ conduction is mainly due to holes, therefore such semiconductor known as P-type Semiconductor.

on the basis of band theory:-

Impurity atom generate an energy level just above the top of valence band, this energy level known as acceptor level.

→ Thus electron from VB easily shift to acceptor level and generate holes in VB.



⇒ Electron from VB also goes to CB generating equal number of holes in VB.

→ Hence holes are Majority charge carrier electrons are Minority charge carrier in P-type Semiconductor.

### Difference between.-

#### I) Intrinsic Semi Conductor

- I) Pure Semi Conductor
- II)  $n_h = n_e$
- III) Conductivity is low
- IV) Conductivity depend only on temperature

#### Extrinsic Semi Conductor.

- I) Impure Semi Conductor
- 2)  $n_h \neq n_e$
- III) conductivity is High.
- IV) Conductivity depend on temperature and concentration of doping.

#### n-type Semi Conductor

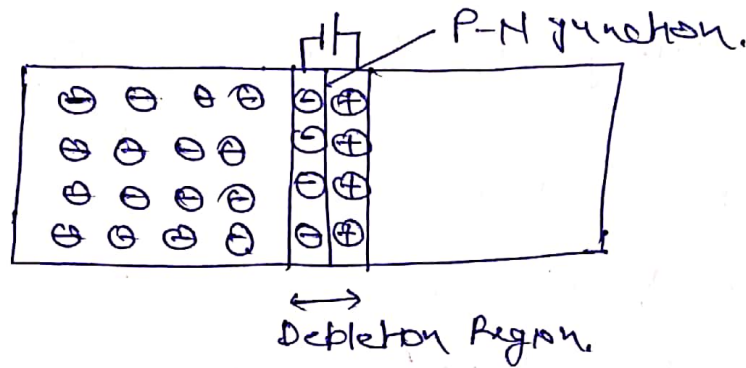
- I) Impurity atom is Pentavalent
- II) Electrons are Majority & Holes are Minority charge carriers.  $n_e > n_h$
- III) Impurity atom is known as donor atom
- IV) Donor level is formed just below the conduction band
- V) Conduction is mainly due to electrons

#### P-type Semi Conductor.

- I) Impurity Atom is Trivalent
- II) Electrons are Minority and holes are Majority charge carriers  $n_h > n_e$
- III) Impurity atom is known as acceptor Atom.
- IV) Acceptor level is formed just above the valence band
- V) Conduction is mainly due to holes.

P-N junction diode:- when P-type crystal is placed in contact with n-type crystal then common surface of contact of P & n type semiconductor is known as P-N junction diode.

Depletion region & Potential barrier:-



When a P-N junction is formed, holes move from p to n region and electrons move from n to p region through the junction (due to difference in concentration of charge carrier in two region).

\* The holes and electrons diffusing towards each other combine near the junction. Thus a layer is formed near the junction (with junction in the middle), which doesn't have any type of free charge carrier and is called depletion region.

\* Depletion region prevent the further movement of electron & holes through the junction.

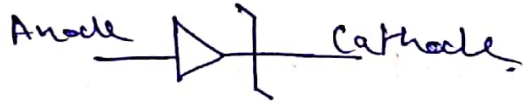
\* Thus a barrier is setup in P-N junction against the movement of charge carrier.

Note → D for Cu it is 0.3V (I) for Si it is 0.7V.

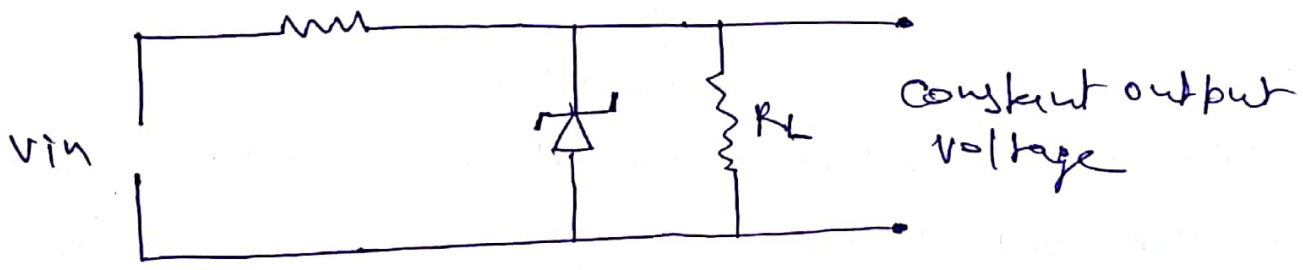
(II) Magnitude of potential barrier depend on nature of semiconductor & doping concentration.

~~The depletion region is about 0.5 cm. barrier for Si~~

Zener Diode:- This diode consists of special heavily doped P-n junction, designed to conduct in reverse direction when a certain specified voltage is reached.



Zener diode as voltage Regulator:- The circuit diagram using zener diode as voltage regulator is shown in Fig.

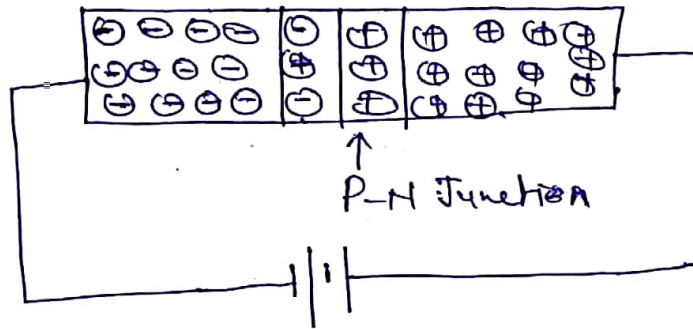


- When input voltage increase, resistance of zener diode decrease (due to breakdown of zener diode) and hence current through zener diode increase.
- AS a result of this large voltage drop occurs across dropping resistance  $R$ . Hence output voltage across  $R_L$  remains constant.
- when input voltage decrease resistance of zener diode increase and hence current through zener diode decrease. So small voltage drop takes place across  $R$ . Hence output voltage  $R_L$  maintained constant.

Biassing:- Applying an External Potential Difference on the faces of junction is called Biassing.

It is of two type -

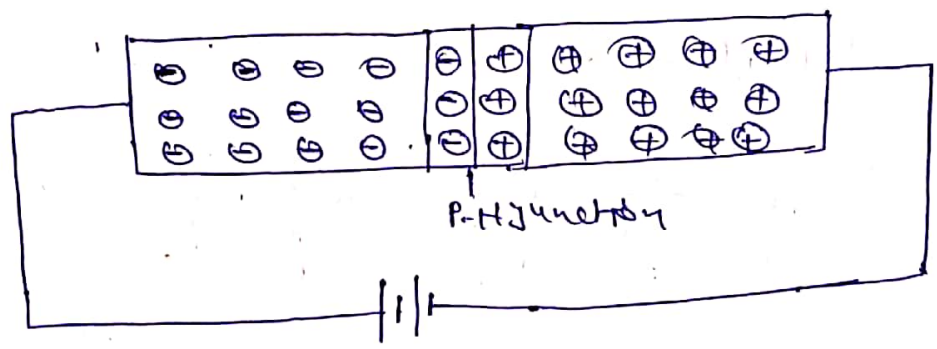
1) Forward Biassing:- In this biassing Negative terminal of the battery is connected to 'n' region and Positive terminal is connected to 'p' region.



Characteristics:-

- i) A current of mA order flows through the junction.
- ii) In Semiconductor current is due to both type of Majority charge carrier.
- iii) In forward biassing diode offer low resistance and width of depletion region decrease.
- iv) Resistance of P-N junction diode is Non-ohmic

11) Reverse Biasing:- In this biasing the terminal of the battery is connected to n-type Semiconductor and -ve terminal is connected to p-type Semiconductor.



Characteristics:- DA small current of  $\mu A$  order flows through the junction due to the movement of charge carrier.

- i) In external circuit current is due to free electrons.
- ii) In reverse biasing P-N junction offers high resistance and width of depletion region increase.
- iii) In reverse biasing current increase suddenly at a particular reverse biased voltage known as breakdown voltage. This phenomenon is known as avalanche breakdown. At this voltage diode may damage.

# P-N Junction (Energy band diagram with Fermi Energy)

## D P-N Junction (No bias)

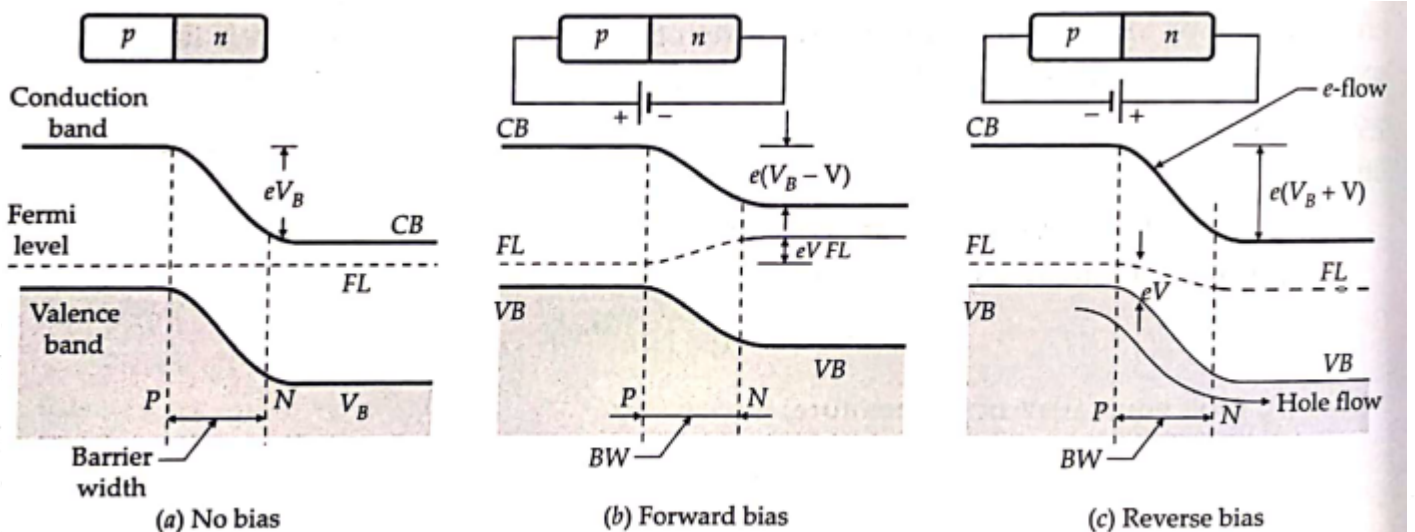
When P-N Junction is in Equilibrium, number of carrier diffusing from P side to N side is equal to number of carrier diffusing from N to P side. Consequently there is no current across junction.

## II) P-N Junction (Forward bias)

Fermi level rise by  $eV$  and their Energy band adjust their position so as to suit the elevation of Fermi level. Due to increase in Energy on N-side potential barrier is reduced to  $e(V_B - V)$ .

## III) P-N Junction (Reverse bias)

It lower the Fermi level on N-side by an amount of  $eV$  raising the barrier height  $e(V_B + V)$  and increasing the depletion width too. Consequently current is very much reduced in Reverse biasing of P-N Junction.



Tripointant  
Kronig Penny Model:-

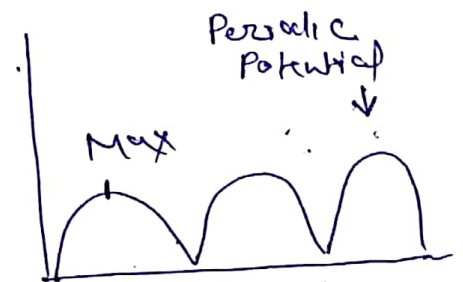
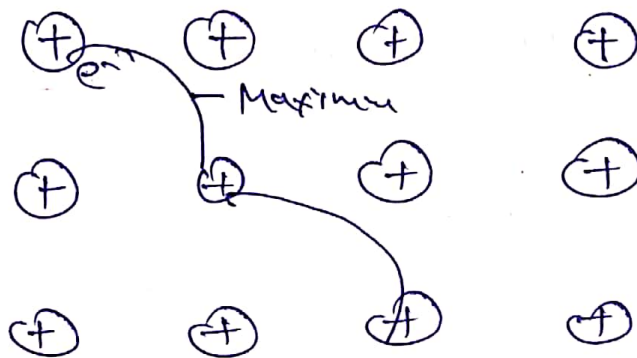
In order to find the allowed energy of electrons in solids, we must solve the Schrodinger Equation for an electron in ~~solids~~ a crystal lattice.

"Kronig & Penny suggested a simplified model consisting of an infinite row of rectangular potential well. Each well represent an approximation to the potential produce by one ion."

OR

Kronig Penny Model is a simplified model for electron in one dimensional periodic potential.

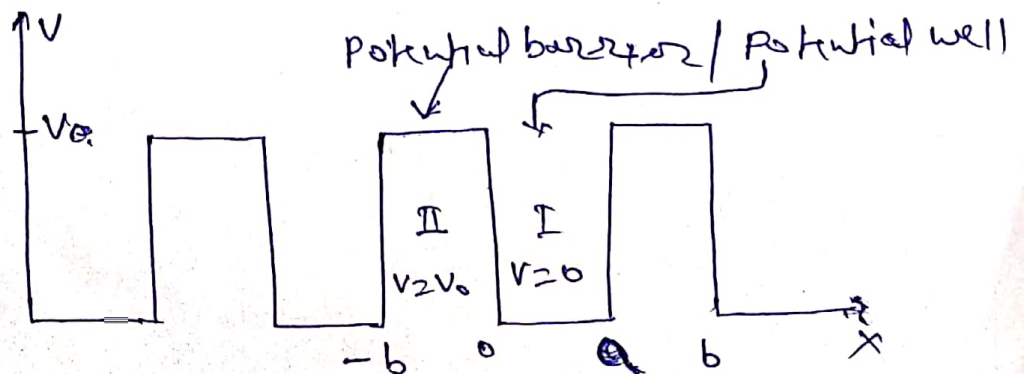
Let us consider a crystal lattice, let us discuss the variation of potential energy in a crystal lattice.



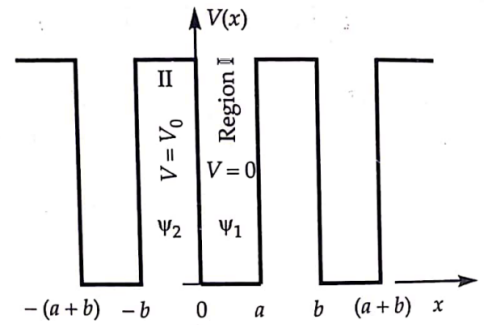
The highest potential is halfway between the atoms.

Kronig Penny consider these periodic potential as rectangular skp.

- Potential barrier potential,  $2 \times$  Maximum
- Potential well potential = 0



$$\left. \begin{array}{l} V=0, \quad 0 < x < a \\ V=V_0, \quad -b < x < 0 \end{array} \right\} \begin{array}{l} \text{Region-I} \\ \text{Region-II} \end{array}$$



For region I Schrodinger wave Eqn<sup>n</sup> is

$$\frac{d^2\psi}{dx^2} + \frac{2\pi^2m}{h^2} (E-V)\psi = 0$$

$$V=0 \Rightarrow \frac{d^2\psi}{dx^2} + \frac{2\pi^2mE\psi}{h^2} = 0 \quad \text{--- (1)}$$

For Region-II Schrodinger wave Eqn<sup>n</sup> is-

$$\frac{d^2\psi}{dx^2} + \frac{2\pi^2m(E-V_0)\psi}{h^2} = 0 \quad \text{--- (2)}$$

$$\text{Let } \alpha^2 = \frac{2\pi^2mE}{h^2} \quad \text{--- (3)} \quad \text{--- } \beta^2 = \frac{2\pi^2m(E-V_0)}{h^2} \quad \text{--- (4)}$$

Using (3), (4) in (1) & (2)

$$\frac{d^2\psi}{dx^2} + \alpha^2\psi = 0 \quad \text{--- (5)}$$

$$\frac{d^2\psi}{dx^2} - \beta^2\psi = 0 \quad \text{--- (6)}$$

$$\text{Also let } \psi = e^{ikx} \quad \text{--- (7)}$$

being a periodic function.

Now General solution of Equation (5) & (6)

will be-

$$\psi_1(x) = A e^{i(\alpha - k)x} + B e^{-i(\alpha + k)x} \quad \text{--- (8)}$$

$$\psi_2(x) = C e^{(\beta - ik)x} + D e^{-(\beta + ik)x} \quad \text{--- (9)}$$

Solving the above Equation we get

$$\boxed{\frac{P}{\alpha a} \sin \alpha a + \cos \alpha a = \cos ka} \quad \text{--- (10)}$$

↓  
Kronig - Penny Model Equation.

where  $P \rightarrow$  Strength with which electrons in a crystal are attracted to the ions on the crystal lattice.

$a$  - Interatomic distance.

Now CASE - I

when  $P = 0$

From (10)  $\cos \alpha a = \cos ka$ .

$$\Rightarrow \alpha a = ka$$

$$\alpha = k$$

$$\alpha^2 = k^2$$

$$\frac{8\pi^2 m E}{h^2} = \left(\frac{2\pi}{\lambda}\right)^2$$

$$\frac{8\pi^2 m E}{h^2} = \left(\frac{2\pi}{h/p}\right)^2$$

$$\boxed{E = \frac{p^2}{2m}}$$

\* value of  $E$  indicate that

Energy of Electron is kinetic Energy.

\* Hence Electron can move freely

\* It indicate Nature of Conductor

CASE - II  $P = \infty$

From (10)

$$\frac{\sin \alpha a}{\alpha a} + \frac{\cos \alpha a}{P} = \frac{\cos ka}{P}$$

$$\Rightarrow \sin \alpha a = 0$$

$$\sin \alpha a = \sin n\pi$$

$$\alpha a = n\pi$$

$$\alpha^2 = \frac{n^2 \pi^2}{a^2}$$

$$\left(\frac{8\pi^2 m E}{h^2}\right) = \frac{n^2 \pi^2}{a^2}$$

$$\boxed{E = \frac{n^2 h^2}{8ma^2}}$$

\*  $n$  - order,  $m$  - mass of  $e^-$

\* This Eqn of Energy is similar to potential well Equation

\* value of  $E$  indicate potential well.

\* Potential well indicate bound electron

\* It indicate Nature of Insulator.

Important  
Brillouin zones! - 1) In one dimensional lattice

Volume around the origin covered by Bragg's Plane called Brillouin zones.

(Note! - Bragg's Plane is perpendicular drawn between two lattice points).

OR

The Electron moving in a perfect periodic lattice potential lattice can have Energy only in the allowed region called zones.

OR

In the  $k$ -space ( $k$  is Propagation constant), the graph between Energy ( $E$ ) &  $k$  is discontinuous, The volume between these discontinuities is known as Brillouin zones.

These discontinuities in one dimensional lattice occurs when  $k = \frac{n\pi}{a}$ , where  $n$  is Negative or Positive integer.

$$k = \pm \frac{\pi}{a}, \pm \frac{2\pi}{a}, \pm \frac{3\pi}{a}, \pm \frac{4\pi}{a} \dots$$

The first reflection and the first Energy gap occurs at  $k = \pm \frac{\pi}{a}$  because the wave reflected from one atom in a linear lattice interfere constructively with the nearest neighbouring atom,

Then the segment  $-\frac{\pi}{a} \leq k < +\frac{\pi}{a}$  called first Brillouin zone.

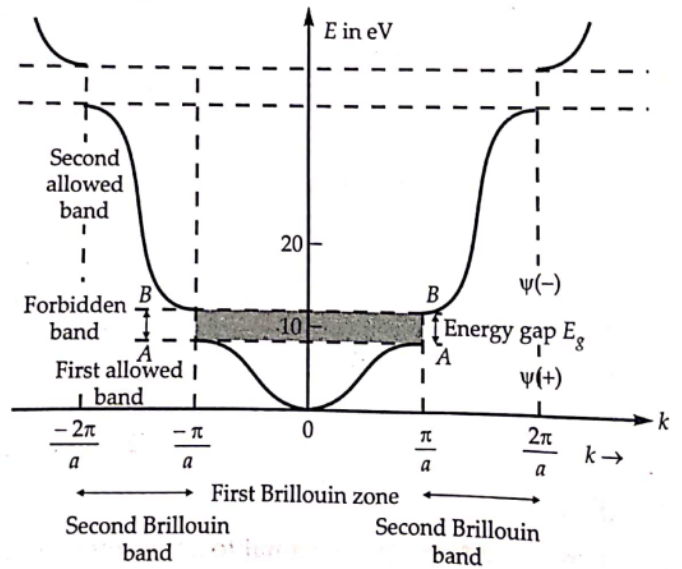
The second Reflection and second Energy gap occurs at  $k = \pm \frac{2\pi}{a}$

$\therefore$   $k$  space between  $-\frac{2\pi}{a}$  to  $-\frac{\pi}{a}$  and  $+\frac{\pi}{a}$  to  $+\frac{2\pi}{a}$

OR  $-\frac{2\pi}{a} \leq k < -\frac{\pi}{a}$  and  $+\frac{\pi}{a} \leq k < +\frac{2\pi}{a}$

Called second Brillouin zone.  $\dots$  etc

The first and second Brillouin zone and energy gap are shown by plotting value of  $E$  vs  $k$  as shown below.



The first and second BZ and energy gap.

## ii) Two dimensional lattice.

Consider  $x-y$  plane be the square ABCD.

The boundaries of square is given by -

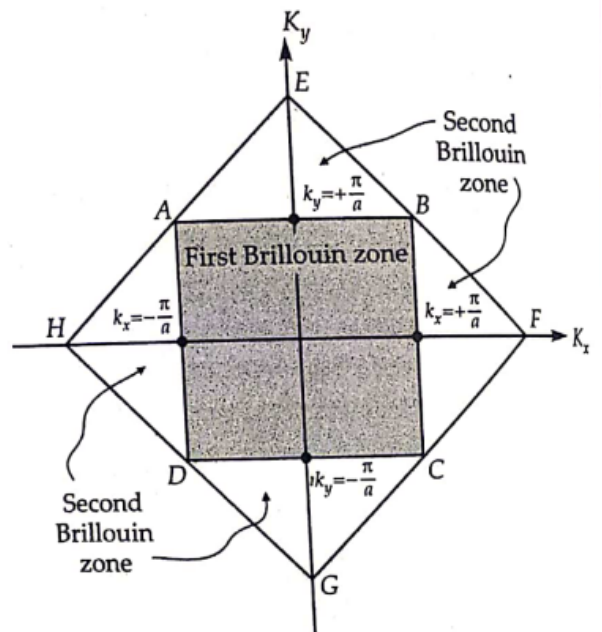
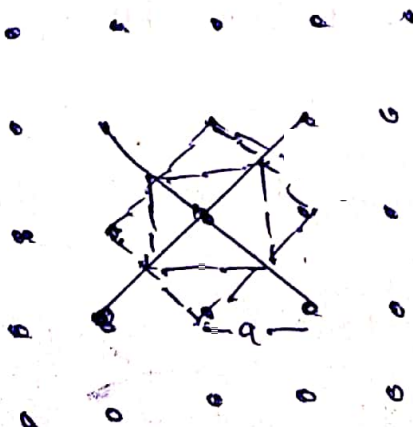
$$k_x = \pm \frac{\pi}{a}, \quad k_y = \pm \frac{\pi}{a}$$

is first Brillouin zone.

Now the boundaries of second square EFGH is given by -

$$k_x = \pm \frac{2\pi}{a}, \quad k_y = \pm \frac{2\pi}{a}$$

is second Brillouin zone.



First and second BZ in 2D lattice.

# Effective Mass! → Important

Mass exhibited by an electron, when inside the Semiconductor is called Effective Mass.

OR

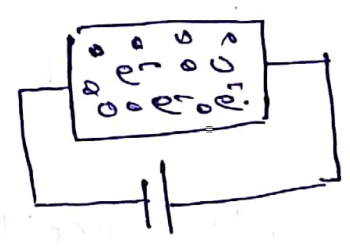
Deviation of electron behaviour in the crystal lattice from free electron behaviour can be taken into account by considering the electron to have Effective Mass  $m^*$ .

$$\frac{1}{m^*} = \frac{4\pi^2}{h^2} \frac{d^2E}{dk^2}$$

$m^* \rightarrow$  Effective Mass.

Electron moving inside the Semiconductor is moving under the influence of atom and other electrons.

Due to this Electrostatic field act on electron known as crystal field. This is called internal Electric field.



This result effective mass of  $e^-$  ( $m^*$ ).

Derivation → According to quantum theory  $e^-$  moving with velocity  $v$  is equivalent to wave packet (group of wave) moving with  $v_g$  (group velocity). i.e.

$$\Rightarrow v = v_g = \frac{d\omega}{dk} \quad \text{--- (1)}$$

$$\text{Also } E = h\nu = \frac{h}{2\pi} \times 2\pi\nu \Rightarrow \boxed{E = \hbar\omega} \quad \text{--- (2)}$$

Diff (2) wrt  $k$  we get.

$$\frac{dE}{dk} = \hbar \frac{d\omega}{dk} \Rightarrow \frac{d\omega}{dk} = \frac{1}{\hbar} \frac{dE}{dk} \quad \text{--- (3)}$$

$$\text{Comparing (1) \& (3) we get } v = \frac{1}{\hbar} \frac{dE}{dk} \quad \text{--- (4)}$$

Now according to de Broglie hypothesis.

$$\lambda = \frac{h}{p}$$

$$p = \frac{h}{\lambda} \Rightarrow \frac{h}{\lambda} \times \frac{2\pi}{2\pi}$$

$$\Rightarrow p = \frac{h}{2\pi} \times \frac{2\pi}{\lambda} = \hbar k$$

$$p = \hbar k$$

Diff wrt  $t$

$$\frac{dp}{dt} = \hbar \frac{dk}{dt}$$

$$\Rightarrow F = \hbar \frac{dk}{dt}$$

(By Newton  
Second Law  $F = \frac{dp}{dt}$ )

$$\Rightarrow \hbar \frac{dk}{dt} = F = eE \quad \text{--- (5)}$$

By (4) Now Diff (4) wrt  $t$  we get.

$$v = \frac{1}{\hbar} \frac{dE}{dk} \Rightarrow \frac{dv}{dt} = \frac{1}{\hbar} \frac{d^2E}{dk^2} \frac{dk}{dt} \quad \text{--- (6)}$$

$$\text{From (5)} \quad \frac{dk}{dt} = \frac{F}{\hbar}$$

Also Acceleration  $a = \frac{dv}{dt}$

$$\text{From (6)} \quad a = \frac{1}{\hbar} \frac{d^2E}{dk^2} \frac{dk}{dt}$$

$$a = \frac{1}{\hbar} \frac{d^2E}{dk^2} \frac{F}{\hbar}$$

$$\frac{a}{F} = \frac{1}{\hbar} \frac{d^2E}{dk^2}$$

Taking  $a/F = m^*$  we get

$$\boxed{m^* = \frac{1}{\hbar} \frac{d^2E}{dk^2}} \quad \text{--- Effective mass}$$

If  $m^* = m$  then electron in crystal behaves as a free electron.

**Example 6.3** The energy near the valence band edge of a crystal is given by  $E = -10^{-39} k^2 \text{ Jm}^2$ .

An electron with wave vector  $10^{10} \hat{k}_x \text{ m}^{-1}$  is removed from an orbital in the completely filled valence band.

Determine effective mass and momentum.

**Solution.** Given :  $E = -10^{-39} k^2$  . Hence  $\frac{d^2 E}{dk^2} = -2 \times 10^{-39}$

Effective mass of the electron

$$m^* = \hbar^2 / \frac{\partial^2 E}{\partial k^2} = - \frac{(1.05 \times 10^{-34})^2}{2 \times 10^{-39}} = -5.5 \times 10^{-30} \text{ kg}$$

$$\text{Momentum of electron} = \hbar k = 1.05 \times 10^{-34} \times 10^{10} \hat{k}_x = 1.05 \times 10^{-24} \hat{k}_x \text{ Js m}^{-1}.$$

### Comparison between Ordinary Diode and Zener Diode

S.No.	Ordinary Diode	Zener Diode
1.	Ordinary diode is operated in forward biased condition.	Zener diode is always operated in reverse biased condition.
2.	It utilizes the forward characteristics.	It utilizes, the reverse characteristics. In forward bias it acts like a PN junction diode.
3.	It does not have any sharp breakdown voltage.	If properly doped, it has a sharp breakdown $V_Z$ .

## Fermi level:-

The Maximum Energy that electron may possess at 0K is the Fermi Energy ( $E_F$ ).

"Fermi level is the surface of the sea of electron at Absolute zero, where no electron will have enough energy to rise above the surface."

Now in order to know how many of Electronic Energy State in the valence band and conduction band will be occupied at different temperature Fermi-factor  $f(E)$  is introduced as-

$$f(E) = \frac{1}{e^{(E-E_F)/k_B T} + 1}$$

where  $k_B$  - Boltzmann constant

$T$  - Temperature in K

$E_F$  - Fermi level of Energy in eV.

$E$  - <sup>given</sup> Energy occupied by Electron under condition of thermal equilibrium.

Now

i)  $E > E_F$  and  $T = 0$   $f(E) = \frac{1}{1 + \exp(\infty)} = 0$

ii)  $E < E_F$  and  $T = 0$   $f(E) = \frac{1}{1 + \exp(-\infty)} = 1$

iii)  $T = T_K$  and  $E = E_F$

$$f(E) = \frac{1}{e^0 + 1} = \frac{1}{2}$$

which means when temperature is not

which means when temperature is not ok but some higher value say  $T=1000\text{K}$  then some electrons covalent bond will break and some electron will transfer to conduction band.

**Example 6.7** Fermi energy of an intrinsic semiconductor is  $0.6\text{ eV}$ . The low lying energy level in the conduction band is  $0.2\text{ eV}$  above the Fermi level. Calculate the probability of occupation of this level by an electron at room temperature.

**Solution.** Given  $E_F = 0.6\text{ eV}$ , then  $E = (0.6 + 0.2)\text{ eV} = 0.8\text{ eV}$

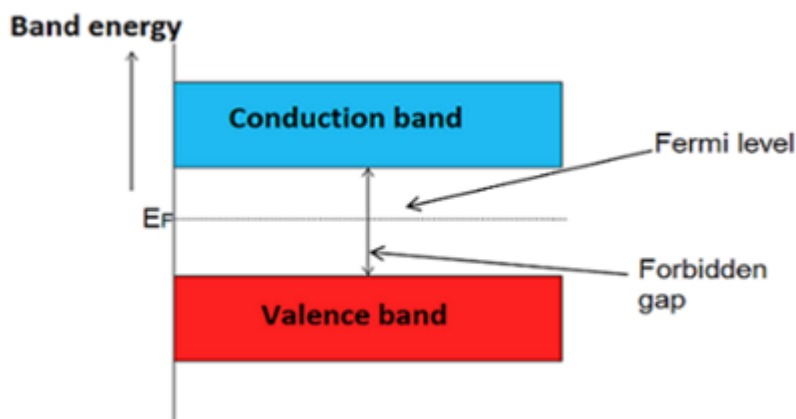
The probability of occupation of an energy level by an electron is given by the F-D distribution law

$$f(E) = \frac{1}{1 + \exp\left(\frac{E - E_F}{k_B T}\right)} = \frac{1}{1 + \exp\left(\frac{(0.8 - 0.6)\text{ eV} \times 1.6 \times 10^{-19}\text{ J}}{1.38 \times 10^{-23}\text{ JK}^{-1} \times 300\text{ K}}\right)}$$

$$= \frac{1}{1 + \exp 7.7} = 0.0004 = 0.04\%$$

### Fermi level in Intrinsic Semiconductor!

- \* In Intrinsic or pure Semiconductor, the Number of holes in valence band is equal to number of electrons in conduction band.
- \* The Probability of occupation of Energy levels in conduction band and valence band are equal.
- \* Fermi level for intrinsic semiconductor lies in the middle of forbidden band.



Fermi level for intrinsic semiconductor is given as -

$$E_F = \frac{E_C + E_V}{2}$$

where  $E_F$  - Fermi level  
 $E_C$  - Conduction band.  
 $E_V$  - valence band.

### Fermi level in extrinsic semiconductor:-

- \* In. extrinsic semiconductor, the number of electron in conduction band and number of holes in valence band are not equal.
- \* Hence, the probability of occupation of energy in conduction band and valence band are not equal.
- \* Fermi level for extrinsic semiconductor lies close to the conduction or valence band.

### D) Fermi level in n-type semiconductor:-

- \* At room temperature, the number of electrons in conduction band is greater than number of holes in valence band.
- \* Probability of occupation of energy levels by electron in the conduction band is greater than the probability of occupation of energy levels by holes in valence band.
- \* Therefore the Fermi level in n-type semiconductor lies close to conduction band.

Fermi level for n-type Semiconductor is given as -

$$E_F = E_C - k_B T \log \frac{N_C}{N_D}$$

$E_F$  - Fermi level

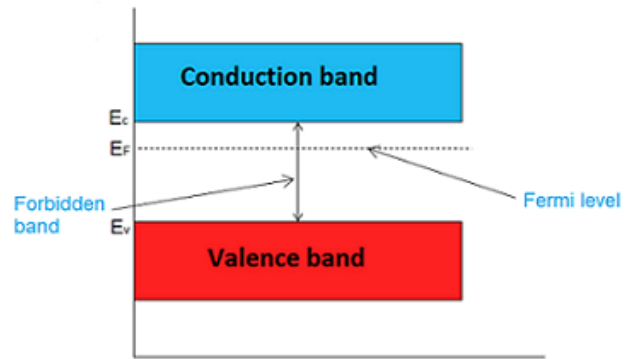
$E_C$  - Conduction band

$k_B$  - Boltzmann Constant

$T$  - Absolute Temperature

$N_C$  - Effective density of state in CB

$N_D$  - Concentration of Donor atom.

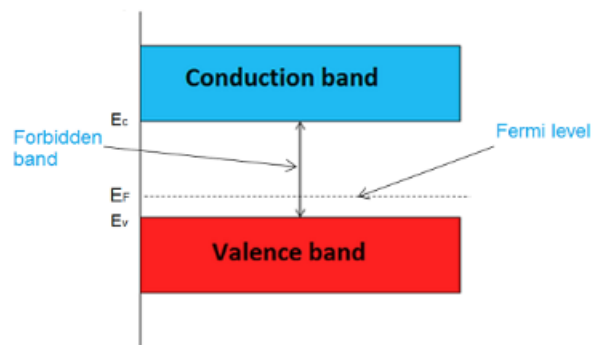


Fermi level in p-type Semiconductor:-

\* Number of holes in valence band is greater than the number of electron in conduction band.

\* Probability of occupation of Energy levels by the holes in the valence band is greater than the probability of occupation of Energy level in the conduction band.

\* Therefore the Fermi level in p-type Semiconductor lies close to valence band.



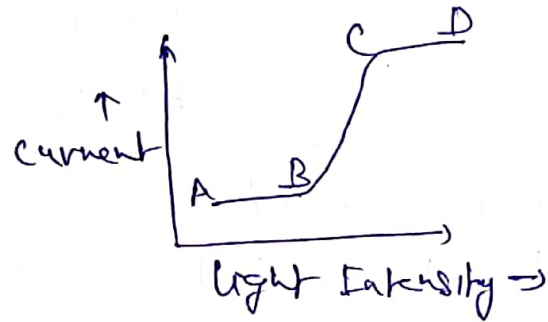
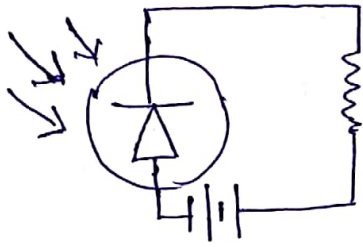
The Fermi level for p-type Semiconductor is given as -

$$E_F = E_V + k_B T \log \frac{N_V}{N_A}$$

$N_V$  - effective density of state in VB

$N_A$  - Concentration of Acceptor ion.

Photo diode  $\rightarrow$  A photodiode is a p-n junction diode, (13)  
 made from photosensitive semiconducting material, whose conductivity is controlled by light falling on it.



It can work in two forms: photodetector as well as solar cell.

Basically, photodiode is a p-n junction diode designed to respond under photon absorption.

There will be absorption only if  $h\nu > E_g$

LED  $\rightarrow$  (Light Emitting Diode)

LED is forward bias p-n junction diode.

If we apply forward bias then, electron-hole pair recombine near the junction, hence recombination energy emitted in the form of photon of energy  $h\nu = E_g$

$\rightarrow$  Here we use direct band gap semiconductor such that most of light must be emitted in the form of photon. Energy of photon depends on the band gap of semiconductor.

$$E_g = h\nu$$

$$E_g = \frac{hc}{\lambda} = \frac{1240 \text{ eV}}{\lambda \text{ (nm)}}$$

\* LED emit incoherent light.

Parameters of LED:-

D Quantum Efficiency of LED:- It is defined as

$$\Phi.E = \frac{\text{Optical Power out}}{\text{Electrical Power in}} \leq 1$$

$\Phi.E (N) = 80\%$  means 80% of electrical power convert into optical power.

Here power is Energy Per second

$$E \text{ (Energy/Per sec)} = nh\nu$$

$n \Rightarrow$  no. of photons per sec.

By this formula we can calculate no. of photons per sec.

1) Slope Efficiency:- It is defined as optical power out divided by input current.

$$\text{Slope efficiency} = \frac{\text{Optical Power out}}{\text{Input current}} \left( \frac{W}{A} \right)$$

Tunnel Diode:- A tunnel diode is a heavily doped Pn junction diode



- Tunnel diode shows negative resistance, so it can be used as amplifier, oscillators and in any switching circuits.
- When voltage value increase, current flow decrease
- Tunnel diode is based on tunnel effect
- Heavily doped P-type & n-type semiconductor in tunnel diode results in a narrow depletion region.